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Discovery and confirmation of the shortest gamma ray burst from a collapsar

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Gamma-ray bursts (GRBs) are among the brightest and most energetic events in the universe. The duration and hardness distribution of GRBs has two clusters¹, now understood to reflect (at least) two different progenitors². Short-hard GRBs (SGRBs; $T_{90} < 2$ s) arise from compact binary mergers, while long-soft GRBs (LGRBs; $T_{90} > 2$ s) have been attributed to the collapse of peculiar massive stars (collapsars)³. The discovery of SN 1998bw/GRB 980425⁴ marked the first association of a LGRB with a collapsar and AT 2017gfo⁵/GRB 170817A/GW170817⁶ marked the first association of a SGRB with a binary neutron star merger, producing also gravitational wave (GW). Here, we present the discovery of ZTF20abwysq (AT2020scz), a fast-fading optical transient in the *Fermi* Satellite and the InterPlanetary Network (IPN) localization regions of GRB 200826A; X-ray and radio emission further confirm that this is the afterglow. Follow-up imaging (at rest-frame 16.5 days) reveals excess emission above the afterglow that cannot be explained as an underlying kilonova (KN), but is consistent with being the supernova (SN). Despite the GRB duration being short (rest-frame T_{90} of 0.65 s), our panchromatic follow-up data confirms a collapsar origin. GRB 200826A is the shortest LGRB found with an associated collapsar; it appears to sit on the brink between a successful and a failed collapsar. Our discovery is consistent with the hypothesis that most collapsars fail to produce ultra-relativistic jets.

On August 26, 2020, at 04:29:52 UT, the Gamma-ray Burst Monitor (GBM) on board the *Fermi* Gamma-ray Space Telescope detected GRB 200826A with duration (T_{90}) of 1.14 ± 0.13 seconds in the 50–300 keV energy range. In addition to *Fermi* (GBM trigger 620108997), GRB 200826A was detected by four other Interplanetary Network (IPN) instruments (see Methods).

The gamma-ray properties alone do not always yield an unambiguous classification. Some SGRBs show afterglow and host properties akin to LGRBs, e.g. Ref. ⁷; some LGRBs show no evidence for collapsars to deep limits akin to SGRBs, e.g. Ref. ⁸. Based solely on T_{90} ⁹, GRB 200826A has a SGRB probability of $65\%_{-11}^{+12}$. Also taking into consideration the E_{peak} parameter of a Comptonized model fit to the single spectrum over the duration of the burst (see

Extended Data Table 5), the probability that GRB 200826A is a SGRB increases to 74% (see Figure 1 and Methods). However, based on rest-frame energetics, GRB 200826A is not consistent with the SGRB population (right panel in Extended Data Figure 1).

Starting 4.2 hours after the GRB, we observed 180 sq. degrees of the *Fermi*-GBM localization with the Zwicky Transient Facility (ZTF)¹⁰ (see Fig. 2). At 17 hours after the GRB, IPN triangulated the source to a smaller region. ZTF20abwysqy was the only candidate that passed our alert filtering scheme and was also inside the IPN region (see Methods). We discovered ZTF20abwysqy at a brightness of $g = 20.86 \pm 0.04$ mag (AB system). The previous upper limit ($g > 21.3$ mag at 5σ) was 17.3 hours before the GRB, as part of the nominal all-sky survey mode. In addition to the spatial coincidence, ZTF20abwysqy was associated with a fading X-ray counterpart¹¹ and variable radio emission¹², confirming ZTF20abwysqy as the afterglow of GRB 200826A (see Extended Data Table 1, Extended Data Table 3 and Extended Data Table 4).

ZTF20abwysqy was discovered in a galaxy with archival detections in Pan-STARRS 1 (PS1)¹³ and Legacy Survey (LS)¹⁴, at a LS photometric redshift of $z_p = 0.71 \pm 0.14$. The offset between the host galaxy’s centroid and the transient is $0.18 \pm 0.05''$, corresponding to a physical distance of 2.09 ± 0.6 kpc. We acquired a Gran Telescopio Canarias (GTC) spectrum of the galaxy; we see strong [OII] and [OIII] lines at $z = 0.748$ (see Extended Data Table 2 and Extended Data Figure 2). With both spectral and photometric spectral energy distribution (SED) fitting, we infer a stellar mass for the host galaxy of $\sim 10^{9.7} M_\odot$ (see Extended Data Figure 2, Extended Data Figure 3 and Extended Data Table 6 in Methods), which is near the distribution peak for LGRBs¹⁵, while below the median for SGRBs based on Ref.¹⁶.

We model the GRB afterglow using the standard synchrotron fireball model to constrain parameters related to the energy and geometry of the GRB central engine (see Methods). Electrons in the circumburst medium are accelerated by the shock wave and reach a power-law energy distribution characterized by the index p , $N(E) \propto E^{-p}$. This results in a SED described by a series of broken power laws, e.g. Ref. ¹⁷ (See Extended Data Figure 4 and Extended Data Figure 5). The associated isotropic kinetic energy of $E_{K,iso} = 6.0^{+51.3}_{-4.4} \times 10^{52}$ erg lies in the top 5% of the $E_{K,iso}$ distribution for SGRBs¹⁶, but within the 90% confidence range of the LGRB energy distribution¹⁸ (see Methods). Our data can only loosely constrain the circumburst density. The upper end of the distribution is consistent with the values found for the median circumburst densities¹⁹ of LGRBs,

while the lower end is more representative of SGRBs¹⁶.

In the collapsar scenario, a high-velocity stripped-envelope SN (SN Ic-BL) should follow the GRB detection³. To test this scenario, we used the Gemini Multi-Object Spectrograph (GMOS-N; see Methods) to acquire *r*- and *i*-band images of ZTF20abwysqy on three different epochs: ~ 28 , ~ 45 , and ~ 80 days after the GRB trigger (epoch 1, 2 and 3 respectively). Using epoch 3 as the reference, we undertook image subtraction using two different subtraction algorithms (see Methods for details). On epoch 1, our Gemini observations show evidence of a transient with an *i*-band magnitude of 25.45 ± 0.15 mag (with an *i*-band 5σ limit of 25.9 mag, see Extended Data Table 1); there is no source in the *r*-band observations up to a 5σ limit of 25.4 mag (see Fig. 3). On epoch 2, we do not detect a source up to a 5σ limit of 25.5 mag in the *i*-band and 25.7 mag in the *r*-band. Thus, at a rest-frame time of ~ 16 days, the foreground extinction-corrected absolute magnitude of ZTF20abwysqy is $M_i = -18.0$ mag.

