

### LJMU Research Online

Lange, R, Moffett, AJ, Driver, SP, Robotham, ASG, Lagos, CDP, Kelvin, LS, Conselice, C, Margalef-Bentabol, B, Alpaslan, M, Baldry, I, Bland-Hawthorn, J, Bremer, M, Brough, S, Cluve, M, Colless, M, Davies, LJM, Häußler, B, Holwerda, BW, Hopkins, AM, Kafle, PR, Kennedy, R, Liske, J, Phillipps, S, Popescu, CC, Taylor, EN, Tuffs, R, Kampen, EV and Wright, AH

Galaxy And Mass Assembly (GAMA): M-star-R-e relations of z=0 bulges, discs and spheroids

http://researchonline.ljmu.ac.uk/id/eprint/3896/

#### **Article**

**Citation** (please note it is advisable to refer to the publisher's version if you intend to cite from this work)

Lange, R, Moffett, AJ, Driver, SP, Robotham, ASG, Lagos, CDP, Kelvin, LS, Conselice, C, Margalef-Bentabol, B, Alpaslan, M, Baldry, I, Bland-Hawthorn, J, Bremer, M, Brough, S, Cluve, M, Colless, M, Davies, LJM, Häußler, B, Holwerda. BW. Hopkins. AM. Kafle. PR. Kennedv. R. Liske. J. Phillipps. S.

LJMU has developed LJMU Research Online for users to access the research output of the University more effectively. Copyright © and Moral Rights for the papers on this site are retained by the individual authors and/or other copyright owners. Users may download and/or print one copy of any article(s) in LJMU Research Online to facilitate their private study or for non-commercial research. You may not engage in further distribution of the material or use it for any profit-making activities or any commercial gain.

The version presented here may differ from the published version or from the version of the record. Please see the repository URL above for details on accessing the published version and note that access may require a subscription.

For more information please contact  $\underline{researchonline@ljmu.ac.uk}$ http://researchonline.ljmu.ac.uk/

# Galaxy And Mass Assembly (GAMA): $\mathcal{M}_{\star} - R_{\rm e}$ relations of z = 0 bulges, discs and spheroids

Rebecca Lange<sup>1\*</sup>, Amanda J. Moffett<sup>1</sup>, Simon P. Driver<sup>1,2</sup>, Aaron S.G. Robotham<sup>1</sup>, Claudia del P. Lagos<sup>1,3</sup>, Lee S. Kelvin<sup>4</sup>, Christopher Conselice <sup>5</sup>,

Berta Margalef-Bentabol<sup>5</sup>, Mehmet Alpaslan <sup>6</sup>, Ivan Baldry<sup>4</sup>, Joss Bland-Hawthorn<sup>7</sup>,

Malcolm Bremer<sup>8</sup>, Sarah Brough<sup>9</sup>, Michelle Cluver<sup>10</sup>, Matthew Colless<sup>11</sup>,

Luke J. M. Davies<sup>1</sup>, Boris Häußler<sup>12</sup>, Benne W. Holwerda<sup>13</sup>, Andrew M. Hopkins<sup>7</sup>, Prajwal R. Kafle<sup>1</sup>, Rebecca Kennedy<sup>5</sup>, Jochen Liske<sup>14</sup>, Steven Phillipps<sup>8</sup>,

Cristina C. Popescu<sup>15,16</sup>, Edward N. Taylor<sup>17</sup>, Richard Tuffs<sup>18</sup>, Eelco van Kampen<sup>19</sup>,

Angus H. Wright<sup>1</sup>

- <sup>1</sup> International Centre for Radio Astronomy Research, University of Western Australia, M468, 35 Stirling Highway, Crawley, WA6009, Australia <sup>2</sup> Scottish Universities' Physics Alliance (SUPA), School of Physics and Astronomy, University of St. Andrews, North Haugh, St. Andrews, KY16 9SS, UK
- <sup>3</sup> Australian Research Council Centre of Excellence for All-sky Astrophysics (CAASTRO), 44 Rosehill Street, Redfern, NSW2016, Australia
- <sup>4</sup> Astrophysics Research Institute, Liverpool John Moores University, IC2, Liverpool Science Park, 146 Brownlow Hill, Liverpool, L3 5RF, UK
- <sup>5</sup> School of Physics and Astronomy, The University of Nottingham, University Park, Nottingham, NG7 2RD, UK
- <sup>6</sup>NASA Ames Research Center, N232, Moffett Field, Mountain View, CA 94035, United States

  <sup>7</sup>Sudaya Institute for Astronomy, School of Physics, A28, University of Sudaya, NSW, 2006, Austrop.
- <sup>7</sup> Sydney Institute for Astronomy, School of Physics A28, University of Sydney, NSW 2006, Australia
- <sup>8</sup> Astrophysics Group, School of Physics, University of Bristol, Bristol BS8 1TL, UK
- <sup>9</sup> Australian Astronomical Observatory, PO Box 915, North Ryde, NSW 1670, Australia
- <sup>10</sup>Department of Physics and Astronomy, University of the Western Cape, Robert Sobukwe Road, Bellville, 7535, South Africa
- <sup>11</sup>Research School of Astronomy and Astrophysics, Australian National University, Canberra, ACT 2611, Australia
- <sup>12</sup>European Southern Observatory, Alonso de Cordova 3107, Vitacura, Santiago, Chile
- <sup>13</sup> University of Leiden, Sterrenwacht Leiden, Niels Bohrweg 2, NL-2333 CA Leiden, The Netherlands
- <sup>14</sup>Universität Hamburg, Hamburger Sternwarte, Gojenbergsweg 112, D-21029 Hamburg, Germany
- <sup>15</sup> Jeremiah Horrocks Institute, University of Central Lancashire, Leighton Building, Preston PR1 2HE, United Kingdom
- <sup>16</sup> The Astronomical Institute of the Romanian Academy, Str. Cutitul de Argint 5, Bucharest, Romania
- <sup>17</sup>School of Physics, the University of Melbourne, VIC 3010, Australia
- <sup>18</sup> Max-Planck-Institut für Kernphysik, Saupfercheckweg 1, D-69117 Heidelberg, Germany
- <sup>19</sup> European Southern Observatory, Karl-Schwarzschild-Str.2, D-85748 Garching, Germany

6 July 2016

#### ABSTRACT

We perform automated bulge + disc decomposition on a sample of ~7500 galaxies from the Galaxy And Mass Assembly (GAMA) survey in the redshift range of 0.002 < z < 0.06 using SIGMA, a wrapper around GALFIT3. To achieve robust profile measurements we use a novel approach of repeatedly fitting the galaxies, varying the input parameters to sample a large fraction of the input parameter space. Using this method we reduce the catastrophic failure rate significantly and verify the confidence in the fit independently of  $\chi^2$ . Additionally, using the median of the final fitting values and the  $16^{th}$  and  $84^{th}$  percentile produces more realistic error estimates than those provided by GALFIT, which are known to be underestimated.

We use the results of our decompositions to analyse the stellar mass – half-light radius relations of bulges, discs and spheroids. We further investigate the association of components with a parent disc or elliptical relation to provide definite z=0 disc and spheroid  $\mathcal{M}_{\star}-R_{\rm e}$  relations. We conclude by comparing our local disc and spheroid  $\mathcal{M}_{\star}-R_{\rm e}$  to simulated data from EAGLE and high redshift data from CANDELS-UDS. We show the potential of using the mass-size relation to study galaxy evolution in both cases but caution that for a fair comparison all data sets need to be processed and analysed in the same manner.

**Key words:** galaxies: fundamental parameters - galaxies: statistics - galaxies: formation - galaxies: elliptical and lenticular - galaxies: spiral

#### 1 INTRODUCTION

At the fundamental level galaxies are multi-component systems (see for example Buta et al. 2010), consisting of at least a spheroid and/or disc. This is most obvious in the Sérsic index – colour plane where the single component Sd and elliptical galaxies occupy distinct peaks with composite galaxies (S0abc) scattered between and around these peaks (see e.g., Driver et al. 2006; Cameron et al. 2009; Kelvin et al. 2012; Lange et al. 2015). These components have very different characteristics with spheroids typically having a featureless appearance and being pressure supported. Discs on the other hand have features such as spiral arms and are rotationally supported. Furthermore bulges are made up of redder stars with moderate to high metallicities and a high  $\alpha$ -element abundance, while discs are made of younger, bluer stars with lower metallicities and typically are dust and gas rich. Spheroids are older, showing little to no star formation and are typically dust and gas depleted (see for example the review by Roberts & Haynes 1994). The simplest explanation for these stark differences is that spheroids and discs form via two distinct mechanisms over two distinct eras (Cook et al. 2009; Driver et al. 2013), i.e. a dynamically "hot mode" (spheroid formation) and "cold mode" (disc formation) evolution.

Traditionally the relative prominence of a bulge component is taken into account when classifying galaxies onto the Hubble sequence (see Hubble 1926, and later revisions by e.g., van den Bergh 1976; Kormendy & Bender 2012), however studying global properties of galaxies by Hubble type could be misleading. For example, numerous evolution mechanisms have been proposed to explain the morphological diversity seen at z=0, such as a (initial) major dissipative event, gas accretion, adiabatic contraction, major and minor mergers and secular processes (see e.g., Hopkins et al. 2010; Trujillo et al. 2011; L'Huillier et al. 2012; Cheung et al. 2013; Sachdeva et al. 2015). Each of these processes potentially acts to modify the prominence of the bulge, disc or other components. This indicates that galaxy components likely follow distinct formation pathways and structure effectively encodes the formation history. Therefore, to study galaxy evolution bulge + disc decomposition is critical.

While the number of studies of large samples which employ bulge + disc decomposition to explore the nature of galaxies and their components is growing, the analysis is challenging (see e.g., Allen et al. 2006; Gadotti 2009; Simard et al. 2011; Bruce et al. 2012; Bruce et al. 2014; Lang et al. 2014; Tasca et al. 2014; Meert et al. 2015; Salo et al. 2015). This is because multi-component fitting is notoriously difficult, especially when trying to automate it for large samples. Nevertheless a number of publicly available codes have now been created to allow bulge + disc decomposition, such as GIM2D (Simard 1998), BUDDA (de Souza et al. 2004), GAL-FIT3 (Peng et al. 2010) and IMFIT (Erwin 2015). Each code has advantages and disadvantages (see Erwin 2015, for example for further discussion), here we elect to use GALFIT3 because of its ability to manage nearby objects, its computational reliability, and its speed.

Many studies that fit 2-component Sérsic light profiles restrict the Sérsic index to n=1 for the disc and in some cases n=4 for the bulge (e.g. Simard et al. 2011; Bruce et al. 2012;

Lackner & Gunn 2012; Meert et al. 2015). This reduces the number of free parameters and ensures the fitting process is more robust but it restricts the possible interpretations of the fitting outcomes, e.g. classical and pseudo bulges can not be differentiated this way. A number of studies now show that the Sérsic index of discs and spheroids (be they pure or component) vary smoothly with mass and luminosity or due to dust or galaxy type (see e.g. Graham & Guzman 2003; Gadotti 2009; Kelvin et al. 2012; Graham 2013; Pastrav et al. 2013a,b). Hence studies where the Sérsic index of the bulge or disc components are fixed may be overly restrictive. Furthermore, to correctly trace a galaxy's formation history a full decomposition of all of its components would be ideal (e.g., the Spitzer Survey of Stellar Structure in Galaxies, S<sup>4</sup>G, Salo et al. 2015). However, this is only viable for very nearby galaxies where all the components can be clearly resolved and hence for relatively small samples (S<sup>4</sup>G is the largest study to date extending to 2352 galaxies for which a number have been fit with more than 2 components). To compare to galaxies at different epochs going beyond a simple bulge and disc decomposition is difficult (Gadotti 2008). There are two reasons, however, why two components might be sufficient, (i) the majority of stellar mass resides in the bulge and disc components for most galaxies, and (ii) some components may simply represent minor perturbations to the underlying disc (e.g. bars, pseudobulges). Such perturbations should arguably be considered secondary rather than primary evolutionary markers.

Here we adopt the stance that bulge and disc components arise from two primary formation pathways (i.e. hot and cold mode evolution, respectively), and that additional components form in secondary formation pathways (i.e., tidal interactions, disc instabilities and perturbations). The likely primary pathways are: monolithic collapse followed by major mergers, which can produce elliptical galaxies by destroying and rearranging any structure previously present in a galaxy, resulting in a smooth light profile (Toomre 1977); and minor mergers and continued gas inflow, which can form or re-grow a disc around a pre-existing spheroid, resulting in a galaxy with two distinct components (see e.g., Steinmetz & Navarro 2002; Kannappan et al. 2009; Wei et al. 2010). A key question worth asking is whether two generic components (spheroids and discs) really can explain the diversity seen, i.e., how many fundamental building blocks and structures are required to adequately reproduce the observed galaxy population? As most of the stellar mass is contained within the bulge and disc how important are tertiary features like bars? Furthermore how many different physical origins do the various spheroids and discs have? Are elliptical galaxies simply naked bulges and are bulges related to high-redshift compact galaxies (e.g. Graham et al. 2015; Berg et al. 2014)? Are the discs of early-types, late-types and irregulars indistinguishable?

We believe that the stellar mass – half-light size (hereafter  $\mathcal{M}_{\star} - R_{\rm e}$ ) relation is a key scaling relation allowing us to address these questions for the following reasons:

• The size of a galaxy is related to its specific angular momentum making the mass and size of a galaxy fundamental observables of conserved quantities (e.g., Romanowsky & Fall 2012).

- The simple assumption that angular momentum is conserved during the initial collapse of the dark matter halo links the angular momentum and mass of a galaxy with its dark matter halo (Fall & Efstathiou 1980; Dalcanton et al. 1997; Mo et al. 1998).
- Hydrodynamical simulations now produce galaxies with realistic sizes and direct comparisons (at different epochs) are possible to study formation and evolution histories of galaxies (see for example the Evolution and Assembly of Galaxies and their Environments simulation suite, EAGLE, Schaye et al. 2015; Crain et al. 2015).
- We can empirically measure and trace the masses and sizes of galaxies and their components over a range of redshifts and in different environments (e.g., with HST as well as high-redshift ground-based surveys and soon with Euclid and WFIRST).

The  $\mathcal{M}_{\star}-R_{\rm e}$  relation therefore represents the next critical diagnostic for galaxy evolution studies beyond simple mass functions (see e.g., Bouwens et al. 2004; van der Wel et al. 2014; Holwerda et al. 2015; Shibuya et al. 2015), enabling us to trace angular momentum build-up and the emergence of the component nature of galaxies while connecting observations to simulations.

Recent studies comparing the  $\mathcal{M}_{\star}-R_{\rm e}$  relation of low and high-redshift are already yielding interesting results. For example, at high-redshift galaxies might look disc-like or elliptical/spheroidal but their physical properties are unlike any discs or ellipticals in the local Universe (see e.g. Bruce et al. 2012; Buitrago et al. 2013; Mortlock et al. 2013). Galaxies at high redshifts are typically more irregular with thick slab-like disc structures and clumpy star-forming regions (Wisnioski et al. 2012). In addition, they can be very compact but massive. In some cases, at redshift  $\sim$ 2 they are a factor of up to 6 times smaller in size than galaxies of the same mass today (Daddi et al. 2005; Trujillo et al. 2007; Buitrago et al. 2008; van Dokkum et al. 2010; Weinzirl et al. 2011).

In this paper we aim to provide a reliable low redshift benchmark of the  $\mathcal{M}_{\star} - R_{\rm e}$  relation for bulges, discs and spheroids. The bulge + disc decomposition sample is derived from a set of galaxies for which detailed morphological information is available (see Moffett et al. 2015). Section 2 describes the data and sample selection, Sections 3 and 4 describe the set up of our bulge + disc decomposition catalogue and component mass estimates. In Section 5 we present the  $\mathcal{M}_{\star} - R_{\rm e}$  relations for bulges, spheroidal and disc galaxies and discuss the association of components with their possible parent populations. We then compare our distributions to the Eagle simulation in Section 6 followed by a comparison of our low redshift  $\mathcal{M}_{\star} - R_{\rm e}$  relation with recent high redshift data from Ultra Deep Survey (UDS) region within the Cosmic Assembly Near-infrared Deep Extragalactic Legacy Survey (CANDELS, Grogin et al. 2011; Koekemoer et al. 2011) in Section 7. Finally in Section 8 we present our summary and conclusions.

Throughout this paper we use data derived from the Galaxy And Mass Assembly (GAMA) survey (Driver et al. 2011; Driver et al. 2016; Liske et al. 2015) with stellar masses

derived from Taylor et al. (2011), sizes derived from Sérsic profile fitting using SIGMA (Kelvin et al. 2012), and for a cosmology given by:  $\Lambda$  Cold Dark Matter universe with  $\Omega_{\rm m}=0.3,~\Omega_{\Lambda}=0.7,~H_0=70{\rm km s^{-1}~Mpc^{-1}}.$ 

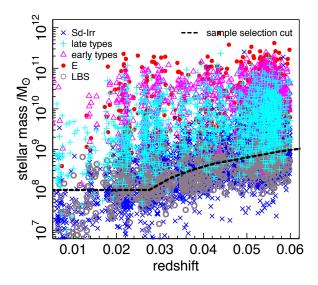
#### 2 DATA

The GAMA survey is an optical spectroscopic and multi-wavelength imaging survey of  $\sim 300000$  galaxies, combining the data of several ground and space based telescopes (Driver et al. 2011; Driver et al. 2016). It is an intermediate survey in respect to depth and survey area (Baldry et al. 2010) and thus fits in between low redshift, wide-field surveys such as SDSS (York et al. 2000) or 2dFGRS (Colless et al. 2003) and narrow deep field surveys like zCOSMOS (Lilly et al. 2007 and see Davies et al. 2014) or DEEP-2 (Davis et al. 2003).