To understand the source of the *i*-band excess, we use Markov chain Monte Carlo (MCMC) and full forward modeling of all multi-band observations excluding the GMOS detection to compare three scenarios: an afterglow only, an afterglow plus a KN, and an afterglow plus a SN. The KN model is based on a best-fit template to AT2017gfo²⁰ scaled by a compilation of SGRB-KN candidates²¹; the collapsar model uses a SN1998bw template²² with stretch and scale parameters drawn from a prior that is consistent with the historical GRB-SN sample²³. We dismiss the afterglow-only and afterglow-plus-KN models because they predict *i*-band flux at the time of the GMOS data point that is too faint by $1.6^{+1.8}_{-0.3}$ mag – inconsistent with the observations at the $\sim 5\sigma$ level (see Extended Data Figure 6 and Extended Data Table 7).

We repeated the analysis of all three scenarios including the GMOS data point in order to do Bayesian model comparison. The Bayes factor between the afterglow-plus-KN and afterglow-only models is ~ 1 , indicating that neither model is favored over the other. However, the Bayes factor between the afterglow-plus-SN and afterglow-only model is $10^{5.5}$, indicating that the afterglow-plus-SN model is strongly favored. Based on the compilation of Ref. ²³ of GRBs with associated SNe, GRB 200826A is the shortest LGRB found with an associated collapsar, with an observed T_{90} of 1.13 s and rest-frame T_{90} of 0.65 s (see Figure 1). The second shortest LGRB with a SN is GRB 040924, with a rest-frame T_{90} of 1.29 s²³.

In the conventional fireball model of a LGRB, the rest-frame duration of the prompt emission, t_γ , is the difference between the duration of the activity of the central engine, t_e , and the time required to break out of the envelope, t_b ²⁴. The short duration of GRB 200826A suggests that $t_e \gtrsim t_b$, which might imply that its central engine is active only very briefly, or that the stellar envelope is unusually thick, as compared to other LGRBs, perhaps requiring a nonstandard progenitor²⁵. Then it is natural to infer that there must be collapsars for which $t_e < t_b$ that fail to produce a fireball. For $t_e \lesssim t_b$, the jet may fail to clear a path for itself and remains cocooned within the exploding star, yet the cocoon itself may still produce a mildly relativistic shock breakout that manifests as a soft, quasi-thermal, low-luminosity GRB (llGRB) like GRB 980425/SN 1998bw or GRB 060218/SN2006aj^{26,27}. For $t_e \ll t_b$, there may be no prompt emission at all.

GRB 200826A, then, may sit on the brink between a successful LGRB and a failed one. The sign of one continuous physical parameter, $t_e - t_b$, switches a collapsar-powered GRB discontinuously between two dominant emission mechanisms: internal shocks in a relativistic jet or mildly relativistic shock breakout at the surface of the exploding star. These two mechanisms correspond to two widely separated and disconnected regions in the hardness-duration-fluence phase space: the traditional border between SGRBs and LGRBs, and the exceptionally long and soft region occupied by llGRBs and X-ray flashes (XRFs). The local rate of LGRBs ($\sim 1 \text{ Gpc}^{-3} \text{ yr}^{-1}$)²⁷ is 2 orders of magnitude lower than the rate of llGRBs ($300 \text{ Gpc}^{-3} \text{ yr}^{-1}$)²⁷ and 3 orders of magnitude lower than the rate of SNe Ic-BL ($4500 \text{ Gpc}^{-3} \text{ yr}^{-1}$)²⁸. It is understood that LGRBs, and to a lesser extent, llGRBs, are collimated and relativistically beamed, suppressing the detection of events in which the jet is pointed away from our line of sight. However, even allowing for a beaming correction to their rates of 100 for LGRBs and 1–10 for llGRBs²⁷ (but no beaming correction for SNe), LGRBs still occur at a rate ($100 \text{ Gpc}^{-3} \text{ yr}^{-1}$) that is up to an order of magnitude lower than that of llGRBs ($300\text{--}3000 \text{ Gpc}^{-3} \text{ yr}^{-1}$) or that of SNe Ic-BL, suggesting that majority of the collapsars fail to produce an ultra-relativistic jet and instead drive a wide-angle or nearly isotropic and only mildly-relativistic cocoon.

Based on indirect evidence from the host, afterglow, and gamma-ray properties, it has been argued that as many as 84% of bursts detected by the *Neil Gehrels Swift Observatory*²⁹ (*Swift*) and 40% of the *Fermi* bursts that are nominally SGRBs ($T_{90} < 2 \text{ s}$) are actually misclassified LGRBs³⁰. If this is correct, then one would expect many more short-duration GRBs with collapsars that are on the edge between success and failure of the jet. The discovery of GRB 200826A, an SGRB

imposter, lends credence to this bold claim, and would suggest that the rates of such short-duration LGRBs is comparable to the rate of llGRBs, up to a few hundred $\text{Gpc}^{-3}\text{yr}^{-1}$. Thus, our discovery upholds the hypothesis that most collapsars fail to produce jets.

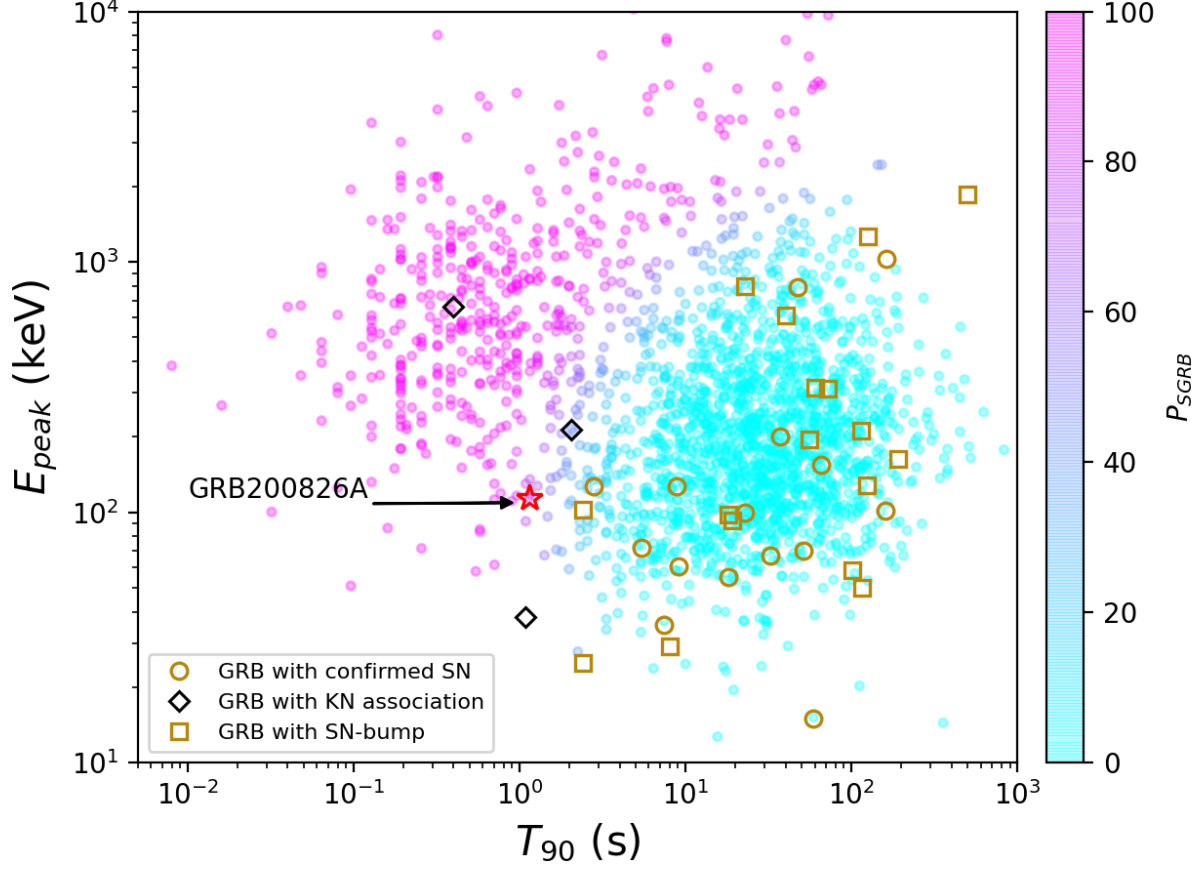


Figure 1: **Gamma-ray properties of GRB 200826A in context.** The peak energy based on a Comptonised fit, E_{peak} (keV), vs. the time-integrated, T_{90} (s), for 2310 *Fermi* GBM GRBs. The data are fit with two log-normal distributions for the two GRB classes. The color of the data points indicate the probability with magenta being 100% SGRB and cyan being 100% LGRB. GRB 200826A is surrounded by a red star with a SGRB probability of 74% (See Methods). Yellow squares show LGRBs with SN-bumps, yellow circles show LGRBs with spectroscopically confirmed SN²³, black diamonds show SGRBs with claimed KN excess²¹.

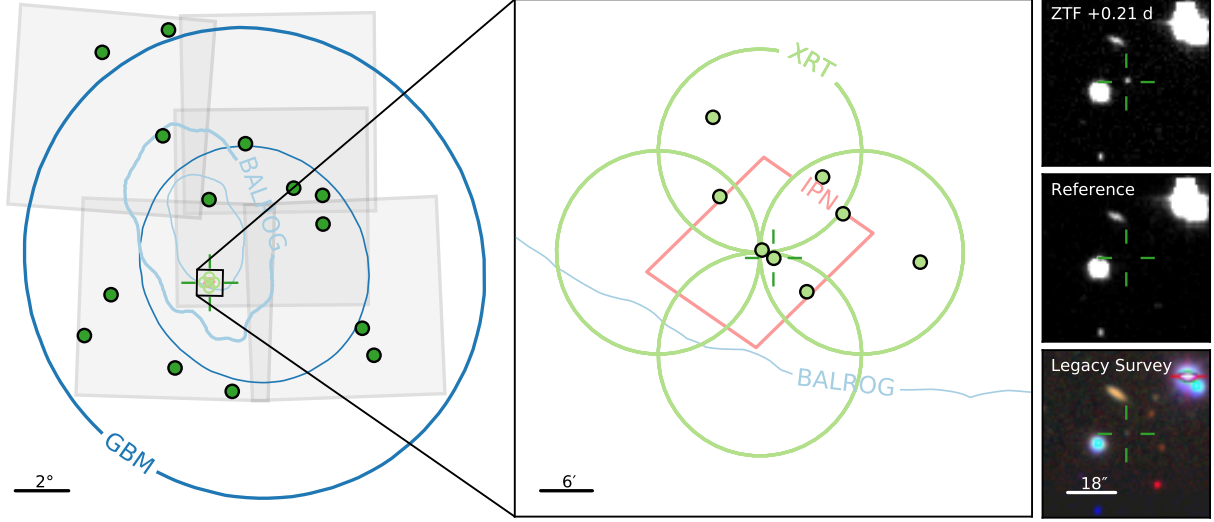


Figure 2: **Discovery of the afterglow of GRB 200826A.** It was found at the position $\alpha = 00^h27^m08.542^s$, $\delta = +34^h01^m38.327^s$ (J2000) with an uncertainty of 0.08.” Contours in the left panel represent the *Fermi* GBM 90% (thick) and 50% (thin) credible regions from the official *Fermi* GBM localization (dark blue) and BALROG (light blue). The filled gray squares show the ZTF fields observed and the dark green dots are the positions of the ZTF Night 1 optical candidates. In the middle panel, 3σ IPN triangulation is shown in pink; the four fields of the four XRT tiled observations are shown as light green circles and the XRT candidates are light green dots. The position of the afterglow is marked by the dark green reticle. The right-hand panels are centered at the position of ZTF20abwysqy. The cutouts, from top to bottom, show the ZTF discovery image, the ZTF stacked reference image, and a false color image showing the host galaxy from Legacy Survey DR8. In the cutouts, North is up and East to the left.