In this paper we use data from the GAMA II (Liske et al. 2015) equatorial regions, which are centered on 9h (G09), 12h (G12) and 14.5h (G15). The three regions are  $12\times5~\mathrm{deg^2}$  in extent and have an r-band Petrosian magnitude limit of r < 19.8 mag. The spectroscopic target selection is derived from SDSS DR 7 (Abazajian et al. 2009, see Baldry et al. 2010 for details) input catalogue and we reach a spectroscopic completeness of  $\geq 98\%$  for the main survey targets. The redshifts (Baldry et al. 2014; Liske et al. 2015) are based on spectra taken with the AAOmega spectrograph at the 3.9m Anglo-Australian-Telescope (Hopkins et al. 2013) located at Siding Spring Observatory. The supporting panchromatic imaging data extends from the FUV to the far-IR via GALEX, SDSS, VISTA, WISE and Herschell (for full details see Driver et al. 2016). All optical and near-IR imaging data has matched aperture photometry (Hill et al. 2011; Liske et al. 2015) Here we focus on the redshifts (SpecCatv27), morphologies (VisMorphv03), optical imaging (ApMatchedv06) and stellar masses (StellarMass-Catv18) data products.

#### 2.1 Sample Selection

We select galaxies with 0.002<z<0.06,  $r_{petro}$  < 19.8 mag and spectra quality NQ> 2 (see Liske et al. 2015) for which visual morphologies have been established following Moffett et al. (2015). To briefly summarise the visual classification procedure: we use three-colour (Hig) postage stamps of the objects to visually inspect them. A simple classification tree is used to sort galaxies in the first instance into bulge and disc dominated, little blue spheroid (LBS) and star/artifact. In the following step the bulge and disc dominated objects are further split into single- and multi-component. Finally the multi-component galaxies are sorted into barred and unbarred. For each galaxy the result was then translated to a Hubble type: E, S0-Sa, SB0-SBa, Sab-Scd, SBab-SBcd and Sd-Irr, plus the additional LBS, Star and Artifact classifications (see Moffett et al. 2015, for further details). Note that throughout this paper the terms early- and late-type refer to our visual classification of objects being bulge or disc dominated and we do not impose any other parameter cuts (e.g. Sérsic index) to classify early- or late-type.



**Figure 1.** Shown is the redshift – stellar mass distribution for the GAMAnear sample. The points are coded according to the Hubble type established in the visual morphology classification. There appears to be no bias towards a particular Hubble type at the higher redshift boundary of our sample.

We perform bulge + disc decomposition in the r-band only as our fitting approach is computationally expensive. For the analysis we do not consider objects classified as stars or artifacts and we excluded one additional galaxy which was too large to be fit robustly. The resulting sample is hereafter called GAMAnear and comprises 7506 galaxies of which 2247 were visually classified as 2-component (S0-Sa, SB0-SBa, Sab-Scd, SBab-SBcd) and 5259 as single component (E, Sd-Irr, LBS) galaxies. Due to the low redshift range this sample extends well below  $10^9 \mathcal{M}_{\odot}$  allowing us to study the low-mass end of the  $\mathcal{M}_{\star}$  -  $R_{\rm e}$  relation. However, because of the limited volume of our survey, this also means we do not have many very high-mass galaxies to study the curvature of the  $\mathcal{M}_{\star} - R_{\rm e}$  relation at higher masses. This is important as the curvature in the (elliptical)  $\mathcal{M}_{\star} - R_{\rm e}$  relation is likely indicative of the assembly history of the galaxy population (see e.g. Bernardi et al. 2011 and references therein).

As shown by (Moffett et al. 2015, their Fig. 5) the fraction of galaxy type by mass behave as expected, e.g. the fraction of early-type galaxies increases with increasing mass and the fraction of late-type galaxies decreases. To check whether the sample selection in this paper is biased we show the sample distribution in the redshift – stellar mass plane in Fig. 1. The dashed line shows the final mass limit used in Sec. 5 to derive the  $\mathcal{M}_{\star} - R_{\rm e}$  relation. The Figure illustrates the mass segregation of the sample with the early-type galaxies being more massive than the late-types. However, there is no clear bias with morphological type or our redshift range, especially at the upper redshift boundary.

#### 3 BULGE + DISC DECOMPOSITION

Obtaining reliable bulge + disc fits is notoriously difficult, particularly in an automated fashion for large samples where the signal-to-noise ratio and resolution varies. Typically 20 to 30% of automatic fits are in some way non-physical. Previous studies have made use of a logical filter (e.g., Allen et al. 2006, see also Simard et al. 2011; Meert et al. 2015) to identify unphysical fits (e.g., component profiles which cross twice, the switching of the bulge and disc components etc., see Allen et al. 2006 for more details) and manage these failures by replacement with a single Sérsic fit. As a first step this reduces the catastrophic failure rate significantly but introduces a bias by removing the subset of two-component systems with poor fits. Following extensive exploration of our data using GALFIT3 (Peng et al. 2010), embedded in SIGMA (Kelvin et al. 2012), we identify five commonly occurring key factors which lead to poor and often catastrophic fitting outcomes. These are summarised below along with our adopted solution:

# (1) Becoming trapped in local minima and/or the limited movement of the final converged solution away from the initial conditions.

The Levenberg-Marquart (LM)  $\chi^2$  minimisation algorithm used by GALFIT3 can get stuck in a local rather than the global  $\chi^2$  minimum, especially when fitting multiple components. One way to overcome this is to vary the initial conditions (i.e., starting points) and repeat the fitting process. Convergence to a common solution, regardless of the starting point, provides confidence that the true minimum has been found. In due course a full Markov Chain Monte Carlo (MCMC) approach, that appropriately samples the prior distribution should be developed but that is beyond the scope of our current investigation at this stage.

## (2) Unphysical solutions, e.g., a scale-length of 0.1'' or a Sérsic index of 20.

In some cases the bulge or disc fits can migrate to the fit limits imposed, and these results are often not physical. While it is tempting to reduce the limits to plausible values this causes a non-physical build-up of the solutions at the limits. Moreover during the path towards convergence it can sometimes be seen that solutions migrate into extreme values and then back again. To minimise the impact of our boundaries we imposed no limits on the parameters, bar the constraint on the centre position, which is set to  $\pm$  5 pixels to account for the oversampling of the PSF (i.e. GAMA pixel size is 0.339" and SDSS FWHM=1.5"). Instead we elect to remove final solutions which settle on extreme values. We can afford to do this since we have multiple fits for each galaxy, i.e., some starting points lead to extreme outcomes but on the whole most converge to plausible values. Note that GAL-FIT does have some inbuilt constraints, such as a maximum Sérsic index of 20.

#### (3) Decision on single or multiple components.

A key problem in galaxy decomposition is to decide how many components are required. Ideally this should be derivable from the independent 1-, 2- or multi-component fits. Experimentation with the Akaike and Bayesian Information Criteria (AIC and BIC respectively) was explored but no obvious automated process for determining the number of components, which agreed with our visual assessments, was identified. This is in part due to the limited information available in single band fitting. Hence we adopt our visual classifications as priors, i.e., E, Sd-Irr and LBS galaxies are taken as single component systems and S0-Sa, SB0-SBa, Sab-Scd and SBab-SBcd as two-component. For completeness, however, we do derive and provide both 1 and 2 component fits for all systems.

#### (4) Reversal of the bulge and disc components

On occasion the initially assigned bulge component migrates to fit the disc component and the disc to the bulge. This effect was first noted in Allen et al. (2006) and can be rectified by switching the components if necessary. Here regardless of the initial parameters we assign the component with the lowest half-light radius as the bulge (i.e., inner, more compact component) and the other as the disc (i.e., outer, more extended component). This can, however, lead to cases where the bulge has a lower Sérsic index than the disc (see Appendix A for our treatment of these cases), which in the majority of cases is an unphysical solution.

## (5) Default galfit errors do not reflect the full complexity and uncertainty in the final fits.

It is known that GALFIT (like other fitting codes) often underestimates the error on the returned parameters (see e.g. Häussler et al. 2007), possibly due to the poor treatment of correlated noise in real images. Essentially the final errors do not provide any indication of fit confidence. By running GALFIT multiple times from a grid of initial conditions we can assess the level of convergence which can be used to provide more realistic error estimates. This reassessment of the errors is probably the most important outcome of our adoption of a grid of initial conditions, providing some certainty for each galaxy as to the robustness of the fit.

The five strategies above proved critical in reducing the catastrophic error rate (as assessed from visual inspection) from  $\sim 20\%$  to  $\sim 5\%$  enabling us to dispense with the need for a logical filter, and most importantly obtain realistic errors.

We recognise that many of the above could also be addressed by improving the minimisation algorithm and implementing an MCMC approach which fully samples the prior distribution. At the present time, however, in the absence of a known prior distribution and limited computing time, we believe our strategies minimise the obvious systematic issues which arise when using the GALFIT3 engine.

# 3.1 Construction of a robust decomposition catalogue

#### 3.1.1 The initial grid and convergence

As stated, for completeness, we perform both single and double (bulge + disc decomposition) component fits in the r-band on all 7506 galaxies in our sample using one or two Sérsic functions, respectively. We do not constrain any fitting parameters, except for the inbuilt limits within GALFIT. Hence in our two component fits the bulge and disc Sérsic indices are not set to any particular value (e.g., 1 and

4) as is often done in other studies. We use the Structural Investigation of Galaxies via Model Analysis (SIGMA, Kelvin et al. 2012) wrapper code for GALFIT (Peng et al. 2010). As a front-end wrapper SIGMA creates cutouts from the GAMA regions, does a local background subtraction and detects objects and stars using SEXTRACTOR (Bertin & Arnouts 1996). To obtain reliable Sérsic fits it is important that local background sky variations are accounted for, yet it is also important to not over-subtract light from the galaxy itself as this will lead to systematic errors in the galaxy flux measurements. Our local background subtraction is in addition to the background subtraction applied during mosaicing of the GAMA data. The grid size used during this additional sky estimation depends on the size of the galaxy and varies from  $32 \times 32$  to  $128 \times 128$  pixels. Using this variable mesh approach was found to be the most robust method to remove small-scale sky variations without removing light from the galaxy (for further details see Kelvin et al. 2012). After the sky subtraction, SIGMA constructs a PSF using PSFEXTRACTOR (Bertin 2013) which is later used to convolve the GALFIT models. The SEXTRACTOR outputs are also used to inform the fitting of neighbouring objects as well as provide initial starting values for the GALFIT run (for full details on SIGMA see Kelvin et al. 2012). During the actual GALFIT routine the primary and all secondary objects are simultaneously modeled using a Sérsic function. In the case of a two-component fit, GALFIT minimises the  $\chi^2$  over two Sérsic functions centred on the primary object while also fitting the secondary objects with a single Sérsic profile.

To identify convergence to the global minimum we use a grid of initial starting points (as previously discussed) for both the bulge and disc components as described below.

- Two component fitting: a total of 88 starting combinations varying input parameters as follows,
- ratio of bulge to disc size (bulge size / disc size,  $R_{SE}$ = SEXTRACTORradius):

```
1:1 (R_{SE}/R_{SE}),
```

1:2  $(0.75 \times R_{SE}/1.5 \times R_{SE})$ ,

1:4  $(0.5 \times R_{SE}/2 \times R_{SE})$ ,

1:9  $(0.33 \times R_{SE}/3 \times R_{SE})$ 

- two sets of component starting Sérsic index:

n=4+1 (bulge + disc)

n=2.5+0.7 (bulge + disc)

- component bulge and disc flux ratio:

60%:40% (bulge: disc) to 10%:90% in steps of 5%

- Single component fitting: a total of 33 combinations of the input parameters
- $-R = 2, 1, 0.5, 0.25 \times R_{SE}$
- Sérsic index n = 1, 2, 3, 4
- total magnitude  $mag = 1,0.8 \times M_{SE}$
- 1 additional model starting with  $R=R_{SE}$  and  $mag=M_{SE}$  and n=2.5

 $R_{SE}$  and  $M_{SE}$  denote the initial size and magnitude values taken from the SEXTRACTOR outputs for the entire galaxy.

Figure 2 shows an example convergence plot for a single component fit (top, G7848) and a double component

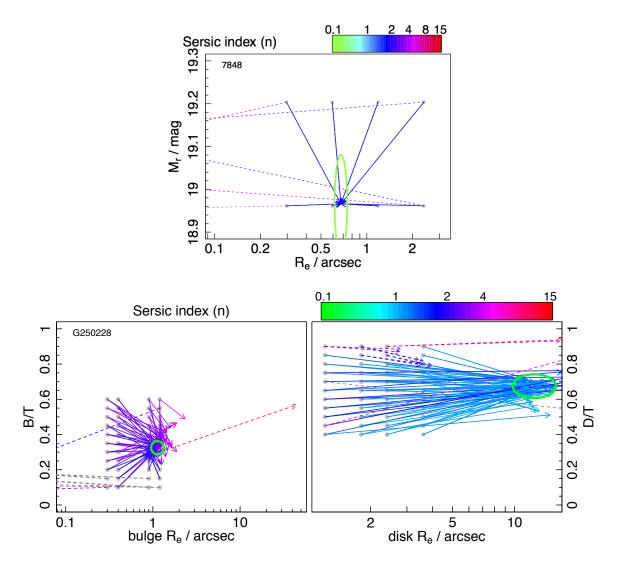
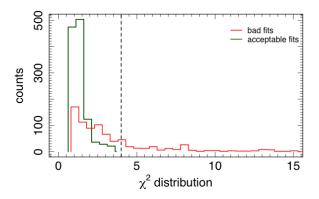


Figure 2. Convergence plot examples for a single component galaxy in the top panel and a double component galaxy in the bottom panel. The plots show the starting (grey points) and end values (arrow head) for several fit parameters making it easy to evaluate how well the galaxy was fit. For the single component fits we show the galaxy's total magnitude, size and Sérsic index. For the two-component fit we show the bulge-to-total (B/T) and disc-to-total (D/T) flux ratios of the components instead of the magnitude. The green ellipse is centred at the median output values and its size corresponds to the adopted error on the median. Note that for the single component fit the error on the magnitude corresponds to our error floor of 0.11 mag. The dashed lines indicate the fitting outcomes we consider to have failed. Note, the arrow colours correspond to the final Sérsic index values only (if it is grey then the Sérsic index is outside the range of the values we considered physical. see Sec.3.1.4) and each grey point has several arrows associated with it due to the combination of starting values of our initial grid. See the text for a detailed description.

fit (bottom, G250228). The plots show the grid of initial conditions (grey points) and vectors pointing to the final solution for each parameter combination. We plot the size versus magnitude plane for the single fits and the size versus component light fraction plane for the double fits (bulge component, left, disc component, right). The colour bar at the top shows the Sérsic indices considered and spans the same range for all convergence plots. The arrows pointing to the final output parameters are coloured according to the final Sérsic index. In practice (e.g. Fig. 2, bottom) not all fits converge to a plausible solution and hence screening is required to remove obvious bad fits. Dashed lines indicate fitting outcomes which were excluded due to bad values (see the screening descriptions below) or a large reduced  $\chi^2$ . If

the (dashed) lines are grey, the final Sérsic index was outside the range displayed in the colour bar. The green ellipse shows the median solution and its size corresponds to the adopted error on the median, i.e. the error is symmetrical and taken to be the average of the  $16^{th}$  and  $84^{th}$  percentile range. We produce convergence plots for all 1 and 2 component fits of our 7506 galaxies. As mentioned previously convergence towards a tight median value is by no means assured and a number of situations need to be managed, including: component flipping, unphysical solutions, and poor quality fits. We refer to this management as screening and define the various steps below.



**Figure 3.** Shown is the reduced  $\chi^2$  distribution of a test sample of 100 two-component galaxies for which 20 fitting outcomes each were visually inspected. The green histogram shows the fits classified as 'good' and the red histogram shows the bad fits. The dashed vertical line is the implemented reduced  $\chi^2$  cut.

#### 3.1.2 Screening via profile switching

Each of our 2247 two-component systems will have 88 model outputs from our grid of initial parameters. When fitting the galaxies component 1 has been assigned the bulge initial parameters and component 2 the disc initial values. Since GAL-FIT components can migrate significantly we ensure, after the fitting has finished, that the more compact component is taken as the bulge and the more extended as the disc. However, we find some cases where, even though the bulge is smaller in size, the disc has the higher Sérsic index. visually inspecting a number of the resulting profiles we find that the more extreme cases typically are bad fits and flag these (see Appendix A). Additionally, we relax the criterion for switching the components and allow bulges with lower Sérsic index than the disc if they are no more then 10% larger than the disc. For 1447 galaxies at least one of the 88 parameter combinations required switching the profiles output by GALFIT.

#### 3.1.3 Screening via rejection of poor quality fits

We also reject fits with poor reduced  $\chi^2$  values. To decide on an appropriate reduced  $\chi^2$  cut we randomly inspected 20 fitting outcomes for each of 100 2-component galaxies. For each fitting outcome we decided (by eye) whether it was acceptable or not based on the light-profile of the model and the resulting residuals. Figure 3 shows the distribution of reduced  $\chi^2$  for these galaxies split into 'good' and 'bad' fitting outcomes, shown as green and red histograms respectively. The vertical dashed line indicates the final cut of reduced  $\chi^2$  = 4 which is left deliberately high to ensure that even for galaxies with a lot of structure we exclude none of the acceptable fit outcomes and have enough outputs to evaluate the 'best' fit. This cut is also implemented for the single Sérsic fits. In total 20505  $(\sim 10\%)$  of the 197736 fitting results were removed from the two-component sample and 5686 ( $\sim 7\%$ ) from the 173547 fitting results of the single component sample.