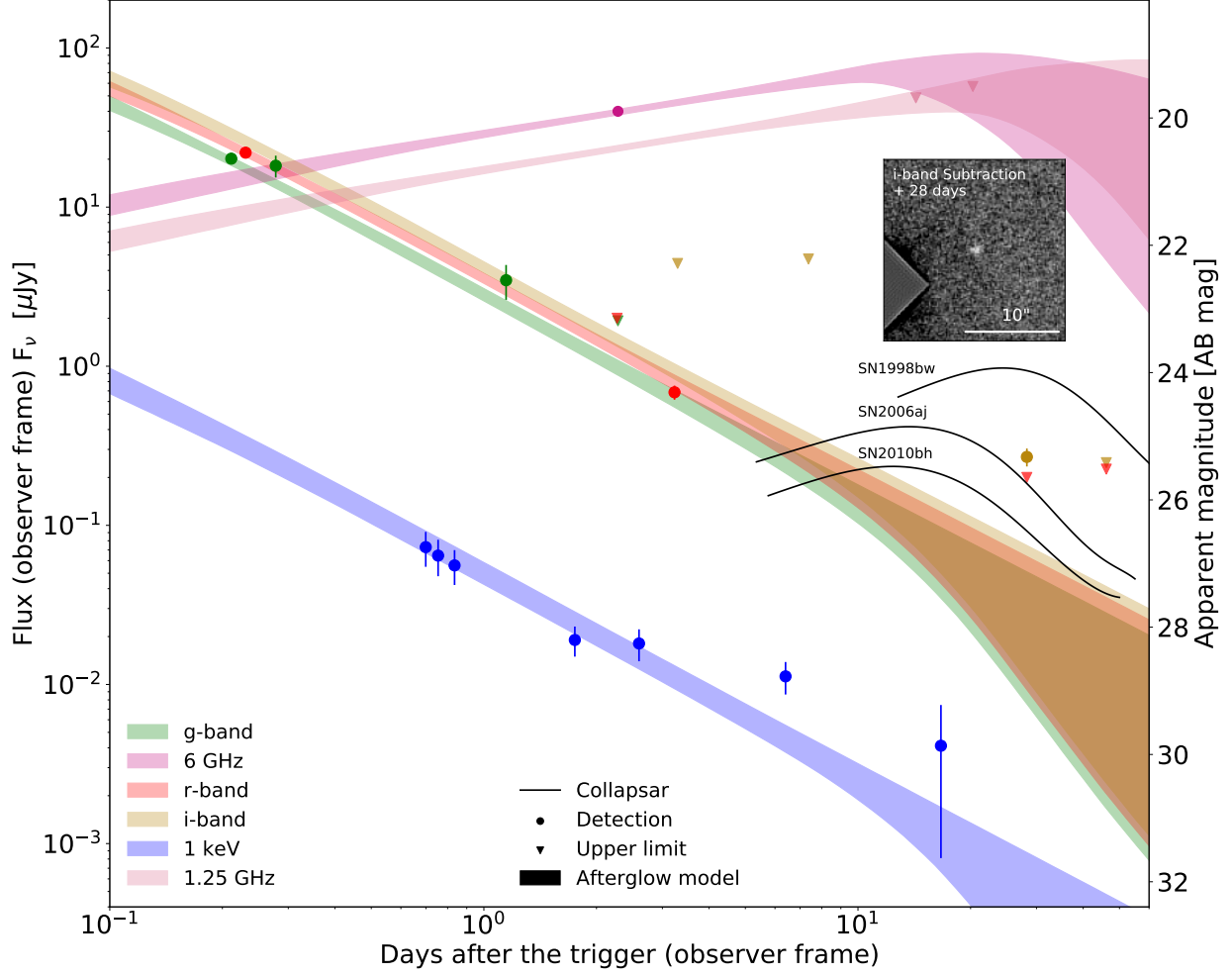


Figure 3: **Panchromatic afterglow and collapsar confirmation.** The available multiwavelength light curve data is over-plotted with the best fits from the `afterglowpy` modeling assuming an ISM-like environment. Detections are shown as circles with their respective error bars, and upper limits are shown as inverted triangles. The optical g -, r - and i - bands are shown in green, red and yellow, the XRT 1 keV data is shown in blue, the VLA in fuchsia, while the GMRT data is presented in pink. We show the K-corrected light curves of three well-studied GRB-SNe (SN1998bw, one of the brightest; SN2006aj and SN2010bh, two of the faintest) with solid black lines. The Gemini GMOS-N i -band detection is shown at day 28.28, as a yellow circle, and is consistent with the collapsar population. We show a cutout of the i -band ZOGY subtraction, revealing our 25.45 mag detection with a ZOGY corrected score of 4.

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Methods

1 Discovery

Background. The traditional paradigm of classifying a GRB based only on its gamma-ray properties is debated^{9,31}. Although the two types of progenitors broadly map to the time duration of the signals, there is no clear boundary in the bimodal distribution. Some SGRBs show afterglow and host properties akin to LGRBs, e.g. Ref. ^{7,34}; some LGRBs show no evidence for collapsars to deep limits akin to SGRBs, e.g. Ref. ^{35,36}.

In the mid 2000s, SGRB-related breakthroughs triggered by the *Neil Gehrels Swift Observatory*²⁹ and HETE-II³⁸ included detections of SGRBs X-ray afterglows, identification of likely host galaxies³⁹, and the first SGRB optical afterglow^{40,41}. Yet, in the first decade of optical follow-up of *Swift* SGRBs, only about 30 optical/NIR afterglows were detected^{16,42} and even fewer had associated redshifts. On 2017 August 17, the joint detection of the binary neutron star (BNS) merger (GW170817) in GWs with LIGO and Virgo⁴⁴, and in gamma-rays by *Fermi* GBM⁴⁵ and INTEGRAL^{46–48}, unequivocally confirmed BNS mergers as at least one of the mechanisms that can produce a SGRB. The merger illuminated the entire electromagnetic spectrum⁴⁹ and the optical/NIR emission provided robust evidence of a radioactively powered KN^{50–59}.

The bulk of our knowledge about SGRBs comes from the well-studied *Swift* SGRB afterglow sample; however, the low *Swift* local rate of SGRBs⁶⁰ has hampered the detection of a GW170817-like event. On the other hand, the *Fermi* GBM is arguably the most prolific engine of discovery for *compact binary mergers*, as it detects ≈ 1 SGRB each week, which is four times the rate of *Swift*^{61,62}, and comparable to the rate of the LIGO and Virgo detectors when in observing mode (≈ 1.5 per week during O3a⁶³). Unfortunately, there is a relative paucity of electromagnetic observations of *Fermi* GRBs and LIGO/Virgo GW events alike because it is challenging to pinpoint them within their positional uncertainties of tens to thousands of square degrees^{64–66}.

In order to better understand the phenomenology of GRBs and compact binary mergers, wide field of view (FOV) optical instruments have looked for *Fermi* GRBs and LIGO/Virgo GW counterparts. In 2013, the Palomar Transient Factory (PTF)⁶⁷ used a 7 deg² camera on the Palomar 48 inch Oschin telescope (P48) to discover the first optical afterglow of a GRB based *solely* on a *Fermi* GBM localization⁶⁸ and subsequently found afterglows of 7 other LGRBs⁶⁹. Now, the Zwicky Transient Facility (ZTF^{10,71,72}), a 47 deg² camera mounted at the P48 telescope, has

enabled searches an order of magnitude faster in areal and volumetric survey speed and has been used to search the coarse error regions of *Fermi* GBM SGRBs⁷³ and LIGO/Virgo events^{74–77}. Both PTF and ZTF have also discovered afterglow-like transients with optical, X-ray, and radio emission, but no gamma-ray counterpart^{78–80}.

Gamma-ray detections. *Fermi* (GBM trigger 620108997)⁸¹, and the IPN instruments – *AGILE*⁸², *INTEGRAL* (SPI-ACS), *Mars-Odyssey* (HEND), *Konus-Wind*⁸³, and *AstroSat*⁸⁴ detected the burst. The GBM localization calculated by the ground software^{85,86} is RA = 4.7, Dec = 35.3 (J2000 degrees, equivalent to J2000 00h 18m, 35d 17') with a statistical uncertainty of 1.7 degrees⁸¹. The BALROG localization is consistent with this⁸⁷. The final GBM localization of the GRB was available a few minutes after the trigger, and it enclosed 339 (63) deg² within the 95% (50%) credible region.

The light curve shows a bright GRB featuring three distinct peaks in the 50–300 keV energy range, with a signal-to-noise ratio (S/N) of ~ 100 and a T_{90} duration of 1.14 ± 0.13 s. This GRB is within the top 36% in terms of fluence (in the 50–300 keV energy band as measured over T_{100}) from bursts recorded in the *Fermi* GBM catalog⁶². Spectral analysis was performed using the RMFIT 4.4.2 light curve and spectral analysis software^a in the interval of T_0 to $T_0 + 1.152$ s encompassing the T_{90} start time and duration. The detector selection for the analysis consisted of the brightest NaI detectors with angle $< 60^\circ$ (N6, N7, N8, N9, NB11) and BGO detector B1. NaI and BGO detector data from ~ 8 keV to ~ 900 keV and ~ 300 keV to ~ 40 MeV respectively was used for the analysis. The Band Model⁸⁸, a standard model for GRB analysis⁸⁹, provided best fit with parameters consistent with the initial analysis⁹⁰, where $E_{peak} = 88.9 \pm 3.2$ keV, $\alpha = -0.26 \pm 0.07$, and $\beta = -2.4 \pm 0.1$. The peak flux in the 64 ms time (10 keV - 1 MeV), measured from $T_0 + 1.152$ s, is 64.3 ± 2.1 ph cm⁻² s⁻¹. The fluence over T_{90} (10 keV to 1 MeV) is $42.6 \pm 0.2 \times 10^{-7}$ erg cm⁻² (the fit merit of the spectral fit is 1.09).

A time resolved spectral analysis was performed from T_0 to $T_0 + 1.152$ s. The data was binned by S/N, with at least 20σ for the S/N in each bin. For 6 of the bins, the best-fit parameters of the Band Model are displayed in Extended Data Table 5. The parameters for the Band model in the initial time interval were poorly constrained so a cutoff power law (Comptonised) model⁹¹ was fitted to this interval. The Band Model was primarily used as this was the best constrained model

^aAvailable at the *Fermi* Science Support Center: <https://fermi.gsfc.nasa.gov/ssc/data/p7rep/analysis/rmfit/>

for the time integrated analysis. The Castor C-stat statistic⁹² was used for its robustness when dealing with low count statistics, while the background was modeled as a first order polynomial over all qualifying detectors, as discussed above, using TTE data products. Spectral evolution of the E_{peak} parameter is observed over the duration of the burst, and is found to exhibit hardness-intensity tracking behavior^{91,92}.

The distribution of GRB T_{90} durations demonstrates a bimodality that is best fit with double log-normal components. This suggests two distinct but overlapping progenitor probabilities, and by extension, classifications for short and long GRBs. Using only the T_{90} , and the parameters derived for the *Fermi* GBM distribution of bursts in Ref.⁹, we determine a SGRB probability of $64.94^{+11.59}_{-11.30}\%$ for GRB 200826A.

According to BATSE⁹³, the overlap of GRB T_{90} duration distributions occurs at ~ 2 s. In the past, GRB hardness ratios, quantifiable for the majority of GRBs in the absence of location information, have been used as a measure for the spectral hardness. Hardness ratio is presented for the GBM catalogue events in Ref.⁶², where the data shows an anti-correlation between hardness and duration; however, there is a large scatter in the data. We derive a T_{90} value of 4.2 s as the duration where an event has equal probability of being in the short or long class⁶² based on 2353 GRBs. We choose the E_{peak} parameter instead of hardness for separating LGRB and SGRB populations as it does not require set energy bands for a comparison to be made⁹⁴. In Figure 1 we presents a log-log plot of the T_{90} and E_{peak} parameter, for which we fit two log-normal distributions, one for each class, and the probability that an event is an SGRB is translated into a colour. From this analysis, the probability that GRB 200826A belongs to the short class is 74%⁹⁵.

GRB 200826A triggered *Konus-Wind* (KW) at $T_{0KW} = 16195.106$ s UT (04:29:55.106). The propagation delay from Earth to Wind is 2.540 s for this GRB; correcting for this factor, the KW trigger time corresponds to the Earth-crossing time 16192.566 s UT (04:29:52.566). The burst light curve shows a multi-peaked pulse which starts at about $T_{0KW} - 0.10$ s and has $T_{100} = 0.97$ s, $T_{90} = 0.67^{+0.13}_{-0.03}$ s, and $T_{50} = 0.28^{+0.04}_{-0.03}$ s measured in the 80–1300 keV band. Considering only the duration, the T_{90} and T_{50} of GRB 200826A are consistent with the short GRB population in the KW sample. The T_{90} and T_{50} durations at which a KW-detected GRB has an equal probability of being short- or long-duration are 2 s and 0.7 s, respectively⁹⁶. The spectral lag between the 20–80 keV and 80–330 keV 16 ms light curves is 30 ± 11 ms, consistent with the bulk of KW short GRBs⁹⁷.

During the burst, KW measured five spectra in the 20 keV–10 MeV band. The first four with 64 ms accumulation time cover the interval from $T_{0\text{KW}}$ to $T_{0\text{KW}} + 0.256$ s and the fifth, from $T_{0\text{KW}} + 0.256$ s to $T_{0\text{KW}} + 8.448$ s. The time-averaged spectrum of the burst (measured from $T_{0\text{KW}}$ s to $T_{0\text{KW}} + 8.448$ s) is best fit in the 20 keV–2 MeV range by the Band function with a low-energy photon index $\alpha = 1.26^{+1.91}_{-1.12}$, a high-energy photon index $\beta = -2.32^{+0.12}_{-0.15}$, and a spectrum peak energy $E_{\text{peak}} = 67^{+13}_{-11}$ keV ($\chi^2/\text{d.o.f.} = 50/59$). The burst had a fluence of $4.60^{+0.71}_{-0.60} \times 10^{-6}$ erg cm $^{-2}$, and a 16-ms peak flux, measured from $T_{0\text{KW}} + 0.544$ s, of $9.81^{+1.83}_{-1.64} \times 10^{-6}$ erg cm $^{-2}$ s $^{-1}$ (both in the 20 keV–10 MeV energy range). Using the spectroscopic redshift $z = 0.7481^{98}$ we have estimated the following rest-frame parameters: the isotropic energy release E_{iso} is $7.17^{+1.11}_{-0.94} \times 10^{51}$ erg, the peak luminosity L_{iso} is $2.67^{+0.50}_{-0.45} \times 10^{52}$ erg s $^{-1}$, and the rest-frame peak energy of the time-integrated spectrum, $E_{\text{peak},z}$ is 117^{+23}_{-19} keV.

With these values, GRB 200826A is within the softest $\sim 1\%$ of KW short GRBs in terms of the observed $E_{\text{peak},z}$ and is within the 1σ prediction band of both the ‘Amati’ and ‘Yonetoku’ relations based on 315 long/soft (Type II) GRBs with known z^{99-101} . Furthermore, in the $E_{\text{iso}}-E_{\text{peak},z}$ plane, GRB 200826A is inconsistent with the short-hard (Type I) GRB population (see Extended Data Figure 1). Thus, despite the short duration, the KW parameters of GRB 200826A imply that it belongs to the long/soft GRB population.

In addition to the IPN detections, the Cadmium Zinc Telluride Imager (CZTI 102) on board AstroSat also detected the burst. We reanalysed the data, combining data from all four quadrants to create 20–200 keV light curves with 0.05 s, 0.1 s and 0.2 s bins. All light curves show a single pulse, with hints of sub-structure in during the rise in the smaller time bins (see Extended Data Figure 1). We process the 0.05 s light curves with the CIFT pipeline 103 , which incorporates better data analysis as compared to the quick-look pipeline and produces more robust results. Our reanalysis yields a peak time of UT 04:29:52.95 - consistent with the Fermi peak. The new value of T_{90} is $0.94^{+0.72}_{-0.18}$ s, significantly shorter than the quick-look values reported in Ref. 84 , but consistent with Fermi and Konus-Wind.