#### 3.1.4 Screening via rejection of unphysical fits

Fig. 4 summarises the derived GALFIT fitted values of all combinations for our bulges and discs showing in the upper panels the bulge (left) and disc (right) sizes, and in the lower panels the bulge (left) and disc (right) Sérsic indices. Taking the top left panel we see that the bulge sizes follow two distinctive bands, one at plausible sizes (i.e., 0.1-10" scales), and one at unphysical sizes (0.001-0.01") given the data resolution of  $\sim 1.5''$ . We reject the fitting results which result in overly compact "bulges" and remove these from further considerations (red dashed line). Overly large bulges or discs are not a prominent problem but we remove obvious outliers based on the distribution of all solutions. Similarly, in Fig. 4 (bottom) we show the Sérsic index distribution. Once again the vertical red dashed lines indicate the division between the fitting results we consider physical and those we consider unphysical and that should therefore be rejected. The limits adopted leading to rejection (red dashed lines) are:

- for bulge sizes:  $R_{\rm e} < 0.01''$  or  $R_{\rm e} > 20''$
- $\bullet$  for disc sizes:  $R_{\rm e} < 0.1^{\prime\prime}$  or  $R_{\rm e} > 200^{\prime\prime}$
- $\bullet$  for bulge and disc Sérsic index: n<0.1 or n>15

These cuts are deliberately permissive and should cut out only the most unrealistic fitting outcomes. For our 2-component galaxy sample, of the 197736 combinations fitted, we reject 17343 ( $\sim 18\%$ ) based on bulge size, 616 based on disc size (< 1%), 18357 ( $\sim 19\%$ ) based on bulge Sérisc index and 11462 ( $\sim 12\%$ ) based on disc Sérsic index. For the single component fits we reject fitting outcomes based on the same limits as the disc size and Sérsic cuts of the 2-component fits. Of the 173547 combinations fit to the single component systems, we reject 2753 ( $\sim 3\%$ ) based on their size and 2355 based on their Sérsic index ( $\sim 3\%$ ). In total 49688 ( $\sim$  25%) fitting results are rejected from our 2-component sample fits and 7757 ( $\sim 4\%$ ) from our single component fits. Note, in many cases fitting results are rejected by more than one criterion (i.e. reduced  $\chi^2$  and/or size and/or Sérsic index).

We also screen our galaxies for various flags, described in Appendix A. However, we only consider two flags important during the component  $\mathcal{M}_{\star} - R_{\rm e}$  relation fits, namely the very high (or low) B/T galaxies and reversed Sérsic index galaxies. We deem the high (and low) B/T galaxies single component systems and move them from our 2-component sample to our single component sample. Galaxies with inverted Sérsic index have bulges with lower n than discs. Visually inspecting several of the profiles we find that in most cases these are bad fits, i.e., we find the disc Sérsic index n > 2. This itself would not be a problem if the errors reflect our confidence in the fit. Many of these profiles, however, have converged to this unphysical solution. We find 182 latetype 2-component systems and 87 early-type 2-component systems have inverted Sérsic index and converged profiles. We remove these galaxies from our component consideration, but use their single component profile fits to establish their global  $\mathcal{M}_{\star} - R_{\rm e}$  relation.

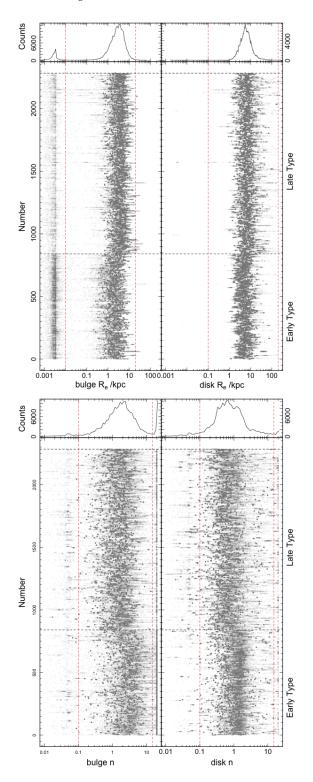


Figure 4. The top panel shows the distribution of the output size for all bulges (left) and discs (right) for the fitted two-component models (no reduced  $\chi^2{\rm cut}$  has been imposed) of each galaxy (indicated as "Number"). The dashed vertical lines show the implemented cuts on the size before the median is established. The bottom panel shows the corresponding distribution for the output Sérsic index for the bulges (left) and discs (right). In all panels the galaxies are sorted by Hubble type with the late-type 2-component systems at the top and early-type 2-component systems at the bottom. The horizontal dashed black line shows where the Hubble type changes.

#### 3.1.5 Final parameter selection

For each galaxy we consider two possible profile fit solutions taken from the remaining fitting results:

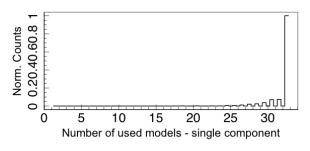
- (i) the minimum  $\chi^2$  model with the associated GALFIT parameters and errors, and
- (ii) the median fit values of the remaining fitting results and the  $16^{th}$  and  $84^{th}$  percentiles (i.e. the 1  $\sigma$  deviation of a normal distribution) as an uncertainty indicator.

While the minimum  $\chi^2$  solution should represent the best formal fit from our grid, the median model is our preferred solution, as the errors on the median reflect the level of convergence and robust errors are critical. Note that the median values are calculated for each output parameter individually and do not directly represent any single solution. In cases where the fitting converged, the median and minimum  $\chi^2$  solution will be almost identical and the  $16^{th}$  and  $84^{th}$  percentile range often is smaller than the GALFIT errors. We therefore adopt an error floor of 10% of the median value, which assures that in almost all cases the median solution is consistent with the minimum  $\chi^2$  solution within the estimated errors.

Fig. 5 shows the histogram of the number of the remaining fitting results used to calculate the median for all single component systems (top) and all 2-component systems (bottom). The single component fits often converge and the histogram peaks at  $\sim 33$  fitting results with only a small tail towards lower numbers. The 2-component fits on the other hand do not converge as often. It is encouraging that the peak is >85 fitting results, however, there is a large fraction of galaxies for which very few solutions remain for the median calculations. In addition, there is a rise towards very low numbers indicating that some galaxies are likely too complex to be fit with two components only. We find that, while the galaxies with low model counts span the whole mass range, most of them lie close to our upper redshift boundary and were classified as late-type double component systems. This shows the inherent difficulty of fitting multiple component systems in poorer image quality regimes. In addition to the tightness of the median errors, we can also use the number of fitting results left for the calculation of the median to help establish our confidence in the fitting results.

#### 3.2 Convergence Examples

We finish this section by presenting five examples which highlight some of the issues encountered and show that the median values present a robust alternative to the minimum  $\chi^2$  solutions. Figure 6 (upper panels) shows the convergence plots (as introduced in Section 3.1.1 and Fig. 2), and diagnostic plots for the median and minimum reduced  $\chi^2$  solutions (middle and lower panels, respectively). The 4 images which make up the diagnostic plots (middle and lower left panels) show, from the top left in a clockwise direction, the SDSS r-band image stamp, the model produced by GALFIT, the residual and the SEXTRACTOR segmentation map overplotted on the SDSS image stamp (the primary object is shown in purple and the secondaries in green). The red and blue ellipse show the  $R_e$  of the bulge and disc, respectively. The yellow ellipse is the original SEXTRACTOR radius and the



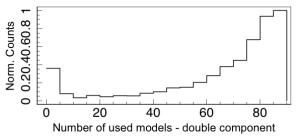


Figure 5. Here we show the distribution of the number of fitting outcomes used to calculate the median. The top panel shows the distribution for single component fits based on single component galaxies only. The bottom plot shows the distribution of the number of 2-component fits used for 2-component galaxies only. It can be seen that the single Sérsic fits generally converge nicely and the 2-component fits have a broader distribution with a spike at very low numbers.

cyan ellipses show the radii for which the surface brightness was evaluated. Also shown (middle and lower right panels) are the 1D light profile comparisons. The black points are the values extracted from within the blue ellipses, the red and green lines are the 1D light profile for the bulge and disc, evaluated from the GALFIT model and the green line is the total light profile. The lower inset panel shows the residual between the model and data. Below we discuss five examples of various fit and convergence outcomes. We only discuss examples for two-component galaxies here since we find that the single Sérsic fits generally converge well (e.g., example a):

#### • Example a: full convergence

Fig. 6a shows the ideal case of full convergence from all combinations of initial parameters. The median and minimum  $\chi^2$  solution diagnostic plots also show that both reached the same answer. The residual images show little structure and the final errors of the median fit are small as indicated by the green ellipse on the convergence plot. For 2-component fits we consider them fully converged when they have more than 80 fitting results remaining after rejection of spurious fits and the error on the median is set to the 10% error floor. This is the case for 423 ( $\sim$  19%) galaxies. Similarly, for single components over 30 fitting results must remain for the median calculation and all errors are set to the 10% error floor. This is true for 4566 ( $\sim$  87%) single component galaxies.

### • Example b: partial convergence

The median and minimum reduced  $\chi^2$  model diagnostic plots in Fig. 6b show good agreement. From the convergence plot it is obvious that many of the solutions found by GAL-

FIT were rejected during the screening process, due to an unphysical Sérsic index or high reduced  $\chi^2$ . The remaining models after screening show convergence resulting in a good solution with tight errors. For 2-component systems we consider good convergence to be reached when we have 60 to 80 solutions remaining, with the errors set to the 10% error floor. This is the case for 297 ( $\sim$  13%) galaxies. Equivalently, for single components 25 to 30 solutions must remain for the median calculation with the errors set to the error floor. We find this true for 220 ( $\sim$  4%) single component galaxies.

#### • Example c: two plausible solutions

The diagnostic plots in Fig. 6c suggest that the median model gives a physically more meaningful two-component solution than the minimum reduced  $\chi^2$  solution which is converging towards a single component solution. The convergence plot, however, highlights that the median solution is not one of the actual solutions found by GALFIT. Nevertheless, the errors on the median (green error ellipse) enclose both solutions. While the median fit can not be considered as robust this uncertainty is fairly reflected in the final errors. To establish whether several plausible solutions have been found, we test how many solutions are near the median. If less than 10% of the solutions of the median for at least one of the size, Sérsic index or B/T values of either the bulge or disc lie within 10% (i.e. the error floor) then we consider the fits to have converged to several plausible solutions which are distinct from the median. This is the case for 205 ( $\sim 5\%$ ) double component systems. For the single component systems we consider the size, Sérsic index and magnitude and find 6 (< 1%) single component galaxies have converged to several distinct solutions.

#### • Example d: no convergence

Fig. 6d shows a case where no obvious single converged solution is found, but the median model returns acceptable parameters with an appropriately broad error distribution. The diagnostic plots also show that the median model returns a physically possible solution with good residuals. The minimum  $\chi^2$  solution, however, returns a fit where the bulge, even though it has a smaller  $R_e$  is the dominant component in the outer parts of the galaxy. Since the errors associated with the median model are large this particular galaxy will not have much influence on the  $\mathcal{M}_{\star} - R_{\rm e}$  relation we fit in Sec. 5.2, but using the median parameters and large error bars means that the galaxy will not be discarded from the sample. To test non-convergence we use the same metric as in example 3, i.e. the percentage of solutions found within 10% of the median. We consider galaxies not clearly converged if more than 10% but less than 50% of the solutions lie close to the median. We test the size, Sérsic index and B/T measurements for the bulge and disc and find for 931 ( $\sim 41\%$ ) 2-component systems at least one of them is not converged. For the single component systems this is the case for 141 ( $\sim 3\%$ ) galaxies.

#### • Example e: no solution

Fig. 6e shows a case where convergence is found, however, all fits are excluded from the final catalogue due to the screening process. No median model diagnostic plot is shown due

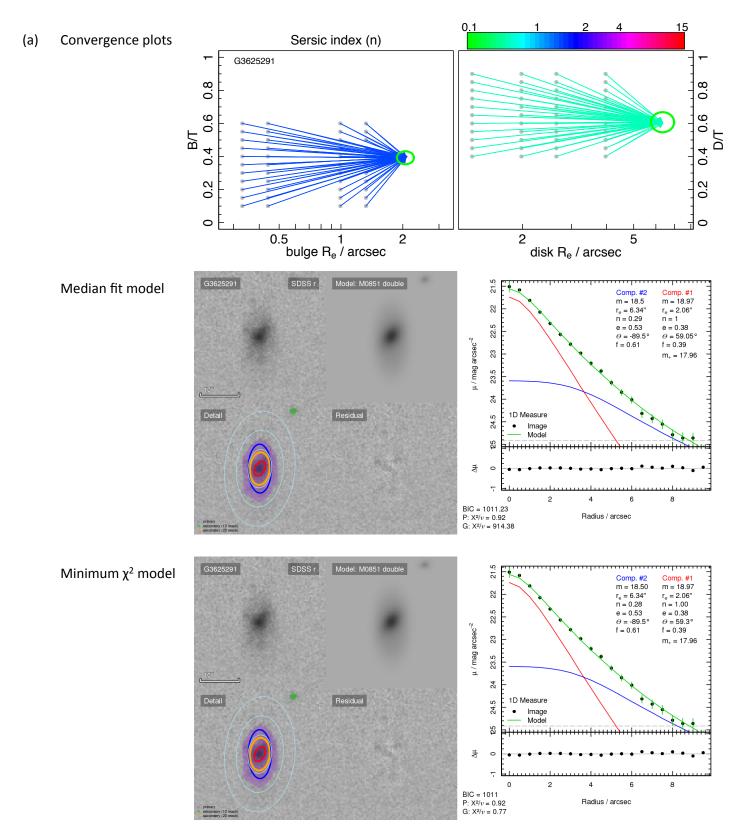


Figure 6. Presented are the convergence plots (top) and corresponding diagnostic fit plots for the median fit model (middle) and minimum reduced  $\chi^2$  solution (bottom). We show four examples ranging from full convergence to no convergence (panels a, b, c, d) and 1 example where all fits are unrealistic (panel e). A detailed discussion can be found in the text.

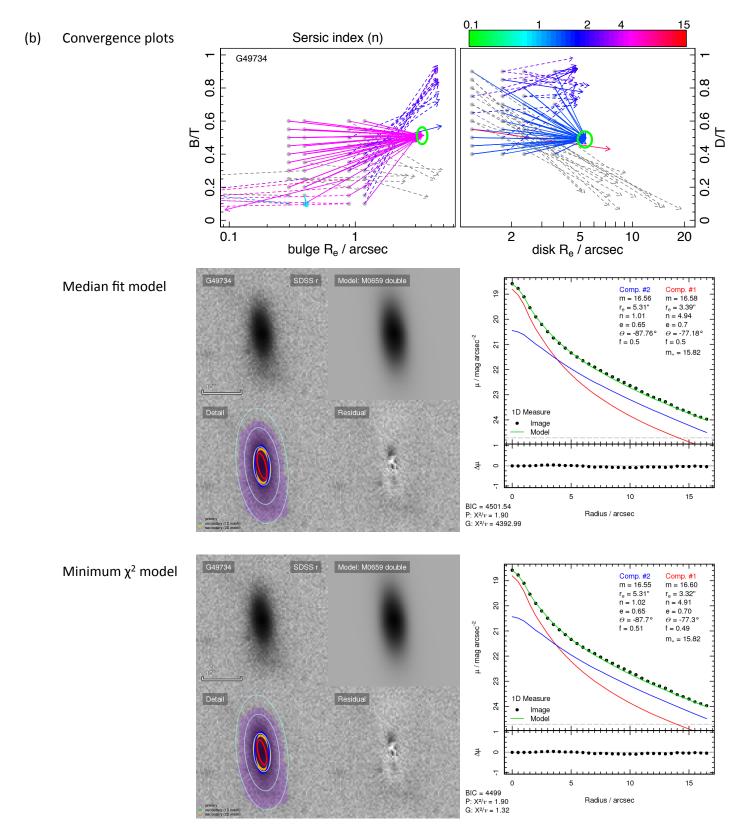


Figure 6 (Cont.).

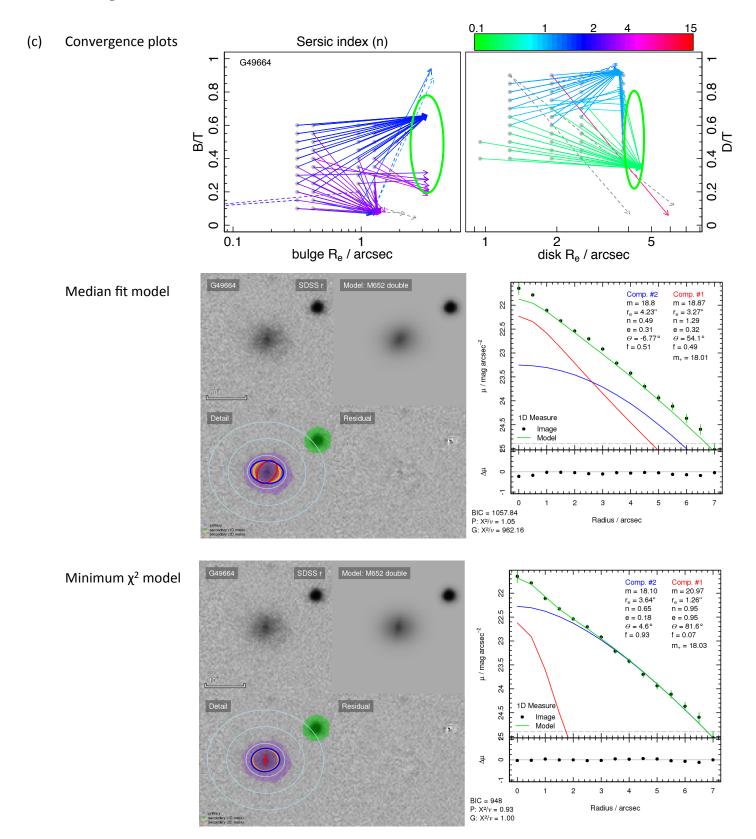


Figure 6 (Cont.).

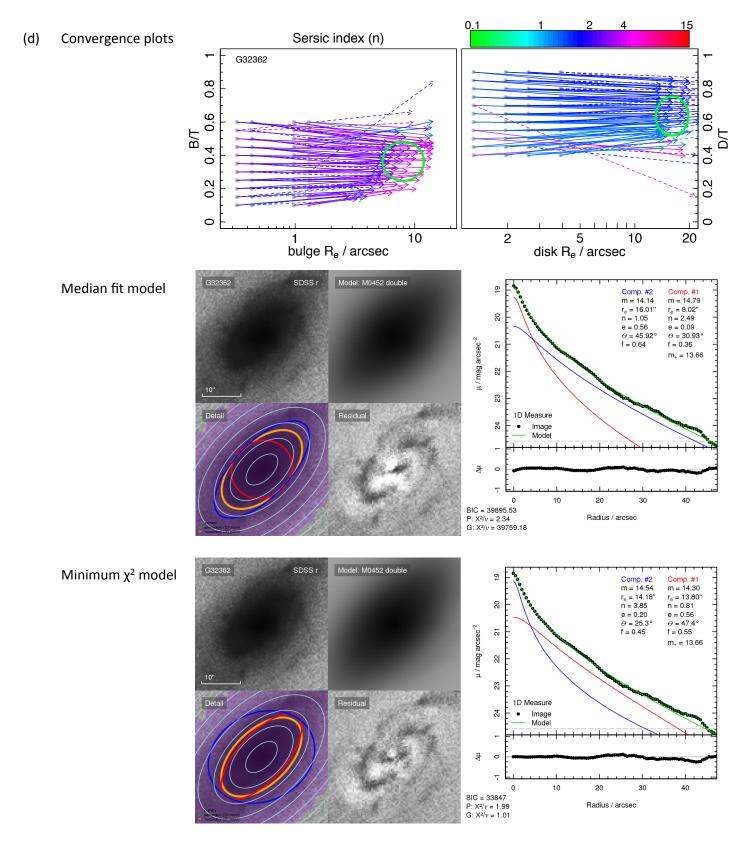
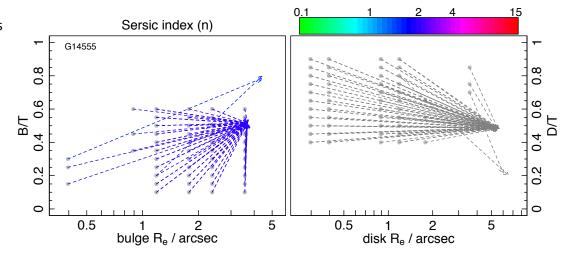


Figure 6 (Cont.).

### 14 R. Lange et al.

(e) Convergence plots



Median fit model

NA

 $Minimum \ \chi^2 \ model$ 

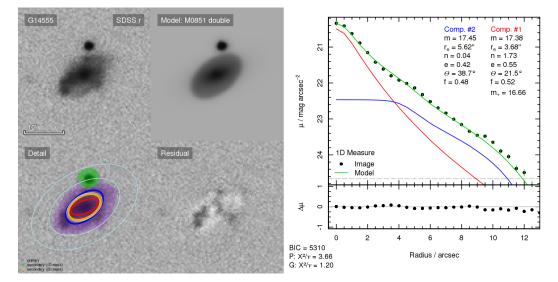


Figure 6 (Cont.).

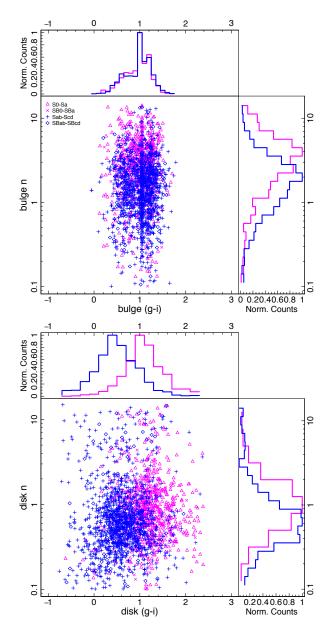


Figure 7. Shown are the (g-i) colour vs Sérsic index distributions for the bulges (top) and discs (bottom) in our two-component galaxy sample. The colour coding in both plots is the same with late-types in blue and early-types in magenta. The bulges of late-type galaxies have smaller Sérsic indices than the early-type bulges. The Sérsic index distribution of the late-type discs also peaks slightly lower than the early-type discs. Late-type discs are also bluer than early-type discs.

to all fit parameters being unrealistic. Only 120 ( $\sim 5\%$ ) of our 2-component galaxies and 129 ( $\sim 2.5\%$ ) of our single component systems fall into this category.

Convergence plots for all systems are available from the GAMA database.

#### 4 COMPONENT MASSES

To derive the  $\mathcal{M}_{\star}-R_{\rm e}$  relations for our galaxy components we now need component mass estimates. For the single Sérsic fits we can directly use the GAMA stellar mass estimates from the StellarMassesv18 catalogue and apply the fluxscale correction (for a detailed description see Taylor et al. 2011, essentially the fluxscale correction accounts for the differences between aperture matched and Sérsic photometry). These masses are based on synthetic stellar population models from the BC03 library (Bruzual & Charlot 2003) with a Chabrier (2003) initial mass function and the Calzetti et al. (2000) dust obscuration law. We find that for our sample the typical error, which has been derived in a Bayesian way (Sections 3.2-3.4 of Taylor et al. 2011), is of the order of  $\sim 0.12$  dex.

For the double component galaxies we calculate the component mass from the component colours (Driver et al. 2006) using the relationship between optical colour (g-i) and mass-to-light ratio as calibrated by Taylor et al. (2011):

$$\log M_*/M_{\odot} = -0.68 + 0.7 \ (g - i) - 0.4 \ (M_i - 4.58) \tag{1}$$

where  $M_i$  is the absolute magnitude in the *i*-band and we use the (g-i) colour of either the bulge or disc to calculate the component mass. The stellar masses derived via Eq. 1 are estimated to be accurate within a factor of two. Note that this equation is sensitive to the evolution of colour and magnitude, however, as our sample has a low redshift range the effects will be negligible.

Ideally we would use bulge and disc colours derived from the bulge + disc decompositions in the g- and i-band, however, this is beyond the scope of this paper. Hence we have to estimate the colours of the components. For this we measure the PSF and total magnitudes of our galaxies in the g, r, and i-band and we then use the GALFIT measured r-band component magnitudes and B/T to estimate the bulge and disc colours.

We measure the core (i.e. PSF) and total magnitudes, which we correct for foreground extinction, in the g, r, and i-band using LAMBDAR, a code developed to measure PSF weighted aperture photometry (for more details see Wright et al. 2016). We then equate the colours measured using the PSF magnitudes to bulge colour measurements (i.e., assuming the bulge has no colour gradient) and combine these colours with our r-band bulge magnitude from GAL-FIT to obtain bulge flux measurements in both g and i (i.e.,  $m_{i,bulge} = m_{r,bulge} - (r - i)_{PSF,bulge}$ ). In cases where the bulge colours could not be measured, we use the median bulge colour of the entire population as there is no significant trend between bulge colour and mass. We derive g and i disc fluxes by assuming that the disc flux in each band is equal to the LAMBDAR total flux minus the previously derived bulge flux.

We examine the disc (g-i) colour distribution and find a small number of extreme outliers  $(>3\sigma)$  from the colour distribution, whose disc (g-i) colours we subsequently replace with the running median disc (g-i) colour.

Finally, we use the derived bulge and disc (g-i) colours and total *i*-band magnitudes to derive stellar mass estimates for each component according to equation 1. The component colour versus Sérsic index distribution is shown in Fig. 7

#### 16 R. Lange et al.

for bulges (top) and discs (bottom). Figure 7 shows that for late-type galaxies bulges are generally redder than discs but for early-type galaxies the colours are very similar (for a detailed study of the wavelength dependence of bulge + disc decompositions in GAMA see Kennedy et al. 2016). But it also highlights the problem galaxies for which the component colour had to be set to the median colour (vertical band in the bulge plot, top) in order to be able to calculate a stellar mass estimate.

#### 5 $\mathcal{M}_{\star} - R_{\rm e}$ RELATIONS

We now present the  $\mathcal{M}_{\star}-R_{\rm e}$  relations, firstly for the different Hubble types and secondly by structural components. We conclude our analysis by presenting a combined  $\mathcal{M}_{\star}-R_{\rm e}$  relation for disc (i.e., Sd-Irr galaxies and disc components) as well as spheroids (i.e., ellipticals and classical bulge components).

Figure 8 shows the selection of the final sample, based on total stellar mass and single Sérsic profile fits, used to fit the  $\mathcal{M}_{\star} - R_{\rm e}$  relation. Note, we do not consider any galaxies below  $\mathcal{M}_* = 10^8 \mathcal{M}_{\odot}$ , since number counts are too low to establish a robust weight (after fluxscale correction this reduces the sample to 6788 galaxies). First we find and exclude outliers from the general mass-size distribution (top panel of Fig. 8). For this we fit the entire sample with a simple power law and remove all galaxies which are more than 3 sigma offset from the best fit linear relation (28 galaxies in total). We then establish the lower mass limit for a volume limited sample at z = 0.06. In the middle panel of Fig. 8 we plot the maximum redshift at which each galaxy can be seen versus its stellar mass. To establish the lower mass limit of a volume limited sample we find the point at which more than 95% of our galaxies could be seen at a redshift of 0.06 (i.e., their maximum redshift is  $z_{max} \ge 0.06$ , indicated as the dashed line). We find a lower mass limit of  $\mathcal{M}_* = 10^9 \mathcal{M}_{\odot}$  (solid blue line), which would reduce our sample size to 3679 galaxies. To include lower mass galaxies we implement a smooth volume and mass limited sample for galaxies below  $\mathcal{M}_* = 10^9 \mathcal{M}_{\odot}$ . For each galaxy we evaluate if their measured redshift is larger than their expected maximum redshift and remove them (1624 galaxies removed). The bottom panel of Fig. 8 shows the resulting sample distribution. All galaxies in red are included in our final sample (5136 total) and all grey points are excluded from the volume limited sample. For all galaxies below  $\mathcal{M}_* = 10^9 \mathcal{M}_{\odot}$ we also calculate a  $V/V_{max}$  weighting based on their redshift and our sample redshift limits of 0.002 < z < 0.06. For galaxies with  $\mathcal{M}_* > 10^9 \mathcal{M}_{\odot}$  the V/V<sub>max</sub> is set to 1. To ensure that the we only include galaxies with good, physical fits we require:

- the fluxscale correction is within 0.5 and 1.5;
- the components have Sérsic indices between 0.3 < n < 10;
- the components are resolved, i.e.  $R_e > 0.5 \times FWHM$  of the PSF which is determined from each galaxy's fit image individually;

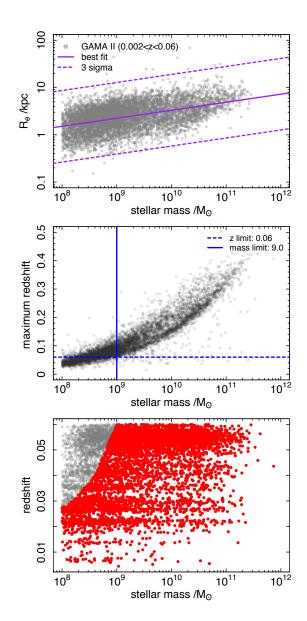


Figure 8. The top panel shows the total stellar mass – half-light size distribution (derived from single Sérsic fits) of the GAMAnear sample. All galaxies more than 3 sigma offset from the line of best fit are removed as outliers from our  $\mathcal{M}_{\star} - R_{\rm e}$  relation fits. The middle panel shows the total stellar mass - maximum redshift distribution of the sample. The blue dashed line shows our redshift limit and the solid blue line the lower mass limit for a volume limited sample. All galaxies to the right of the mass limit and above the redshift limit are included in the volume limited sample.

The bottom panel shows the total stellar mass - redshift distribution of our sample. For all galaxies below  $10^9 \mathcal{M}_{\odot}$  we have implemented a smooth volume limited sample selection. Our final sample is highlighted in red.

 at least 5 solutions remained to calculate the median fit parameters.

This reduces the sample size to 2669 single component and 1470 double component systems.

We adopt a simple power law, following Shen et al. (2003) and Lange et al. (2015), to fit the  $\mathcal{M}_{\star} - R_{\rm e}$  relation:

$$R_e = a \left(\frac{\mathcal{M}_*}{10^{10} \mathcal{M}_{\odot}}\right)^b, \tag{2}$$

where  $R_e$  is the effective half-light radius in kpc and  $\mathcal{M}_*$  is the mass of the galaxy. To perform the actual fitting we utilise the HYPERFIT package (Robotham & Obreschkow 2015) which estimates the  $\mathcal{M}_\star - R_e$  relation via Bayesian inference for each morphological group and component. During fitting we assume uniform priors and each galaxy is weighted by its  $V/V_{max}$  and the (convergence) errors for each individual galaxy are fully taken into account during the fitting process.

#### 5.1 Global $\mathcal{M}_{\star}$ – $R_{\rm e}$ relations by Hubble type

To establish the global (i.e. single component Sérsic fit)  $\mathcal{M}_{\star} - R_{\rm e}$  relation by Hubble type we have grouped the GAMAnear sample into 5 populations:

- 1564 late-type single component galaxies (including 40 high/low B/T galaxies),
- 890 late-type multi-component systems (comprised of Sab-Scd and SBab-SBcd galaxies) of which 708 have also good 2-component fits,
- 580 early-type multi-component systems (which include S0-Sa and SB0-SBa galaxies) of which 493 have also good 2-component fits,
- 806 early-type single component (including 33 high B/T galaxies), and
- 372 Little Blue Spheroids (LBS)

The resulting global  $\mathcal{M}_{\star} - R_{\rm e}$  relations are shown in panel (i) of Fig. 9, from left to right the plots are (a) Sd-Irr, (b) visually late-type multi-component systems, (c) visually early-type multi-component systems, (d) ellipticals and e) LBS. The fit parameters can be found in Table 1 (i).

We find that the single Sérsic  $\mathcal{M}_{\star} - R_{\rm e}$  relation fits to the different morphological types lie on almost parallel lines (i.e., comparable gradients but offset in normalisation). Most two-component systems are more massive than Sd-Irr galaxies, but compared at the same mass we find that two component systems are smaller than Sd-Irr galaxies. Compared to the ellipticals, however, we find that two-component systems are larger at a given mass. This corroborates the composite nature of these galaxies, i.e. the disc surrounding the bulge makes their global  $R_e$  appear larger than ellipticals but smaller than Sd-Irr galaxies at a given mass, with the offset between the Sd-Irr and elliptical relation depending on the relative dominance of the disc or bulge component. The LBS galaxies on the other hand are our smallest and least massive population. The slope of their  $\mathcal{M}_{\star} - R_{\rm e}$  relation is flatter than that of any of the other morphological types. Nevertheless, their sizes are mostly consistent with an extension of the elliptical population. In fact, within the errors the LBS relation is consistent with the low-mass elliptical relation ( $\mathcal{M}_* < 10^{10} \mathcal{M}_{\odot}$ ), see Table 1 (ii).

Here our morphological subdivisions are finer than in our previous work (L15) but broadly agree. In detail our

Hubble type (i)	a /kpc	b
Sd-Irr S(B)ab-S(B)cd S(B)0-S(B)a E LBS	$6.347 \pm 0.174$ $5.285 \pm 0.098$ $2.574 \pm 0.051$ $2.114 \pm 0.035$ $2.366 \pm 0.166$	$\begin{array}{c} 0.327 \pm 0.008 \\ 0.333 \pm 0.009 \\ 0.326 \pm 0.015 \\ 0.329 \pm 0.01 \\ 0.289 \pm 0.019 \end{array}$
(ii) E $(M_* \ge 10^{10} M_{\odot})$ E $(M_* \ge 2 \times 10^{10} M_{\odot})$ E $(M_* < 10^{10} M_{\odot})$ E $(M_* < 2 \times 10^{10} M_{\odot})$	$1.382 \pm 0.065$ $0.999 \pm 0.089$ $1.978 \pm 0.077$ $2.108 \pm 0.041$	$0.643 \pm 0.032$ $0.786 \pm 0.048$ $0.265 \pm 0.022$ $0.326 \pm 0.012$
Structural Components (iii)		
late type disc late type bulge early type disc early type bulge	$6.939 \pm 0.17$ $4.041 \pm 0.129$ $4.55 \pm 0.097$ $1.836 \pm 0.054$	$0.245 \pm 0.008$ $0.339 \pm 0.014$ $0.247 \pm 0.015$ $0.267 \pm 0.026$
Combined Case (iv)		
all discs all discs + LTB final $z = 0$ discs E + ETB final $z = 0$ spheroids	$5.56 \pm 0.075$ $5.125 \pm 0.065$ $5.141 \pm 0.063$ $2.033 \pm 0.028$ $2.063 \pm 0.029$	$\begin{array}{c} 0.274 \pm 0.004 \\ 0.263 \pm 0.004 \\ 0.274 \pm 0.004 \\ 0.318 \pm 0.009 \\ 0.263 \pm 0.005 \end{array}$
global late-types	$4.104\pm0.044$	$0.208 \pm 0.004$

**Table 1.** The regression fit parameters to Eq. 2 for the different Hubble types and structural components as well as a combined early and late type relation (see Fig.9).

late-type	a	b
r band	$3.971 \pm 1.745$	$0.204 \pm 0.018$
early-type	a	b
r band $M_* > 2 \times 10^{10} M_{\odot}$	$1.819 \pm 1.186$ $1.390 \pm 1.557$	$0.46 \pm 0.023$ $0.624 \pm 0.033$

**Table 2.** Fitting parameters taken from L15 (Table 2, 3 and B2) for the morphological late- and early-type  $\mathcal{M}_{\star}-R_{\rm e}$  relation.