Optical. We ingested the GBM localization map into the GROWTH Target-of-Opportunity (ToO) marshal, an interactive tool design to plan and schedule ToO observations for ZTF 73 . The observation plan generated by the ToO marshal relies on `gwemopt` $^{104-106}$, a code that optimizes the telescope scheduling process for skymaps with a healpix format, like the *Fermi*-GBM maps. The `gwemopt` procedure involves slicing the skymap into predefined tiles of the size and shape of the ZTF

field-of-view, determining which fields have the highest enclosed probability, and optimizing observations based on airmass and visibility windows. For this purpose, we used a modified version of the “greedy” algorithm described in Ref. ¹⁰⁷, implemented within `gwemopt`. The resulting optimized plan for the optical follow-up of GRB 200826A consisted of four primary ZTF fields and one secondary field^b. The fields were observed in the *r*- and *g*-band for 300 s each, starting 4.9 hours after the GBM detection. The observing plan for the first night covered 186 deg², corresponding to 77% of the GBM region (see Fig.2). Once the *Konus-Wind* data became available on ground at about 17:55 UT, the 288 arcmin² IPN error box was derived using *Konus*, GBM, and HEND data, which allowed the X-ray Telescope (XRT¹⁰⁸) on board *Swift* to initiate a ToO at about 20:45 UT. The IPN box was published later at about 21:30 UT¹⁰⁹.

Reference images of the fields are then subtracted from the ZTF ToO observations and any high significance difference ($>5\sigma$) generates an *alert*^{72,110} that contains relevant information about the transient. We queried the stream of alerts using three different tools: the GROWTH marshal¹¹¹, the Kowalski infrastructure^{112c} and AMPEL^{113–116d}. Our filtering scheme has been described in previous SGRB and GW searches^{73,76}, but we summarize the main points here. We aim to identify sources that 1) are spatially coincident with the skymap, 2) are detected only after the GBM trigger, 3) are far from known bright sources, 4) are spatially distinct from Pan-STARRS 1 (PS1¹³) stars (based on Ref. ¹¹⁸), 5) have at least two detections separated by at least 15 minutes to avoid moving objects, 6) have a real-bogus score (RB^{112,119}) greater than 0.15, and 7) showed an increase in their flux relative to the reference image. For each candidates passing these filters, we visually inspected the light curve, cross-matched with the PS1-DR2 catalog to check for previous activity, and cross-matched against the Wide-field Infrared Survey Explorer (WISE¹²⁰) catalog to determine whether it was consistent with an active galactic nucleus (AGN) based on its position in the WISE color space^{121–123}. Additionally, we searched the Minor Planet Center (MPC) to ensure that our potential counterparts were not consistent with known solar system objects. Forced photometry¹²⁴ was performed using data from the 10 nights previous to the trigger, to remove young SNe from the sample. A final quality check was done by querying for alerts around each candidate, and rejecting transients that have multiple alerts within a radius of 15” as these alerts could suggest artifacts, ghosts, or slow-moving objects near the stationary points¹²⁵.

^bThe ZTF secondary fields are strategically located to cover the chip-gaps of the fields in the primary grid

^c<https://github.com/dmitryduev/kowalski>

^d<https://github.com/AmpelProject>

After the first night of observations 28195 alerts were generated in the region and 14 sources passed every stage of our filtering criteria¹²⁶. The SGRB afterglow, ZTF20abwysqy¹²⁷, was not in our first selection as there was a ZTF alert 11'' NE; however, it passed all the other filtering criteria. Two of the ZTF ToO fields triggered during the first night covered $> 99\%$ of the triangulated IPN region. However, none of the candidates reported in Ref. ¹²⁶ fell there. For our second night of observations, we scheduled 600 s observations in r - and g -band and queried the database.

Probability of Chance Coincidence. We roughly calculate the probability of chance coincidence (p) of this optical transient to be independent of GRB 200826A. For this we follow Ref. ¹²⁸ and use Poissonian statistics to derive the probability for one or more events to be randomly coincident:

$$p = 1 - e^{-\lambda}$$

where λ is the product of three different values, $\lambda = \prod_{i=1}^3 r_i$. For this study, we consider r_1 to be the time window between the last ZTF non-detection (0.72 days before the trigger time) and the time of the detection (0.21 days, see Extended Data Table 1), resulting in $r_1 = 0.93$ days. The parameter r_2 is the rate of *Fermi* GRBs, which from the latest GBM catalog⁶² gives ~ 0.65 bursts per day. Finally r_3 is the ratio between the total IPN area (288 arcmin²) and the sky. This derivation gives a $p = 1.16e - 6$, and allows us to rule out a random association between the GRB and the afterglow at the 4.87σ level.

2 Follow-up

The summary of the follow-up results can be found in Extended Data Table 1, Extended Data Table 3, and Extended Data Table 4.

Optical/NIR. Here we present our optical and near infrared (NIR) follow-up results. In addition to our observations, MASTER¹²⁹ and Kitab follow-up with a clear filter^{130,131} led to upper limits of 18.3 mag and 20.4 mag respectively.

Las Cumbres Observatory. We performed follow-up of ZTF20abwysqy with the Spectral camera mounted at the 2-m Faulkes Telescope North (FTN), located at Haleakala Observatory. Starting on 2020-08-28 11:32:26 UT we acquired three sets of images with 300s exposures each in g - and r -bands through the LCO observation portal^e. Furthermore, we obtained a second

^e<https://observe.lco.global/>

(reference) epoch in g - and r -bands on 2020-09-02 UT. Our images were reduced by an automatic subtraction pipeline. Our pipeline retrieves images from LCO, stacks them, extracts sources from the image using `Source Extractor`¹³² and performs photometric calibration of sources using the PS1 catalog. Then, the pipeline performs image subtraction using the High Order Transform of Psf ANd Template Subtraction code (HOTPANTS¹³³) to subtract a PSF-scaled reference image aligned using `SCAMP`¹³⁴. Our photometry measurements reported in g - and r -band were determined after host subtraction using the LCO reference images in the same filters acquired on 2020-09-05.

Hale Telescope. We obtained two epochs of dithered J-band imaging with the Wide Field Infrared Camera (WIRC¹³⁵) on the Palomar Hale 200-in (P200) telescope on 2020-08-28 UT and on 2020-09-04 UT, spending an hour integrating on target during each epoch. We reduced the images using the image reduction pipeline described in Ref. ¹³⁶. Images were aligned and stacked using `SWarp`¹³⁷ and calibrated against the 2MASS catalog. Image subtraction between the two epochs was performed using the method described in Ref. ¹³⁸ to derive flux measurements and its uncertainty for the first epoch.

Gran Telescopio Canarias. The 10.4m Gran Telescopio Canarias (GTC) (+OSIRIS) obtained spectroscopic observations starting on August 30, 04:30 UT. Two 1200s exposures were gathered with the R1000B grism and one 1200s exposure was obtained with the R2500I grism, in order to cover the entire 3700–10000 Å range. The slit was placed covering the position of the potential host galaxy. Standard routines from the Image Reduction and Analysis Facility (IRAF) were used to reduce the data.

Lowell Discovery Telescope. We used the Large Monolithic Imager (LMI¹³⁹) mounted on the 4.3m Lowell Discovery Telescope (LDT) to observe the optical transient ZTF20abwysqy on three different nights: August 29, September 13, and September 19 (3.2, 18.1 and 24.2 days after the GRB trigger). Observations were conducted with an average airmass of 1.0 while the seeing varied from 1.1'' for the first night to 1.4'' and 1.6'' for the second and third night, respectively. We took 8 exposures of 180 s in the r -band on August 29. Images were taken with different filters on September 13: 5 exposures of 180 s in u -band, 4 exposures of 180 s in g -band, 6 exposures of 180 s in i -band and 6 exposures of 180 s in z -band. Finally, 10 exposure of 150 s in the i -band and r -band were taken during the last night of observations on September 19. We used standard procedures to perform bias and flat-field correction. The astrometry was calibrated against the SDSS catalog (release DR16¹⁴⁰) and frames were aligned using `SCAMP` and stacked with `SWarp`.

After stacking the images, we extracted sources using `Source Extractor` and the magnitudes were calibrated against 45 PS1 stars in the field on average. We used HOTPANTS to perform image subtraction between the first and third epochs and found a source in the r -band at the location of ZTF20abwysky with a magnitude of 24.46 ± 0.12 mag (see Extended Data Table 1). We determine the magnitudes of the host galaxy using the second epoch of observations. To verify this result, we used the 80 days i -band GMOS-North (see description in the paragraph below) as a reference and the HOTPANTS subtraction shows a source with r -band magnitude of 24.76 ± 0.23 mag, consistent with the result using the LDT reference.

Gemini Observatory. We acquired images of the transient location on September 23, October 10, and November 7 corresponding to 28.28, 46.15 and 80.23 days after the GRB trigger respectively, which we denote as epoch 1, 2, and 3. We used GMOS-North, mounted on the Gemini North 8-meter telescope on Mauna Kea, under the approved Director’s Discretionary Time (DDT) proposal DD-104 (P.I.: L. Singer). The host galaxy coordinates, in the Gemini images, are $\alpha = 00^h 27^m 08.5557^s$, $\delta = +34^d 01^m 38.634^s$.

The first set of observations (epoch 1) was scheduled to be closest in time to a SN1998bw-like peak without suffering from moon illumination. All three sets of observations consisted of 14 200 s r - and i -band exposures, with a position angle of 45° to avoid blooming from neighboring bright stars. The average airmass for the observations was ~ 1.0 and the seeing was stable throughout the three epochs, at $0.7''$, $0.5''$, and $0.85''$. The images were later reduced using DRAGONS^{f141}, a Python-based data reduction platform provided by the Gemini Observatory.

We extracted sources using `Source Extractor` and determined a photometric zero point using 23 PS1 stars in the field. Using the ZOGY¹⁴² algorithm-based python pipeline, we performed image differencing on Gemini data. The pipeline makes use of `Source Extractor`, `PSFEX`¹⁴³, `SWarp`, `SCAMP`, and `PyZOGY`¹⁴⁴ to perform image subtraction. We took a $1300 \text{ pix} \times 1300 \text{ pix}$ sized cutout for both the science and reference images, centered at the ZTF20abzwysqy position in the image in order to achieve good subtraction quality, as the images were affected by background variation. Sources were extracted using `Source Extractor` from both science and reference images with a $5\text{-}\sigma$ detection threshold. The resulting catalogues were fed to `SCAMP`, which calculates and corrects for the astrometric errors in both science and reference catalogues with the help of a Gaia data release 2 catalogue¹⁴⁵ for the same field. Using `SWarp`, images are

^f<https://dragons.readthedocs.io/>

then re-sampled to subtract the background and the point-spread function (PSF) of the images. `Source Extractor`-generated weight maps were used as input to the `SWarp` to generate variances, which when added with the images' Poisson noise in quadrature, which results in the rms image. We select good sources from `Source Extractor` catalogues based on their S/N, full width at half-maximum (FWHM) and `Source Extractor` flag values. Sources which raised a `Source Extractor` flag were discarded. These sources were used to match the flux level of the images. We used `PSFEx` to extract the PSF model of the re-sampled images with the `Source Extractor` catalogue as input to `PSFEx`. The rms images, re-sampled images, PSF models and astrometric uncertainty are used as input to `PyZOGY`, which generate the difference image and a corrected score image as final products¹⁴². A source is detected with a ZOGY corrected score of 4 in the *i*-band, at the location of the transient. With ZOGY, we derive an *i*-band magnitude of 25.49 ± 0.12 mag for epoch 1. No source is detected in the *r*-band of epoch 1 nor in any band in epoch 2.

We have confirmed the ZOGY-based results using an independent image subtraction pipeline, `FPipe`¹⁴⁶, which is based on empirically measuring the PSFs of the science and reference images and matching them using the common PSF method (CPM¹⁴⁷). We detect the source with an *i*-band brightness of 25.45 ± 15 mag in epoch 1. The results agree within uncertainties with the ZOGY-based subtractions. We do not detect the source in the *r*-band up to a $5\text{-}\sigma$ limit of $r > 25.6$ mag for epoch 1. No source is detected in either filter during epoch 2 up to a 5σ upper limit of $i > 25.4$ mag and $r > 25.5$ mag¹⁴⁸ (see Figure 3).

Therefore, using two independent image subtraction pipelines, we confirm the detection of a source in the *i*-band, at the location of the transient.

UV/X-ray. *Swift* began observing the IPN localization region of GRB 200826A 0.7 days after the trigger. The last two observations were ~ 4 ks and triggered as ToO observations. The XRT data was reduced by the the online reduction pipeline^{g149,150}. As the hardness ratio remains constant within error bars, we assume a single absorbed power law spectrum, with a Galactic neutral hydrogen column¹⁵¹ of to $6.02 \times 10^{20} \text{cm}^{-2}$. We convert the count rates (see Extended Data Table 3) to flux density at an energy of 1 keV, using the parameters derived by the *Swift* pipeline: a photon index of $\Gamma_X = 1.5^{+0.7}_{-0.5}$, an intrinsic host absorption of $n_{H,int} = 6^{+32}_{-6} \times 10^{20} \text{cm}^{-2}$, and an unabsorbed counts-to-flux conversion factor of $4.27 \times 10^{-11} \text{ erg cm}^{-2} \text{ ct}^{-1}$.