Sd-Irr class has a steeper  $\mathcal{M}_{\star}-R_{\rm e}$  relation than the late-type relation in L15. This is an effect of our sample selection. In fact, fitting an  $\mathcal{M}_{\star}-R_{\rm e}$  relation to a combined sample of Sd-Irr and all two-component systems (see Table 1 (iv)) results in a relation fully consistent with the late-type relation in L15 (reproduced in Table 2). Our elliptical class has a shallower  $\mathcal{M}_{\star}-R_{\rm e}$  relation compared to the early-type relation in L15. This is also largely a sample selection effect, caused by the relative increase in the number of low-mass to high-mass ellipticals within the sample. Fitting a high-mass elliptical relation (see Table 1 (ii)), similar to L15, with  $\mathcal{M}_{\star}>10^{10}\mathcal{M}_{\odot}$  we again find good agreement with our earlier work.

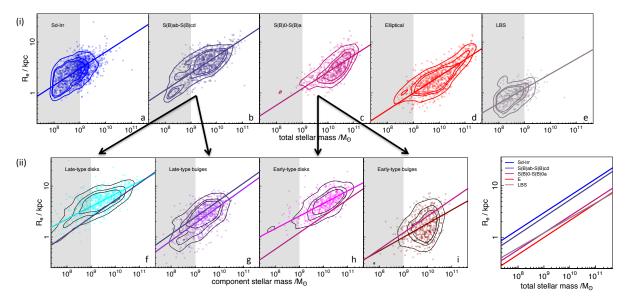


Figure 9. The top panel (i) shows the global  $\mathcal{M}_{\star} - R_e$  relation for: a) Sd-Irr (blue), b) late-type multi-component (dark purple), c) early-type multi-component (rose), d) ellipticals (red), and e) LBS (grey) galaxies.

Panel (ii) shows, from left to right, the component  $\mathcal{M}_{\star} - R_{\rm e}$  relation for late-type discs (cyan) and bulges (purple), plots f) and g) respectively. We also plot the global S(B)ab-S(B)cd relation from panel b) in comparison. The early-type discs (magenta) and bulges (dark red) are shown in plots h) and i). Again we plot the global S(B)0-S(B)a (shown in panel c) in comparison. The grey shaded areas indicate where our smooth volume limited sample selection starts and galaxies are up-weighted by their V/V<sub>max</sub>. The black lines are the 90th, 68th and 50th percentiles of the respective mass—size distributions. The arrows show from which population the components were derived.

Finally we show all global  $\mathcal{M}_{\star} - R_{\rm e}$  relations in comparison in the far right plot in panel (ii). This highlights the similarities between the different populations, i.e. the relations are parallel but offset from each other depending on their bulge fraction.

#### 5.2 $\mathcal{M}_{\star} - R_{\rm e}$ relations by galaxy component

The distributions and fits to the structural component  $\mathcal{M}_{\star} - R_{\rm e}$  relations are shown in panel (ii) of Fig. 9 and the fitting parameters can be found in Table 1 (iii).

From left to right panel (ii) shows the late-type discs and bulges (LTD and LTB, plots f and g, respectively) followed by early-type discs and bulges (ETD and ETB, plots h and i). In each panel we also show the global  $\mathcal{M}_{\star} - R_{\rm e}$  relation of the population they were derived from, which is indicated by the large black arrows. Not surprisingly we find that generally discs are larger and bulges are smaller than the global single Sérsic fits of the population they were derived from. Additionally our data shows that the component  $\mathcal{M}_{\star} - R_{\rm e}$ relations are typically less curved than the global relation they were derived from. However, the late-type discs are an exception to this and show a slight upward curvature at high masses. This is in qualitative agreement with results from Bernardi et al. (2014) who found that the discs of late-type galaxies cannot be fit with a single power law, whereas the bulges of early-type galaxies follow a pure power law which is not exhibited by their parent population.

We now wish to establish whether these component fits are consistent with the Sd-Irr or E  $\mathcal{M}_{\star}-R_{\rm e}$  relations. In particular we wish to test the following hypotheses:

- (a) ETD and Sd-Irr galaxies are associated
- (b) ETB and ellipticals are associated
- (c) LTD and Sd-Irr galaxies are associated
- (d) LTB and ellipticals are associated
- (e) LTB and Sd-Irr galaxies are associated
- (f) LBS and ellipticals are associated

- (g) LBS and Sd-Irr galaxies are associated
- (h) S(B)ab-S(B)cd systems fit with a single component (see Sec. 5.3) and Sd-Irr galaxies are associated

In Fig. 10 we visualise the affiliation of various components with either the Sd-Irr (left) or ellipticals (right) by plotting their relative deviations, defined as

$$(R_{\text{observed}} - R_{\text{predicted}})/R_{\text{predicted}}$$
 (3)

where  $R_{\text{observed}}$  represents the sizes of the tested populations, and  $R_{\text{predicted}}$  represents the predicted size using either the Sd-Irr or elliptical  $\mathcal{M}_{\star}$  –  $R_{\rm e}$  relation. In both panels the area under the curve has been normalised and each population is colour coded as shown in the legend with the dashed, vertical black line showing the peak of the deviations for the Sd-Irr galaxies and ellipticals, i.e. the populations to which we compare. The disc components, late-type bulges and the global late-type populations all broadly align with the Sd-Irr relation. The late-type discs tend to higher deviations indicating that they are larger than the Sd-Irr population, however, their peak deviation is close to 0. The good agreement of the late-type bulges with the Sd-Irr relation hints at their possible 'pseudo'bulge nature which is corroborated by their Sérsic index distribution. On the other hand dust in latet-type galaxies can artificially lower the Sérsic index of LTB, making them appear larger and thus fall within the parameter space of the Sd-Irr relation. Ellipticals, early-type bulges and LBS on the other hand do not align with the Sd-Irr relation and their distributions are completely offset. Conversely, we find that the LBS and early-type bulges visually align with the elliptical  $\mathcal{M}_{\star}$  –  $R_{\rm e}$  relation. They have similar peaks,

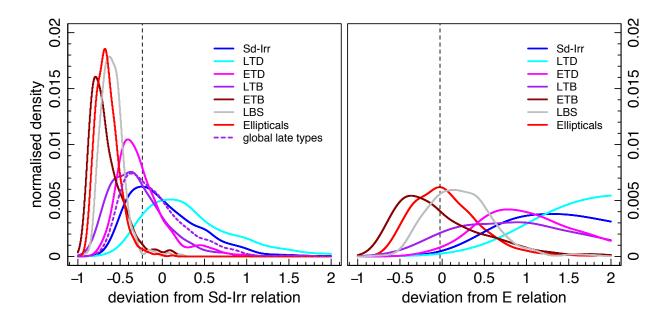


Figure 10. Shown are the relative fractional deviations of the different components from the best fit  $\mathcal{M}_{\star} - R_{\rm e}$  relation of the tested parent population. The left plot shows the deviation of the data from the Sd-Irr  $\mathcal{M}_{\star} - R_{\rm e}$  relation and the right plot shows the deviation from the elliptical relation. In both plots the area under the curve has been normalised over the range of deviations shown in the plot.

however, the early-type bulges have a broader distribution, which is in good agreement with the larger scatter observed in Fig. 9. Whether this is intrinsic or a byproduct of the decomposition is unclear. Late-type bulges, discs and Sd-Irr galaxies are offset from the elliptical distribution and do not follow their relation.

In addition to the qualitative nature of Figure 10 we also perform a two sample Kolmogorov-Smirnoff-test (KStest) for each hypothesis. For this we compare the different samples to the  $\mathcal{M}_{\star} - R_{\rm e}$  relation of either Sd-Irr or elliptical galaxies. Since the KS-test in essence compares the cumulative distributions of two populations in one dimension it does not take into account the spread of our data around the  $\mathcal{M}_{\star}-R_{\mathrm{e}}$  relation. Hence we decided to bin our data in  $\Delta \log(M_{\star}) = 0.2$  steps to establish the median mass and size of the bin and we use the median bin mass to calculate the expected size based on either the Sd-Irr or elliptical  $\mathcal{M}_{\star} - R_{\rm e}$ relation. This also allows us to test whether the expected size distribution from the  $\mathcal{M}_{\star} - R_{\rm e}$  relation fit agrees with the observed median size distribution of the sample, even for cases which have only little or no overlap in the mass-size plane (e.g. LBS galaxies and the elliptical  $\mathcal{M}_{\star} - R_{\rm e}$  relation). The resulting KS statistics are shown in Table 3. For our tested assumptions, combining the results of Fig. 10 and the KS-test, we conclude:

- a,b,c,e,f,h = True
- $\bullet \ d,g = False$

# 5.3 Should S(B)ab-S(B)cd systems be described as single or multi-component?

In the previous section we found that both late-type discs and late-type bulges are associated with the Sd-Irr relation.

Case	D	p-values
common anto sea Cd Inn and E se	lation	
components vs Sd-Irr and E re	eration	
a) ETD vs Sd-Irr	0.4	0.418
b) ETB vs E	0.38	0.66
c) LTD vs Sd-Irr	0.29	0.635
d) LTB vs E	0.5	0.1
e) LTB vs Sd-Irr	0.42	0.256
f) LBS vs E	0.43	0.575
g) LBS vs Sd-Irr	1	0.001
global fits vs Sd-Irr and E rela	tion	
Sd-Irr vs Sd-Irr	0.33	0.73
global late-types vs Sd-Irr	0.23	0.898
global early-types vs Sd-Irr	0.57	0.019
E vs E	0.2	0.994
global late-types vs E	0.62	0.013
global early-types vs E	0.29	0.635

**Table 3.** Shown are the D- and p-values for a two tailed KS-test on the hypothesis stated in Sec. 5.2.

Additionally, the Sérsic index distribution of these components (Fig. 7, bulge, top and disc, bottom) show that late-type discs extend to lower Sérsic indices than typically expected. In fact, the LTD in our sample have a median Sérsic index of  $n \sim 0.6$  (i.e., a more Gaussian-like light profile). The Sérsic index distribution of our late-type bulges peaks at  $n \sim 2$  which is much lower than would be expected for an intermediate to high-mass classical bulge where  $n \sim 4$ . There are two possible reasons, (i) we are seeing the effects of dust

affecting both bulge and disc, and/or (ii) late-type systems are composed of pseudo-bulges and discs.

Seeing a bulge or disc through dust has the effect of lowering the Sérsic index as well as making them appear larger (Pastrav et al. 2013a,b). On the other hand, if the late-type galaxies do indeed contain pseudo-bulges, which are arguably perturbations of the disc, this begs the question as to whether the late-type galaxies should or should not be decomposed into 2 components.

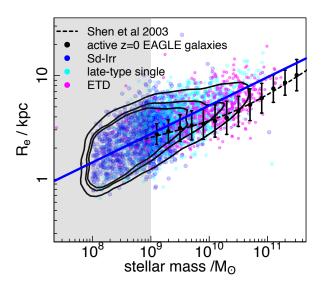
Figure 10 (left panel) and a KS-test show that the global (i.e. single Sérsic fit) late-type  $\mathcal{M}_{\star} - R_{\rm e}$  relation can also be associated with the Sd-Irr relation and a decomposition of the late-type 2-component systems is not strictly necessary. Another issue to consider here is that if S(B)ab-S(B)cd galaxies are truly 2-component systems than fitting their light profile with a single component only would bias our size estimation to larger sizes (Bernardi et al. 2014), especially for brighter and larger galaxies. However, the majority of our S(B)ab-S(B)cd systems are comparatively small when considering this effect found by Bernardi et al. (2014). Furthermore fitting a final disc relation using either global or component fits for the S(B)ab-S(B)cd galaxies we find that the resulting  $\mathcal{M}_{\star} - R_{\rm e}$  relations are nearly identical (these relations are also given in Table 1 for reference). Hence, as considering either single or two-component sizes has little effect on the  $\mathcal{M}_{\star} - R_{\rm e}$  relation and as we cannot conclusively tell the difference between a pseudo-bulge and a classical bulge without kinematic data, we opt to use the global (single component) S(B)ab-S(B)cd galaxies for our final z = 0disc relation to avoid over-interpreting our results. This is also in concordance with other recent studies where late-type two component galaxies are considered 'bulgeless' discs (for example Sachdeva et al. 2015).

#### 5.4 Combined disc and spheroid $\mathcal{M}_{\star} - R_{\rm e}$ relations

In this section we now aim to establish the definitive z=0 disc and spheroid  $\mathcal{M}_{\star}-R_{\rm e}$  relation composed of associated global and component populations as identified in the previous section. In summary, we consider the following populations to be associated:

- i) Sd-Irr and late-type galaxies and the discs of early-type galaxies and,
- ii) Ellipticals, early-type bulges and LBS.

We show our final combined  $\mathcal{M}_{\star} - R_{\rm e}$  relations for discs (top) and spheroids (bottom) in Figure 11. The data points are colour coded by the population they belong to. The solid lines show our final disc and spheroid  $\mathcal{M}_{\star} - R_{\rm e}$  relation fits. The dashed red line in the lower panel shows our high mass elliptical relation, which would be more appropriate to use for comparisons with high redshift data (see Sec. 7). In comparison we also show simulated z = 0 eagle galaxies which are actively star-forming (top) and passive (bottom) as black points with error bars. We discuss the selection and comparison of the EAGLE data in the next Section. Additionally, we show the  $\mathcal{M}_{\star} - R_{\rm e}$  relation for late- and early-type galaxies by Shen et al. (2003) as a dashed black line. We show these relations as they were used to calibrate the simulated data from EAGLE. Note, we plot the Shen et al. (2003) relations only over the mass range for which they were established and we have corrected them for the size-wavelength dependence



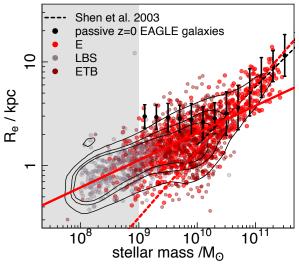


Figure 11. Shown are the final z=0 disc (top) and spheroid (bottom)  $\mathcal{M}_{+} - R_{e}$  relations. The disc relation consists of Sd-Irr galaxies, early-type disc components and reconstituted late-type 2-component systems. The spheroid relation is composed of ellipticals, early-type bulges and LBS. The solid coloured lines are the final disc and spheroid  $\mathcal{M}_{\star}-R_{\mathrm{e}}$  relation fits to the data, the dashed red line shows the high mass elliptical only () relation, the grey shaded areas indicate where our smooth volume limited sample selection starts and galaxies are up-weighted by their  $V/V_{max}$ . The black (solid) lines are the 90th, 68th and 50th percentiles of the respective mass-size distributions. Additionally we show the active (top) and passive (bottom) z = 0 galaxies from the EAGLE simulation (black points + error bars) and the  $\mathcal{M}_{\star} - R_{\rm e}$  relation from Shen et al. (2003) corrected for waveband and circularised radius (dashed black line, only shown over the mass range in which it was established).

using the equations given in L15 (their Table 4), assuming  $\mathcal{M}_{\star}=10^{10}\mathcal{M}_{\odot}$ . Additionally, we also correct for the fact that Shen et al. (2003) use a circularised radius<sup>1</sup>. The  $\mathcal{M}_{\star}-R_{\rm e}$  relation parameters can be found in Table 1 (iv).

<sup>&</sup>lt;sup>1</sup>  $R_c = \sqrt{b/a}R_e$ , where b and a are the semi-minor and semi-

For the combined disc populations (Fig. 11, top) we find that the relation flattens considerably from the Sd-Irr only relation, an effect of including the high-mass early-type discs and late-type galaxies. If our assumption is true that the Sd-Irr galaxies, early-type discs, and 'reconstituted' latetype galaxies are related than this could hint at a possible change in the slope of the late-type  $\mathcal{M}_{\star} - R_{\rm e}$  at high masses. Comparing our  $\mathcal{M}_{\star} - R_{\rm e}$  relation to the Shen et al. (2003) relation it is obvious that our data follows an opposite trend at higher masses and does not turn up but down. This effect arises because we compare a component relation to a global relation, we see a similar steepening of our data if only global sizes are considered in the  $\mathcal{M}_{\star} - R_{\rm e}$  distribution instead. In addition, since our data extends to lower masses than the Shen et al. (2003) analysis we can also see that the  $\mathcal{M}_{\star}$  –  $R_{\rm e}$  relation does not actually flatten out at lower masses. Overall this shows that a linear fit (in log-log space) to the data is sufficient to describe the disc  $\mathcal{M}_{\star} - R_{\rm e}$  relation over the observed stellar mass range.

Compared to the elliptical only relation, the slope of the combined elliptical, LBS and ETB  $\mathcal{M}_{\star} - R_{\rm e}$  relation also flattens and essentially lies parallel to the final disc relation. We also see a turn off in the data, albeit in the opposite sense to the combined disc relation, i.e. the data flattens at the low-mass end. For the spheroid  $\mathcal{M}_{\star} - R_{\rm e}$  relation this flattening is caused by the inclusion of more low-mass components (mainly LBS), this is corroborated by a comparison of our data to the Shen et al. (2003) relation. We see a very good agreement over the mass range in which the Shen et al. (2003)  $\mathcal{M}_{\star} - R_{\rm e}$  relation was established. However, a clear turn-off, or flattening can be seen at low masses. That the flattening is caused by low-mass galaxies is further supported by comparing the final spheroid relation to the low-mass elliptical only relation, see Table 1 (ii), which are within the errors identical. Additionally we fit the  $\mathcal{M}_{\star} - R_{\rm e}$ relation to ellipticals and early-type bulges only and find that this relation is, within the errors, identical to the elliptical only relation, supporting the notion that early-type bulges are indeed classical bulges and "elliptical-like". This begs the question whether the LBS are all indeed early-type galaxies. As shown in Fig. 10 the distribution of the relative deviations of LBS from the elliptical  $\mathcal{M}_{\star} - R_{\rm e}$  relation is consistent with them being associated. Their broad distribution, however, which looks similar to a top-hat function, extends to high deviations which suggests that not all LBS galaxies are the same. Without higher resolution imaging data we cannot yet conclusively say if LBS are indeed all early-type galaxies. Preliminary visual inspection of LBSs available in the higher resolution VIKING imaging shows that a significant fraction of them are actually two-component systems.

If the LBS and ETBs, as well as the high- and low-mass ellipticals are all indeed the same population, then a curved fit, i.e. a double power law, is necessary to describe the  $\mathcal{M}_{\star}-R_{\rm e}$  relation. In lieu of a definitive answer we recommend considering the high- and low-mass early-type populations separately. In Table 1 (ii) we provide a  $\mathcal{M}_{\star}-R_{\rm e}$  relation for high- and low mass ellipticals considering two different mass separators at  $M_{\star}=10^{10}M_{\odot}$  and  $M_{\star}=2\times10^{10}M_{\odot}$ .

major axes. Using the average ellipticity of our sample we get  $R_e = 1.33 \times R_c$ .

The combined  $\mathcal{M}_{\star}-R_{\rm e}$  relation to all ellipticals, ETBs and LBS in Table 1 (iv) should only be considered for samples that primarily contain galaxies with  $M_{\star} \leq 10^{10} M_{\odot}$ . Furthermore, it should be noted that previous studies have reported a second deviation of the early-type  $\mathcal{M}_{\star}-R_{\rm e}$  relation from a simple power law at very high masses ( $\mathcal{M}_{\star} \sim 2 \times 10^{11} \mathcal{M}_{\odot}$ , e.g., Bernardi et al. 2007; Hyde & Bernardi 2009; Bernardi et al. 2011, 2014). Due to the limited volume we survey, we do not see this deviation. However, as this second change in slope is likely linked to the formation history of galaxies it needs to be taken into account when studying the high-mass end of the  $\mathcal{M}_{\star}-R_{\rm e}$  relation.

Finally we include the caveat that the flattening of both the disc and spheroid population could be real or due to miss-classification in our visual morphology sample or a significant bias in our bulge + disc decomposition. Ultimately deeper data, such as that provided by VST KiDS and a fit in all available wavelengths to establish robust masses should clarify whether the flattening is real or an artifact of our methodology.

## 6 DOES THE SIZE DISTRIBUTION POSE A PROBLEM FOR SIMULATIONS?

In this section we have a first look at comparing our final disc and spheroid relations with data from the EAGLE simulation (Schaye et al. 2015; Crain et al. 2015). EAGLE is a suite of cosmological hydrodynamical simulations performed at two numerical resolutions, in periodic volumes with a range of sizes, and using a variety of subgrid implementations to model physical processes below the resolution limit. The subgrid parameters governing energetic feedback mechanisms of the EAGLE reference model were calibrated to the z=0 galaxy stellar mass function, galaxy stellar mass black hole mass relation, and galaxy stellar mass - size relations (see Crain et al. 2015 for details and motivation). The EAGLE reference model reproduces many observed galaxy relations that were not part of the calibration set, such as the evolution of the galaxy stellar mass function (Furlong et al. 2015b), of galaxy sizes (Furlong et al. 2015a), of their optical colours (Trayford et al. 2015), and of their atomic (Bahé et al. 2016) and molecular gas content (Lagos et al. 2015), among others, and thus is an excellent testbed to compare with our observations.

We use the public database of Eagle described in McAlpine et al. (2015). In particular, we focus our attention on the reference model of EAGLE run in a cubic volume of length 100 comoving Mpc on a side with  $2 \times 1504^3$  dark matter and gas particles (particle masses are  $9.7 \times 10^6 \mathcal{M}_{\odot}$ and  $1.81 \times 10^6 \mathcal{M}_{\odot}$ , respectively, which help to reach a physical resolution of 0.7 kpc. One of the notable aspects of EAGLE is the plethora of sub-grid baryonic physics included in the model: (i) radiative cooling and photo-heating rates, (ii) star formation, (iii) stellar evolution and metal enrichment, (iv) stellar feedback, and (v) black hole growth and AGN feedback. These physical models are the key ingredient to reproduce a large set of properties of the observed galaxy population in the local Universe. For more details of the simulation we refer the reader to Schaye et al. (2015).

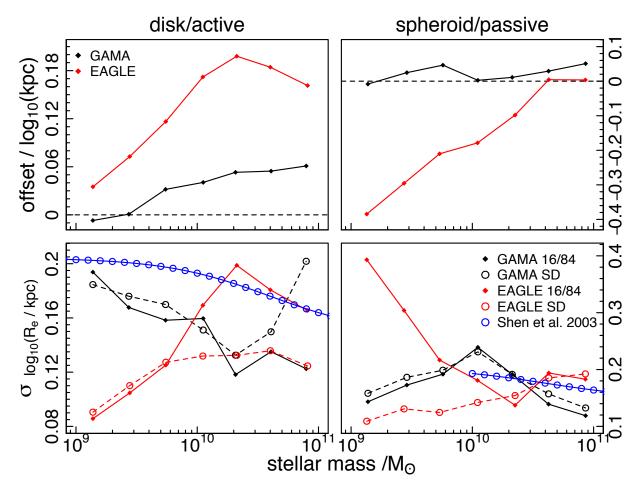


Figure 12. Shown is the offset (top) of the Gaussian distribution fit of the observational and simulated data from the disc (left) and spheroid (right)  $\mathcal{M}_{\star} - R_{\rm e}$  relation. The bottom panel shows the corresponding scatter  $\sigma_{\log(R_e)}$  (SD, circle plus dashed lines) as well as the  $16^{th}/84^{th}$  percentile scatter (16/84, filled symbols plus solid line) of the distribution. Additionally we show the  $\sigma$ -mass relation from Shen et al. 2003 for comparison.

To compare EAGLE data with our observations we use the masses and sizes of z=0 active and passive galaxies, calculated from data in the public database, following the method described in Furlong et al. (2015b). Briefly, the size is the mean half-mass radius taken within a 100kpc aperture projected along the x-y, x-z, and y-z axes. Note that this definition of size differs from the size used to calibrate the simulation against the Shen et al. (2003) late type galaxies. The galaxy sizes used in the calibration are based on the Sérsic scale length obtained from fitting Sérsic profiles to the surface density profiles of the simulated galaxies. To separate active and passive galaxies Furlong et al. (2015b,a) use a specific star-formation rate cut of  $0.01~Gyr^{-1}$ , which is approximately one decade below the observed main sequence of star formation. For the following comparison there are two notable caveats:

- (i) we compare our measurements of the disc and spheroid  $\mathcal{M}_{\star}-R_{\rm e}$  relation for the half-light radius to the half-mass radius of the simulated galaxies, and
- (ii) we separate active and passive simulated galaxies to compare to our discs and bulges, respectively. We caution the reader that although a correlation between being bulge (disc)-dominated and being passive (active) is expected,

they are not necessarily the same populations.

Figure 11 shows our final disc and spheroidal  $\mathcal{M}_{\star} - R_{\rm e}$ relations. We plot the EAGLE data as black points and the error bars indicate the 16<sup>th</sup> and 84<sup>th</sup> percentile of their distribution. We also show the Shen et al. (2003) relation for n < 2.5 and n > 2.5 galaxies as the dashed line. The difference in the shape of the disc  $\mathcal{M}_{\star} - R_{\rm e}$  relation (Fig. 11, top) is in part caused by comparing disc only components (our data) with global sizes (EAGLE). This is because also considering a bulge component in a global profile fit has the effect of lowering the half-mass radius due to their typically smaller size and higher concentration compared to discs. For the spheroid relation (Fig. 11, bottom) we agree well with the EAGLE data down to  $\mathcal{M}_* \sim 10^{10} \mathcal{M}_{\odot}$ . Below this mass limit we see a much less marked change of slope compared to the EAGLE data. This could be due to comparing passive galaxies in EAGLE with our spheroid sample that contains LBS galaxies. The latter are arguably active systems. On the other hand, at least in part, this could be explained by the known limitations of the simulation, e.g. the high fraction of passive low mass galaxies at z = 0 due to the finite sampling of the star-formation in low mass galaxies. A one-to-one comparison, in which EAGLE galaxies are analysed with the same pipeline applied to GAMA galaxies is needed to shed light on this issue. A further point to consider is the shape of the  $\mathcal{M}_{\star}-R_{\rm e}$  relation at very high masses  $(\mathcal{M}_{*}\geqslant 10^{11.3}\mathcal{M}_{\odot})$  which has been shown to deviate from a pure power law (Bernardi et al. 2007; Hyde & Bernardi 2009; Bernardi et al. 2011, 2014, e.g.,). As mentioned previously, due to our small survey volume, we do not sample the high-mass end of the  $\mathcal{M}_{\star}-R_{\rm e}$  relation well and cannot confirm the curvature of the relation. However, to ensure that simulations return realistic galaxy sizes at all masses, not only the mass range over which they were matched, a detailed comparison with the relevant studies is necessary.

Figure 12 summarises the findings of Figure 11. Here we compare the distribution of the vertical scatter for the GAMA and EAGLE data from our disc (left) and spheroid (right)  $\mathcal{M}_{\star} - R_{\rm e}$  relation. We divide our sample in bins of 0.3 dex in stellar mass and fit a Gaussian to the distributions (see Figures B1 and B2) to study the scatter of the data from the  $\mathcal{M}_{\star}-R_{\rm e}$  relation at fixed stellar mass. The top panel of Fig. 12 shows the offset ( $\mu$  of the Gaussian fit) of the  $\log(R_e)$  scatter from the  $\mathcal{M}_{\star} - R_e$  relation. Note that for the high mass spheroid/passive galaxy sample we actually compare to the high mass  $(\mathcal{M}_* > 10^{10} \mathcal{M}_{\odot})$  elliptical  $\mathcal{M}_{\star} - R_{\rm e}$  relation. This does not change the scatter of the data but has the effect of moving the offset from the  $\mathcal{M}_{\star} - R_{\rm e}$  relation to  $\sim 0$ . The bottom panel shows the  $\sigma_{\log(R_e)}$  versus stellar mass distribution. We show the sigma (standard deviation, SD) of the best fit Gaussian to the underlying distribution as well as the 16<sup>th</sup> and 84<sup>th</sup> percentile scatter of the GAMA and EAGLE data. In cases where the underlying distributions are somewhat skewed the SD and 16/84<sup>th</sup> percentile sigma do not agree well and could even hint at a possible bimodality in the underlying distribution. We also include the scatter versus mass relation from Shen et al. (2003), corrected from  $\log_e$  to  $\log_{10}$ , for comparison. Note that the Shen et al. (2003) scatter-mass relation is based on the combined scatter of their early- and late-type  $\mathcal{M}_{\star} - R_{\rm e}$  relations. In Appendix B we show the histograms of the scatter of the data from the  $\mathcal{M}_{\star} - R_{\rm e}$  relation for each mass bin with the best fit distribution overplotted.

Examining Fig. 12 the offset in the modes of the GAMA and EAGLE data is expected, as seen in Figure 11. This is due to the EAGLE simulation being calibrated using the Shen et al. (2003) relations. To some extent the variance (or sigma values as indicated on the Figure) is of more interest as these have not been explicitly tailored in the simulation to match the data distributions.

We focus our analysis on the sigma derived from the Gaussian fits, i.e., the points labelled SD in Fig. 12. We find that for the disc/active population the variance in the simulated data is almost always smaller than the observations. For the spheroid/passive populations there is a divide at  $\mathcal{M}_* \sim 2 \times 10^{10} \mathcal{M}_{\odot}$  with the simulated data having a smaller variance for less massive systems compared to the data and a comparable variance for more massive systems. This is somewhat surprising as the dark matter spin distributions are known to be quite broad. If coupling is strong one would expect the distribution of specific angular momentum and disc sizes to be comparably broad. Because we are comparing light against mass and components against classes we should be careful with interpreting any deviations. Clearly

an improved comparison can be made from bulge + disc decompositions of the EAGLE images which would place the observational and simulated data onto the same footing. Although gri images of EAGLE galaxies with  $\mathcal{M}_* > 10^{10} \mathcal{M}_{\odot}$  have been made publicly available from their database, it is still not sufficient to perform an analysis like the one done here for GAMA. Images of individual bands, preferably to lower stellar masses, would be required for this. For the moment we consider Figures 11 and 12 to provide a good demonstration of the potential of the mass-size plane for comparing observational and simulated data.

# 7 COMPARISON WITH HIGH REDSHIFT DATA

There is a well know discrepancy between the  $\mathcal{M}_{\star}$  –  $R_{\rm e}$ relation of high redshift galaxies (z>1) and the local  $\mathcal{M}_{\star} - R_{\rm e}$  relation, with the high redshift galaxies at the same mass being smaller than their low redshift counterparts (e.g. Daddi et al. 2005; Longhetti et al. 2007; van Dokkum et al. 2008; Trujillo et al. 2006, 2007). This itself is not a problem since an evolution in the  $\mathcal{M}_{\star} - R_{\rm e}$ relation is expected due to galaxies at later times being formed through less dissipative events (i.e. low redshift progenitors are gas-poorer than high redshift progenitors, see e.g. Hopkins et al. 2009). Hence at lower redshift newly formed galaxies are expected to be bigger than their high redshift counterparts. However, the lack of old massive compact systems at low redshift means that the galaxies observed at high redshift must have grown by a factor of up to  $\sim 5$ -6 to end up on the local  $\mathcal{M}_{\star} - R_{\rm e}$  relation. One might argue that the measurements of high redshift galaxies are inherently difficult and the observed size growth is biased by systematics. However, several studies have shown that, even considering all the uncertainties in the mass and size measurements of the high-z galaxies, the size growth is real. Even for the most unfavorable cases the high redshift galaxies lie well below the present day  $\mathcal{M}_{\star} - R_{\rm e}$  relation (see e.g. Buitrago et al. 2013; Weinzirl et al. 2011).

In this section we briefly compare published measurements of high-redshift galaxies from the CANDELS-UDS field (Mortlock et al. 2013, 2015; Margalef-Bentabol et al. 2016), against our local disc and spheroid  $\mathcal{M}_{\star} - R_{\rm e}$  relations. To briefly summarise, the CANDLES-UDS data contains 1132 galaxies with  $\mathcal{M}_{\star} > 10^{10} \mathcal{M}_{\odot}$  and 1 < z < 3. Of these 683 are fit with a single component Sérsic profile while 449 with a bulge + disc profile. For the two-component fits the disc is set to n = 1 and the bulge n is free unless the fit failed in which case the bulge was reset to either n = 1 or n = 4 (Margalef-Bentabol et al. 2016).

The CANDELS-UDS data were obtained using the Hubble Space Telescope's Wide Field Camera 3 (WFC3/IR) F160W (H) band, which is well matched in physical resolution and rest-wavelength to the GAMA low-redshift SDSS data. Here we restrict ourselves to the redshift range 1 < z < 1.5 where the H-band equates to a comparable rest wavelength of 640-800nm. For the HST data stellar masses were derived by the CANDELS team using BC03 stellar populations, a Chabrier IMF and with the same  $\Lambda$ CDM cosmology as our analysis. One

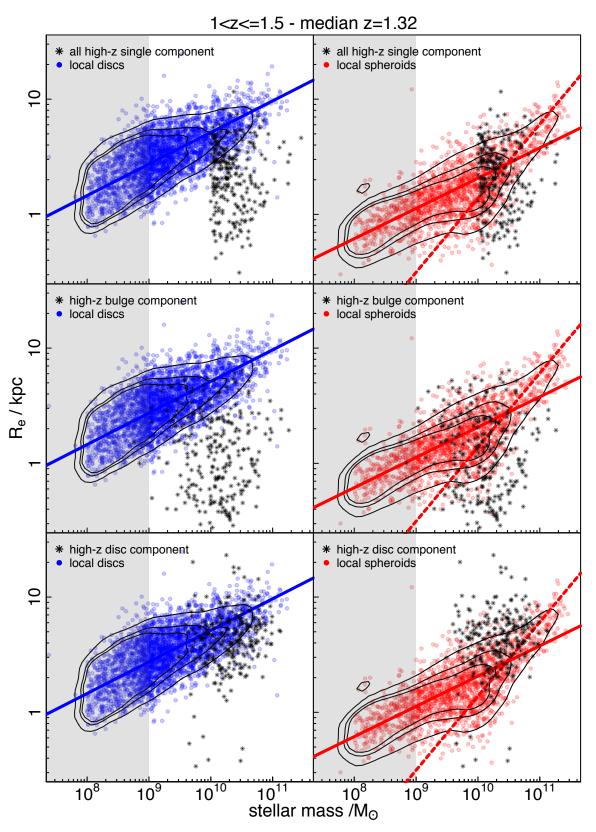


Figure 13. We show the local disc (blue, left side) and spheroid (red, right side)  $\mathcal{M}_{\star} - R_{\rm e}$  distributions in comparison to the high redshift CANDELS data (1 < z < 1.5) for our three high-redshift populations. The solid blue and red lines are our final disc and spheroid  $\mathcal{M}_{\star} - R_{\rm e}$  relation as presented in Sec. 5.4. The dashed red line is the local high mass ( $\mathcal{M}_{*} > 10^{10} \mathcal{M}_{\odot}$ ) elliptical  $\mathcal{M}_{\star} - R_{\rm e}$  relation as given in Table 1 (ii). It is immediately obvious that most data do not agree with the disc  $\mathcal{M}_{\star} - R_{\rm e}$  relation. On the other hand there is good agreement between our local spheroids and the high-redshift single component systems and the high-redshift bulge components (with a compact extension of the high-redshift bulge components).

minor difference worth highlighting between the HST and GAMA analysis, is the derivation of component masses. For the HST data this was based on the single band bulge-to-total ratio rather than component colours (see Section 4).

The high-redshift CANDELS-UDS sample naturally divides into three distinct populations: single component systems, bulges of two-component systems, and discs of two-component systems. We now explore whether any of these three populations follow a similar  $\mathcal{M}_{\star}-R_{\rm e}$  relation to the local benchmarks.

In Fig. 13 we show the  $\mathcal{M}_{\star} - R_{\rm e}$  plane with the various high-redshift samples overlaid on the local benchmark data (left panel against discs and right panel against spheroids). Note that we show both the local spheroid  $\mathcal{M}_{\star} - R_{\rm e}$  (solid red line) and the high mass elliptical  $\mathcal{M}_{\star}$  –  $R_{\rm e}$  relation  $(\mathcal{M}_* > 10^{10} \mathcal{M}_{\odot}, \text{ dashed red line})$ . For our comparison we concentrate on the high mass elliptical relation since this describes our low redshift data better in the mass range observed for high redshift galaxies. It is immediately apparent that none of the populations provide a good association with the low-redshift disc benchmark (left panel), with all systems being significantly more compact. However, the high-redshift discs lie the closest and fall just below the low redshift relation (bottom left panel). On the other hand, we see what appears to be a fairly close association between the high redshift single component systems and the low redshift spheroid benchmark (top right panel). The high redshift bulges also overlap with the low-redshift-spheroid relation (middle right panel), but extend to significantly lower sizes as well.

To elucidate these issues further, we show in Fig. 14 the deviation distribution. We follow our methodology from before (see Section 5.2) and use Eq. 3 to calculate the deviation from the local benchmark relations. Here  $R_{
m observed}$  represents the sizes of the CANDELS-UDS data, and  $R_{\rm predicted}$ represents the predicted size using either the low redshift disc relation or low redshift high mass  $(\mathcal{M}_* > 10^{10} \mathcal{M}_{\odot})$  elliptical relation  $\mathcal{M}_{\star}-R_{\rm e}$  relation. On Fig. 14 each of the three populations are shown compared against either the local disc (left panel) or local spheroid (right panel) relations. The dark green dot-dashed lines show the scatter of the local data about the disc and spheroid relations. Note that the low redshift spheroid deviation (dark green dashed line, right panel) is only evaluated down to  $\mathcal{M}_* = 10^{10} \mathcal{M}_{\odot}$  to give a fair comparison, since a high mass only  $\mathcal{M}_{\star} - R_{\rm e}$  relation is used for evaluation. The solid black line shows the high-redshift single component sample while the dotted lines show the high-redshift discs (cyan) and bulges (magenta). Fig. 14 reiterates our findings from Fig. 13. Compared to the local spheroid benchmark, two populations show a plausible fit, the single component systems, and the high redshift bulges, albeit with a greater spread potentially indicative of the greater measurement error associated with fitting high redshift data. Whereas none of the distributions agree with the disc  $\mathcal{M}_{\star} - R_{\rm e}$  relation. However, the high redshift disc components have the largest overlap with our local disc  $\mathcal{M}_{\star} - R_{\rm e}$  relation.

In addition to the qualitative nature of Figures 13 and 14, we also perform a KS-test, as described in Sec. 3, to establish the association of the high redshift data with

our local disc and high mass spheroid  $\mathcal{M}_{\star} - R_{\rm e}$  relation. The resulting test statistics are shown in Table 4, and corroborate our visual inspections.

Putting aside external observations (e.g. visual morphology used as priors), our analysis suggests that the majority of high-redshift systems overlap satisfactorily with the low-redshift spheroid  $\mathcal{M}_{\star} - R_{\rm e}$  relation, with the exception of the high-redshift disc components. The obvious and simplest conclusion, is that we are essentially seeing bulge and spheroid formation/emergence at high-redshift, with some two-component systems existing, which adhere reasonable closely to the z=0 spheroid relations but with the highredshift discs somewhat more compact than their z = 0counterparts. This observation meshes well with the notion of rapid spheroid formation at high-redshift (z > 1.5; Tacchella et al. 2015) followed by disc growth at intermediate to lower-z (z < 1.5; e.g., Sachdeva et al. 2015), i.e., two-phase evolution as described in Driver et al. (2013). The obvious objections, however, are that observations of high-redshift systems generally show them to be visually clumpy, vigorously star-forming, and exhibiting clear evidence of systemic rotation, and hence are often described as disc-like (see for example Förster Schreiber et al. 2006, 2009; Green et al. 2010, and the recent review by Glazebrook 2013). However, these observations also show strong vertical velocity dispersions (e.g. Barro et al. 2014), verging on or exceeding the Toomre stability criterion (Toomre 1964), i.e. if these are discs they are highly unstable and unlike any type of disc seen locally.

At some level there is a semantic issue worth raising: When exactly does a spheroid become a spheroid, or a disc become a disc? Even in the ideal scenario of an entirely isolated collapsing gas cloud, it is likely to go through several star-burst phases, fragmentation and merging of these fragments before finally resembling what we consider a classical elliptical. Exactly at what point should we start calling such a system an elliptical, at the moment of first collapse or only after all star-formation has ceased and the system becomes dynamically relaxed? If the critical criteria are along the lines of the Hubble classification then clearly the high-redshift systems are not spheroids, however, it is also clear these are not conventional discs (smoothly rotating systems with minimal velocity dispersions and aspect ratios of 1:10), and the use of this terminology is equally misleading.

The mass-size relation essentially maps fundamental (conservable) quantities of mass and angular momentum. In this sense the mass-size relation is quite powerful and appears to be arguing that the majority of systems at high-redshift are, if not spheroids, proto-spheroids in the process of settling into spheroids. Those systems that do appear to have two-components exhibit bulges consistent or slightly smaller than low-redshift bulges (which could be due to increased nuclear activity making the bulges appear more compact), and discs which are offset to lower sizes, however disc growth is expected to continue to lower redshifts.

While the above paints a consistent and tantalising picture the caveats at this stage are significant. The analysis of the CANDLES-UDS data has been conducted by an independent group using distinct methods and strategies which could introduce systematic offsets. Distances are also based on photometric redshifts for the vast majority of the high-

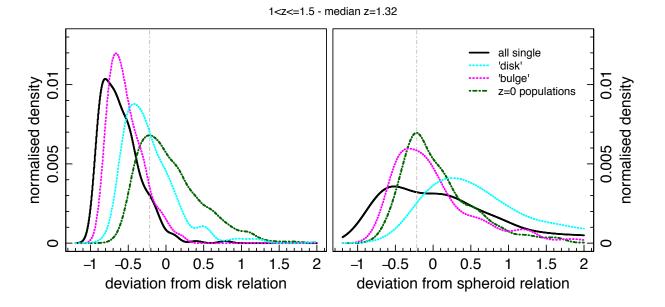


Figure 14. Shown is the same comparison as Fig. 10 but for the CANDELS high redshift data ( $1 < z \le 1.5$ ) compared to our final disc (left) and high mass elliptical (right,  $\mathcal{M}_* > 10^{10} \mathcal{M}_{\odot}$ )  $\mathcal{M}_{\star} - R_{\rm e}$  relations.

Case	D	p-values
(i) high redshift vs local disc	$\mathcal{M}_{\star} - I$	$R_{\rm e}$ relation
a) all high-redshift single	1	0.002
b) disc	0.62	
c) bulge	1	0
(ii) high redshift vs local hig	h mass s	pheroid $\mathcal{M}_{\star} - R_{\mathrm{e}}$ relation
a) all high-redshift single	0.5	0.474
b) disc	0.75	0.019
c) bulge	0.4	0.418

**Table 4.** Shown are the D- and p-values for a two tailed KS-test on the hypothesis that high redshift  $1 \leqslant z \leqslant 1.5$  galaxies and components are associated with either the local disc (i) or spheroid (ii)  $\mathcal{M}_{\star} - R_{\rm e}$  relation.

redshift sample. The sample size is also relatively small (subject to cosmic variance effects), and spans a particularly narrow mass range. Nevertheless, as we push the depth and area boundaries with facilities such as Euclid and WFIRST, the  $\mathcal{M}_{\star}-R_{\rm e}$  scaling relation shows great promise for providing not only a connection to the hydrodynamical simulations but also as a bridge between the low and high-redshift Universe.

#### 8 SUMMARY AND CONCLUSION

We have presented our bulge-disc decomposition catalogue for 7506 galaxies from the GAMA survey in the redshift range of 0.002 < z < 0.06 (Sec. 3). To overcome the limitations of the LM minimisation algorithm used in GALFIT, which can get trapped in local minima (especially for 2-component fits), we repeatedly fit our galaxy sample with varying starting points to map out the parameter space. For the single component galaxies we use a set of 33 combinations of the starting parameters, and for the 2-component fits we use a set of 88 combinations. We implement a screening process to prune bad fits and determine the final fitting values and errors from the median of the acceptable fits. We use the  $16^{th}$  and 84<sup>th</sup> percentile of the remaining output parameter distribution to determine the error on the median model combined with a 10% error floor. Through this strategy we reduce our catastrophic failure rate from  $\sim 20\%$  to  $\sim 5\%$ .

We then presented the  $\mathcal{M}_{\star}-R_{\rm e}$  relations of our sample by Hubble type and component with the component masses based on an estimation from the bulge and disc colours. Next we explored the association of the bulge and disc components with either the Sd-Irr or elliptical  $\mathcal{M}_{\star}-R_{\rm e}$  relation. We find that S(B)ab-S(B)cd galaxies likely consist of a disc plus a pseudo-bulge. Considering that a pseudo-bulge is a perturbation of the disc we decide that our late-type 2-component systems are best represented by a single component Sérsic fit. Thus we associate elliptical, early-type bulges and LBS for the spheroid  $\mathcal{M}_{\star}-R_{\rm e}$  relation and Sd-Irr, single component fit S(B)ab-S(B)cd galaxies and early-type discs for the final disc  $\mathcal{M}_{\star}-R_{\rm e}$  relation, which we provide as a definitive low redshift benchmark:

$$\begin{split} R_e &= 5.141 \Big(\frac{\mathcal{M}_\star}{10^{10}\mathcal{M}_\odot}\Big)^{0.274} \text{ for discs and,} \\ R_e &= 2.063 \Big(\frac{\mathcal{M}_\star}{10^{10}\mathcal{M}_\odot}\Big)^{0.263} \text{ for spheroids.} \end{split}$$

However, we caution the reader that the spheroid relation is heavily dominated by low mass galaxies. If a comparison to high mass spheroids is needed then the high mass elliptical  $\mathcal{M}_{\star} - R_{\rm e}$  relation (see Table 1) should be used in lieu of a curved spheroid relation.

Next we used our local disc and spheroid  $\mathcal{M}_{\star} - R_{\rm e}$  distributions to compare to data from the EAGLE simulation. We find a qualitatively good agreement between the sizes of the EAGLE data and our  $\mathcal{M}_{\star} - R_{\rm e}$  relations. This is not surprising as the sizes in EAGLE were calibrated using the (Shen et al. 2003)  $\mathcal{M}_{\star} - R_{\rm e}$  distribution. Hence the variance is of more interest as this has not been explicitly matched between the observed and simulated data. Comparing the scatter of the observed and simulated data we find that in almost all cases the simulated data has a smaller scatter which is unexpected, considering that the dark matter spin distribution is known to be fairly broad and we would expect the sizes and angular momentum distributions to have similarly broad distributions. Since we are comparing half-light sizes to half-mass sizes and components versus active and passive galaxies, we caution the reader to not over-interpret this comparisons. Instead we would like to highlight the potential of using the mass-size plane to compare observational and simulated data.

Finally, we compare our local  $\mathcal{M}_{\star}-R_{\rm e}$  relations to high redshift data from the CANDELS-UDS field. We concentrate on available data in a redshift range of  $1 < z \le 1.5$  with available single and 2-component fits (Mortlock et al. 2013; Margalef-Bentabol et al. 2016). We generally find that low-mass high redshift galaxies agree better with the local  $\mathcal{M}_{\star}-R_{\rm e}$  distributions than high-mass high redshift galaxies. Furthermore, high redshift systems, with the exception of disc components, more closely follow that of our local spheroid relation. The high redshift discs on the other hand follow the local disc  $\mathcal{M}_{\star}-R_{\rm e}$  relation, albeit offset to slightly smaller size. We interpret this as evidence for spheroid formation at high redshift and propose that further disc formation and/or growth does not occur until later times.

#### 9 ACKNOWLEDGEMENTS

RL would like to acknowledge funding from the International Centre for Radio Astronomy Research and the University of Western Australia. SB acknowledges the funding support from the Australian Research Council through a Future Fellowship (FT140101166). CL is funded by a Discovery Early Career Researcher Award (DE150100618).

GAMA is a joint European-Australasian project based around a spectroscopic campaign using the Anglo-Australian Telescope. The GAMA input catalogue is based on data taken from the Sloan Digital Sky Survey and the UKIRT Infrared Deep Sky Survey. Complementary imaging of the GAMA regions is being obtained by a number of independent survey programs including GALEX MIS, VST KiDS, VISTA VIKING, WISE, Herschel-ATLAS, GMRT and ASKAP providing UV to radio coverage. The VISTA

VIKING data used in this paper is based on observations made with ESO Telescopes at the La Silla Paranal Observatory under programme ID 179.A-2004. GAMA is funded by the STFC (UK), the ARC (Australia), the AAO, and the participating institutions. The GAMA website is http://www.gama-survey.org/.

#### REFERENCES

Abazajian K. N. et al., 2009, ApJS, 182, 543

Allen P. D., Driver S. P., Graham A. W., Cameron E., Liske J., de Propris R., 2006, MNRAS, 371, 2

Bahé Y. M. et al., 2016, MNRAS, 456, 1115

Baldry I. K. et al., 2014, MNRAS, 441, 2440

Baldry I. K. et al., 2010, MNRAS, 404, 86

Barro G. et al., 2014, ApJ, 795, 145

Berg T. A. M., Simard L., Mendel Trevor J., Ellison S. L., 2014, MNRAS, 440, L66

Bernardi M., Hyde J. B., Sheth R. K., Miller C. J., Nichol R. C., 2007, AJ, 133, 1741

Bernardi M., Meert A., Vikram V., Huertas-Company M., Mei S., Shankar F., Sheth R. K., 2014, MNRAS, 443, 874

Bernardi M., Roche N., Shankar F., Sheth R. K., 2011, MNRAS, 412, L6

Bertin E., 2013, PSF Extractor. Ascl:1301.001

Bertin E., Arnouts S., 1996, A&A Supplement Series, 117, 393

Bouwens R. J., Illingworth G. D., Blakeslee J. P., Broadhurst T. J., Franx M., 2004, ApJ, 611, L1

Bruce V. A. et al., 2012, MNRAS, 427, 1666

Bruce V. A. et al., 2014, MNRAS, 444, 1660

Bruzual G., Charlot S., 2003, MNRAS, 344, 1000

Buitrago F., Trujillo I., Conselice C. J., Bouwens R. J., Dickinson M., Yan H., 2008, ApJL, 687, L61

Buitrago F., Trujillo I., Conselice C. J., Häußler B., 2013, MN-RAS, 428, 1460

Buta R. J. et al., 2010, ApJS, 190, 147

Calzetti D., Armus L., Bohlin R. C., Kinney A. L., Koornneef J., Storchi-Bergmann T., 2000, ApJ, 533, 682

Cameron E., Driver S., Graham A., Liske J., 2009, ApJ, 699, 105
 Chabrier G., 2003, ApJ, 586, L133

Cheung E. et al., 2013, ApJ, 779, 162

Colless M. et al., 2003, ArXiv Astrophysics e-prints: astro-ph/0306581

Cook M., Lapi A., Granato G. L., 2009, MNRAS, 397, 534

Crain R. A. et al., 2015, MNRAS, 450, 1937

Daddi E. et al., 2005, ApJ, 626, 680

Dalcanton J. J., Spergel D. N., Summers F. J., 1997, ApJ, 482, 659

Davies L. J. M. et al., 2014, ArXiv e-prints: 1409.3574

Davis M. et al., 2003, Proc. SPIE, 4834, 161

de Souza R. E., Gadotti D. A., dos Anjos S., 2004, ApJS, 153, 411

Driver S. P. et al., 2006, MNRAS, 368, 414

Driver S. P. et al., 2011, MNRAS, 413, 971

Driver S. P., Robotham A. S. G., Bland-Hawthorn J., Brown M., Hopkins A., Liske J., Phillipps S., Wilkins S., 2013, MNRAS, 430, 2622

Driver S. P. et al., 2016, MNRAS, 455, 3911

Erwin P., 2015, ApJ, 799, 226

Fall S. M., Efstathiou G., 1980, MNRAS, 193, 189

Förster Schreiber N. M. et al., 2009, ApJ, 706, 1364

Förster Schreiber N. M. et al., 2006, ApJ, 645, 1062

Furlong M. et al., 2015a, ArXiv:1510.05645

Furlong M. et al., 2015b, MNRAS, 450, 4486

Gadotti D. A., 2008, MNRAS, 384, 420

Gadotti D. A., 2009, MNRAS, 393, 1531

Glazebrook K., 2013, PASA, 30, e056

Graham A. W., 2013, A review of elliptical and disc galaxy structure, and modern scaling laws, T.D.Oswalt, W.C.Keel, eds. Springer Publishing, pp. 91–140

Graham A. W., Dullo B. T., Savorgnan G. A. D., 2015, ApJ, 804, 32

Graham A. W., Guzman R., 2003, AJ, 125, 2936

Green A. W. et al., 2010, Nature, 467, 684

Grogin N. A. et al., 2011, ApJS, 197, 35

Häussler B. et al., 2007, ApJS, 172, 615

Hill D. T. et al., 2011, MNRAS, 412, 765

Holwerda B. W., Bouwens R., Oesch P., Smit R., Illingworth G., Labbe I., 2015, ApJ, 808, 6

Hopkins A. M. et al., 2013, MNRAS, 430, 2047

Hopkins P. F., Bundy K., Hernquist L., Wuyts S., Cox T. J., 2010, MNRAS, 401, 1099

Hopkins P. F., Bundy K., Murray N., Quataert E., Lauer T. R., Ma C.-P., 2009, MNRAS, 398, 898

Hubble E. P., 1926, ApJ, 64, 321

Hyde J. B., Bernardi M., 2009, MNRAS, 394, 1978

Kannappan S. J., Guie J. M., Baker A. J., 2009, AJ, 138, 579

Kelvin L. S. et al., 2012, MNRAS, 421, 1007

Kennedy R. et al., 2016, in prep.

Koekemoer A. M. et al., 2011, ApJS, 197, 36

Kormendy J., Bender R., 2012, ApJS, 198, 2

Lackner C. N., Gunn J. E., 2012, MNRAS, 421, 2277

Lagos C. d. P. et al., 2015, MNRAS, 452, 3815

Lang P. et al., 2014, ApJ, 788, 11

Lange R. et al., 2015, MNRAS, 447, 2603

L'Huillier B., Combes F., Semelin B., 2012, A&A, 544, A68

Lilly S. J. et al., 2007, ApJS, 172, 70

Liske J. et al., 2015, MNRAS, 452, 2087

Longhetti M. et al., 2007, MNRAS, 374, 614

Margalef-Bentabol B. et al., 2016, in prep.

McAlpine S. et al., 2015, ArXiv:1510.01320

Meert A., Vikram V., Bernardi M., 2015, MNRAS, 446, 3943

Mo H. J., Mao S., White S. D. M., 1998, MNRAS, 295, 319

Moffett A. J. et al., 2015, ArXiv e-prints

Mortlock A. et al., 2015, MNRAS, 447, 2

Mortlock A. et al., 2013, MNRAS, 433, 1185

Pastrav B. A., Popescu C. C., Tuffs R. J., Sansom A. E., 2013a, A&A, 553, 23

Pastrav B. A., Popescu C. C., Tuffs R. J., Sansom A. E., 2013b, A&A, eprint arXiv:1308.0458

Peng C. Y., Ho L. C., Impey C. D., Rix H.-W., 2010, AJ, 139,

Roberts M., Haynes M., 1994, araa, 32, 115

Robotham A. S. G., Obreschkow D., 2015, PASA, 32, 33

Romanowsky A. J., Fall S. M., 2012, ApJS, 203, 17

Sachdeva S., Gadotti D. A., Saha K., Singh H. P., 2015, MNRAS, 451, 2

Salo H. et al., 2015, ApJS, 219, 4

Schaye J. et al., 2015, MNRAS, 446, 521

Shen S., Mo H. J., White S. D. M., Blanton M. R., Kauffmann G., Voges W., Brinkmann J., Csabai I., 2003, MNRAS, 343, 978

Shibuya T., Ouchi M., Harikane Y., 2015, ApJS, 219, 15

Simard L., 1998, in Astronomical Society of the Pacific Conference Series, Vol. 145, Astronomical Data Analysis Software and Systems VII, Albrecht R., Hook R. N., Bushouse H. A., eds., p. 108

Simard L., Mendel J. T., Patton D. R., Ellison S. L., McConnachie A. W., 2011, ApJS, 196

Steinmetz M., Navarro J. F., 2002, New A, 7, 155

Tacchella S. et al., 2015, Science, 348, 314

Tasca L. A. M. et al., 2014, A&A, 564, L12

Taylor E. N. et al., 2011, MNRAS, 418, 1587 Toomre A., 1964, ApJ, 139, 1217 Toomre A., 1977, in Evolution of Galaxies and Stellar Populations, Tinsley B. M., Larson D. Campbell R. B. G., eds., p. 401

Trayford J. W. et al., 2015, MNRAS, 452, 2879

Trujillo I., Conselice C. J., Bundy K., Cooper M. C., Eisenhardt P., Ellis R. S., 2007, MNRAS, 382, 109

Trujillo I., Ferreras I., Rosa I. G. D., 2011, MNRAS, 415, 3903

Trujillo I. et al., 2006, ApJ, 650, 18

van den Bergh S., 1976, ApJ, 206, 883

van der Wel A. et al., 2014, ApJ, 788, 28

van Dokkum P. G. et al., 2008, ApJL, 677, L5

van Dokkum P. G. et al., 2010, ApJ, 709, 1018

Wei L. H., Kannappan S. J., Vogel S. N., Baker A. J., 2010, ApJ, 708, 841

Weinzirl T. et al., 2011, ApJ, 743, 87

Wisnioski E., Glazebrook K., Blake C., Poole G. B., Green A. W., Wyder T., Martin C., 2012, MNRAS, 422, 3339

Wright A. H., Robotham A. S., Bourne N., and Steve J. Maddox S. P. D., Dunne L., Cluver M., team G., 2016, in prep.

York D. G. et al., 2000, AJ, 120, 1579

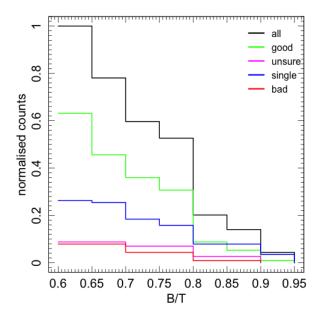


Figure A1. The distribution of galaxies with a median B/T > 0.6 and the visual decision whether they are 2-component, single component systems, or if the fit failed.

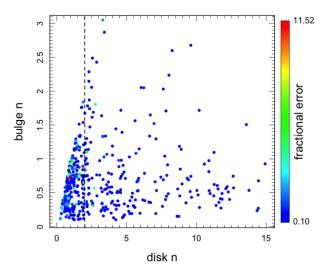


Figure A2. The disc versus bulge Sérsic index distribution for galaxies with bulge n < disc n. The points are colour coded by the fractional error on the disc Sérsic index. It is obvious that many of the fits for these galaxies converged (i.e. 10% error). We remove all converged and disc n > 2 galaxies from our final sample since we consider their fits unphysical.

#### APPENDIX A: FLAGGING

As described in Section 2.1 our sample of 7506~0.002 < z < 0.06 galaxies has been classified onto the Hubble type system as described in Moffett et al. 2015. The sample contains 860~E, 826~S(B)0-S(B)a, 1421~S(B)abc, 3531~Sd-Irr systems and 868~Little Blue Spheroids (LBS). The E, LBS and Sd-Irr's we consider single component systems, best fit by a single Sérsic profile, while we assume the remainder to be best described by a two-component, double Sérsic profile.

Following the decision on one or two components we flag the resulting fits for various criteria:

- (i) very high or low B/T (B/T > 0.8 and < 0.1 respectively)
- (ii) disc n > bulge n
- (iii) B/T reverses between minimum reduced  $\chi^2$  and median model
- (iv) bulge and disc position angle offset
- (v) the minimum reduced  $\chi^2$  solution is an outlier to the median value

Flags (i) through (iv) are only evaluated on the sample of two-component galaxies. Flag (v) is evaluated for all galaxies (i.e., single and two-component).

For our two-component sample we have a total of 962 galaxies with at least one flag ( $\sim 44\%$ ). This drops drastically when checking for galaxies with several flags, and we find only 206 galaxies with more than one flag raised. Our single component sample has a total of 164 flagged galaxies ( $\sim 3\%$ ).

To check whether any of the flags are more likely to produce unsatisfactory fits we visually inspect a random sample of 50 flagged two-component galaxies. We found that in many cases the median fit is acceptable and only flags (i) and (ii) are more likely to yield potentially bad fits.

#### (i) galaxies with high (low) B/T

We selected galaxies with a median B/T > 0.6 and visually inspected their minimum reduced  $\chi^2$  and median fit results as well as the convergence plots. For each galaxy we decided whether it is better fit with a double component or single component or if it has a bad or uncertain fit (i.e., the fit has bad apertures and no solution can be found, or it is unclear whether a 2-component fit is appropriate).

Fig. A1 shows the distribution of high B/T galaxies. We set all 46 galaxies with a B/T > 0.8 to a single component fit. Additionally, since our mean B/T error is 0.1, we also consider all galaxies with a B/T < 0.1 to be a single component galaxy. This adds another 42 galaxies for which the 2-component fit is considered not appropriate.

#### (ii) disc Sérsic index > bulge Sérsic index

As mentioned in Sec. 3.1.2 during the set up of our final B/T decomposition catalogue we screen galaxies for bulge and disc component switching. To do this we assume that the bulge  $R_e$  is smaller than the disc  $R_e$  and swap the assigned component for those galaxies where this is not the case. However, for galaxies where the bulge  $R_e$  is up to 10% larger than the disc  $R_e$  we also check whether the disc n is larger than the bulge n and only swap the assigned component if this is the case.

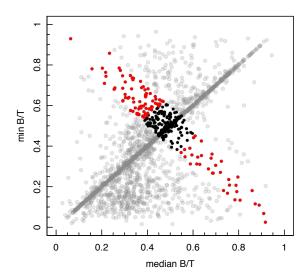


Figure A3. The median vs minimum reduced  $\chi^2$  distribution of the B/T in grey. The black points show the galaxies for which the sum of the min and median models lies between 0.9 and 1.1. The red points show the galaxies for which additionally the absolute difference between the min and median models B/T is larger than 0.15.

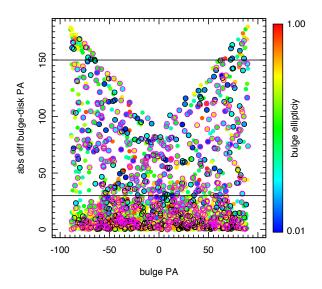


Figure A4. The absolute difference between bulge and disc position angle versus the bulge position angle. The points are coloured by the bulge ellipticity (e=1-b/a). The two vertical lines at 30 and 150 degrees show the range between which we flag the galaxies. The black circles show galaxies with bulge e<0.3 and the pink circles show galaxies with a disc e<0.3, which are excluded from the flagging.

However, even though the swapping ensures that our bulge is smaller than the disc it is not guaranteed that the bulge n is larger than the disc n. We check our 2-component fits and find 443 galaxies where the bulge n is smaller than the disc n. Fig. A2 shows the bulge versus disc n distribution coloured by the percentage error for the 443 flagged galaxies. Visually inspecting a number of the resulting fits we find that galaxies with a disc n > 2 and/or converged fits are typically bad and we remove them from consideration of the component samples. In total we remove 100 S(B)0-S(B)0a galaxies and 215 S(B)ab-S(B)cd galaxies.

### (iii) bulge and disc fraction reverse between median and minimum reduced $\chi^2$ model

We have 116 galaxies for which the median and minimum reduced  $\chi^2$  models have opposite B/T values. To establish this sample we select galaxies where the sum of the median and min B/T is between 0.9 and 1.1. The spread of this sum should include our inherent uncertainty in establishing the B/T, which is around 0.1-0.15. However, since this criterion alone also flags galaxies with a B/T=0.5 we add a second criterion that the absolute difference between the median and min B/T has to be larger than 0.15 (i.e. larger than our average uncertainty in establishing a B/T). Figure A3 shows the B/T distribution for the median versus minimum reduced  $\chi^2$  model, and highlights the process of flagging the galaxies with a reversed B/T between the median and minimum reduced  $\chi^2$  models.

#### (iv) bulge and disc position angle more than 30 degrees offset

This flag is defined by identifying all galaxies for which the bulge and disc position angle (PA) differs by more than 30 degrees. However, for round bulges or discs the PA is not very meaningful and even a large offset between bulge and disc PA is not indicative of a problem with the fitting results. To exclude these galaxies we set a second condition for this flag, namely that the bulge or disc ellipticity has to be large enough to have a clearly identifiable preferred major axis, i.e. e=1-b/a>0.3. This reduces the flagged galaxies to 100 objects.

Figure A4 shows the distribution of the absolute PA offset vs bulge PA, with the points coloured by their ellipticity. The horizontal lines show the angle offset between which we flag, and the objects circled show galaxies with e<0.3 which are excluded from the flagging.

#### (v) minimum reduced $\chi^2$ solution represents an outlier to the median model

Finally we flag galaxies for which the minimum reduced  $\chi^2$  solution represents and outlier, i.e. where the error on the median values does not include the minimum reduced  $\chi^2$  solution. For minimum reduced  $\chi^2$  solutions where the parameters are smaller than the median values we consider the error on the lower end, and vice versa where the median is smaller we consider the error on the higher end.

We tested size, magnitude and Sérsic index distributions for outliers. We only consider the minimum reduced  $\chi^2$  values to be a true outlier, and possibly better fit, if both the disc and bulge values fall outside the error range. Additionally we check that the error on the median is larger than our assumed error floor, otherwise we consider the various fits of the galaxy to

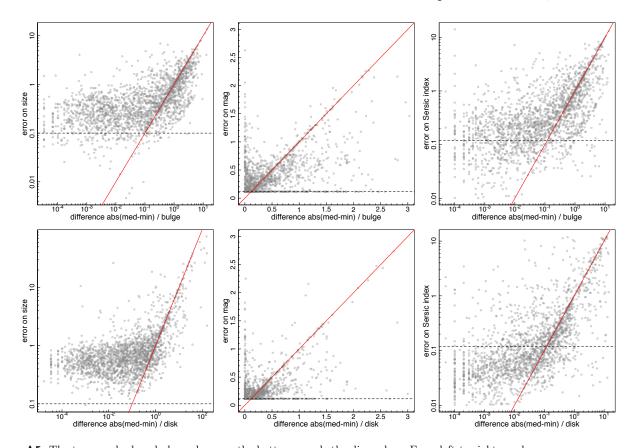


Figure A5. The top panels show bulge values, on the bottom panels the disc values. From left to right we show: difference in median and minimum reduced  $\chi^2$  model sizes versus the median size error difference in median and minimum reduced  $\chi^2$  model magnitudes versus the median magnitude error difference in median and minimum reduced  $\chi^2$  model Sérsic indices versus the median Sérsic index error The red lines show the 1:1 correspondence, the black dashed lines indicate our lower limits on the errors. Everything to the right of the red lines and above the dashed lines is flagged.

have converged and thus the minimum reduced  $\chi^2$  solution likely represents a failed fit.

Figure A5 shows the distribution of the absolute difference between the median and minimum reduced  $\chi^2$  values versus the error for the median values. The left hand panel shows the bulges and the right hand panel the disc distribution. From top to bottom we investigate size, magnitude, and Sérsic index.

The red lines show the 1:1 correspondence, i.e. where the minimum reduced  $\chi^2$  value lies on the edge of the median error distribution. The horizontal dashed lines shows our error floor, and we assume that most models converged to give rise to such a low error.

All galaxies to the right of the red lines and above the dashed black lines are potentially bad. We do, however, only flag the ones that are bad for both the disc and bulge values. This gives rise to 54 galaxies with flagged sizes, 393 galaxies with flagged magnitudes, and 88 galaxies with flagged Sérsic indices. In total this results in 449 galaxies for which the minimum reduced  $\chi^2$  solution is an outlier to at least one of the three fitting parameters tested.

The equivalent test on our single component fits finds 73 (42) spheroid (disc) sizes, 27 (30) spheroid (disc) magnitudes, and 62 (48) spheroid (disc) Sérsic indices flagged. In total this equates to 90 (74) spheroid (disc) galaxies with at least one of the parameters flagged ( $\leq 5$ ).

### APPENDIX B: SIMULATION COMPARISONS

In Figures B1 and B2 we show the distribution of the vertical scatter from the  $\mathcal{M}_{\star} - R_{\rm e}$  relation for discs and spheroids. The black histogram shows the distribution of our data and the blue line is a fit of a normal distribution to it. The red histograms show the distribution of the z=0 EAGLE data with the red line being the normal distribution fit to it.

Figure 12 is essentially derived from these plots where the standard deviation (SD) points refer to the sigma of the Gaussian fit and the  $16^{th}$  and  $84^{th}$  percentiles are the width of the underlying distribution. Typically these measurements of sigma should be the same, however, in some cases the distributions have tails towards higher deviations which causes the  $16^{th}/84^{th}$  percentile and SD measurements to be different (as seen in Fig. 12).

32

Figure B1. a) Distribution of the half-light radius scatter of discs for bins in stellar mass (shown at the top of each panel) for our data. The blue line shows the normal distribution fit to the observed data (black histogram) and the red line is the normal distribution fit for the simulated data (red histogram), with the caveat that we use the half-mass radius for star forming galaxies in the case of EAGLE. The legend on the left shows the best fit values for the observed data and on the right for the simulated data.

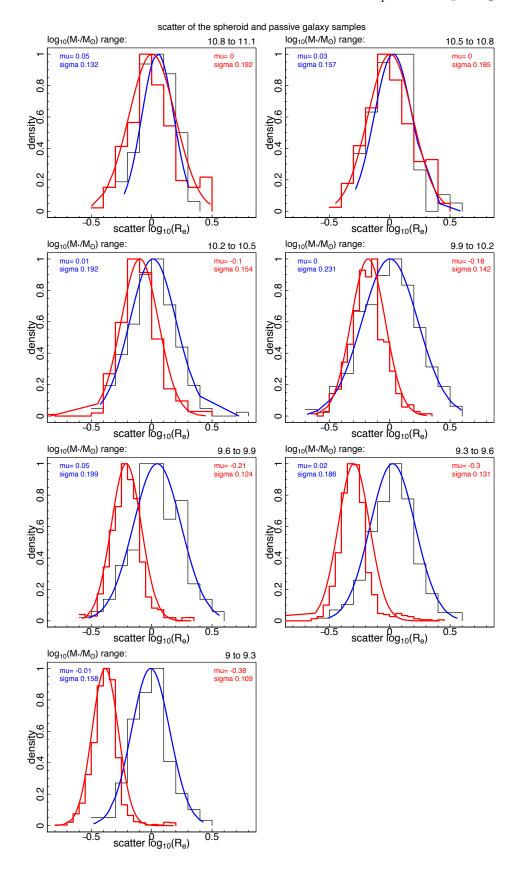


Figure B2. Same as Fig. B1 but shown is the  $log(R_e)$  scatter for the spheroid component distribution and in the case of EAGLE we show the scatter of the half-mass radius of passive galaxies.