^ghttps://www.swift.ac.uk/xrt_curves/00021028/

The first UV/Optical Telescope (UVOT) observations of ZTF20abwysqy was ~ 1.6 days after the burst and the detection in the white filter was associated to the underlying galaxy¹¹.

Radio. For the afterglow modeling, we additionally use data reported by Ref. ¹². They measure a flux density of $\sim 40 \mu\text{Jy}$ at a mean frequency of 6 GHz using the Karl G. Jansky Very Large Array (VLA), 2.28 days after the trigger. We observed GRB 200826A in band-5 (1050–1450 MHz) of the upgraded Giant Metrewave Radio Telescope (uGMRT¹⁵⁴) on 2020-09-09 15:41:30.5 UT and 2020-09-14 23:38:25.5 UT (~ 14.5 and ~ 19.8 days after the burst respectively) under the approved DDT proposal ddtC147 (P.I.: Poonam Chandra). The data was recorded with the GMRT Wideband Backend (GWB) correlator with a bandwidth of 400 MHz divided in 2048 channels, with the central frequency of 1250 MHz. The total on-source time was ~ 75 mins with overheads of ~ 30 mins. 3C48 was used as the flux and bandpass calibrator and J0029+349 was used as the phase calibrator. We used standard data reduction procedures in Common Astronomy Software Applications (CASA)¹⁵⁵ for analysing the data. The dead antennas were first flagged by manual inspection and the end channels were flagged due to low gain. The automatic flagging algorithms incorporated into the CASA task *flagdata* were used to remove most of the Radio Frequency Interference (RFI). Any remaining corrupted data were then flagged manually. The calibrated data was then imaged and self-calibrated to get the final image. The synthesized beam for the final image was $\sim 4'' \times 2''$. We found no evidence of radio emission at the GRB position on both days. The 3σ upper limits were $48.6 \mu\text{Jy}/\text{beam}$ for the observations taken on 2020-09-09¹⁵⁶ and $57.4 \mu\text{Jy}/\text{beam}$ for the observations taken on 2020-09-14.

3 The Host Galaxy

Our GTC spectra of the host galaxy show strong [OII] and [OIII] features at $z = 0.748$, in agreement with a previous report based on data from the Large Binocular Telescope Observatory (LBTO)⁹⁸.

We modeled the GTC spectrum with the Penalized Pixel-Fitting (pPXF)¹⁵⁷ to infer the properties of the host galaxy’s stellar populations, resulting in a fit with $\chi^2_\nu = 1.011$. The pPXF results confirmed the [OII] features at 3726Å and [OIII] lines at 4959Å and 5007Å, indicators of recent star-formation and young, hot stars. The stellar population age derived from the weighted ages of the templates is $1.514^{+3.83}_{-0.85}$ Gyr, with evidence for a younger ~ 0.1 Gyr population (see Figure Extended Data Figure 2 and Extended Data Table 6).

Additionally, we corrected the magnitudes of the host galaxy (see Extended Data Table 2) using foreground extinction maps¹⁵⁸ and fed them to `Prospector`¹⁵⁹, to model the SED of the host galaxy. `Prospector` uses `fspy`^{160,161} to generate multiple stellar populations, and fits the observed photometry to determine the formed mass of the galaxy, age, and intrinsic extinction among other parameters, using the WMAP9 cosmology¹⁶² internally. We fitted the photometry to a galaxy using the Chabrier¹⁶³ initial mass function (IMF), the Calzetti¹⁶⁴ extinction curves for the dust around old stars, and a star formation history (SFH) with the form of $te^{-t/\tau}$. Our results using the nested sampling `dynesty`¹⁶⁵ algorithm, gave a galactic stellar mass distribution of $M_{gal} = 4.64 \pm 1.67 \times 10^9 M_{\odot}$, a mass-weighted galactic age of $t_{gal} = 1.08^{+1.28}_{-0.72}$ Gyr, a metallicity of $\log(Z/Z_{\odot}) = -1.06^{+0.67}_{-0.48}$, a dust extinction of $A_V = 0.34^{+0.33}_{-0.22}$, and a star formation rate (SFR) of $4.01^{+41.87}_{-3.59} M_{\odot} \text{ yr}^{-1}$ (see the resulting SED in Extended Data Figure 2 and the posterior probability distributions for the free parameters in Extended Data Figure 3). We derived the SFR and weighted-mass age following similar studies on GRB host galaxies^{166–168}, and the stellar mass from the mass fraction derived with `Prospector`.

4 Modeling

For the modeling of the multiwavelength emission, all optical and NIR observations were from difference images and were corrected for foreground extinction¹⁵⁸.

The afterglow. In the standard synchrotron fireball model a power-law energy distribution characterized by the index p , $N(E) \propto E^{-p}$, results in a SED described by a series of broken power laws^{17,170}. The frequencies at which the broadband SED presents its breaks are the self-absorption frequency ν_a , the synchrotron frequency ν_m , and the cooling frequency ν_c .

The temporal decline of both the optical and X-ray data during the first four days can be described by a single power law model which suggests the jet-break has not yet occurred. Our ZTF g -band observations are the most constraining in the optical therefore we use these to estimate the temporal decline rate $\alpha_o = -1.05 \pm 0.13$. We use the 1 keV XRT data to find $\alpha_x = -0.89 \pm 0.07$. The similar slopes between the optical and X-ray observations suggest the location of the cooling break frequency, ν_c , lies beyond X-ray frequencies. Therefore, we estimate the spectral index at ~ 1 day between the X-ray and optical as $\beta_{ox} = -0.67 \pm 0.02$.

We now use α_o and α_x to estimate the power law index of the electron energy distribution, p , and to determine the circumburst density profile. In a constant density (ISM-like) medium

$\alpha_{\text{ISM}} = \frac{3(1-p)}{4}$ ($\nu < \nu_c$), which gives $p_o = 2.44 \pm 0.17$ and $p_x = 2.17 \pm 0.10$ ¹⁷¹. In the wind-like scenario $\alpha_{\text{wind}} = \frac{(1-3p)}{4}$, which gives $p_o = 1.78 \pm 0.17$ and $p_x = 1.50 \pm 0.10$. The optical to X-ray spectral index β_{ox} gives an estimate for p of 2.34 ± 0.04 ($\beta = \frac{1-p}{2}$). Theoretical studies of relativistic collisionless shocks predict $p \gtrsim 2$ and particularly $p \sim 2.2$ in the ultra-relativistic limit¹⁷². Given the low values of p for the wind-like scenario we choose to assume an ISM-like density profile and $p \sim 2.4$ throughout this work.

We note the possibility of a wind-like environment with $p < 2$ and a spectral break ν_c between the optical and X-rays. Such a scenario would not change the predicted optical and NIR emission and it would be unable to account for the late GMOS *i*-band detection without an additional component. However, this scenario may allow for a better fit of the afterglow X-ray evolution, which would decay at a rate shallower than the optical by $\Delta\alpha = 0.25$. We do not consider this scenario further because of the theoretically disfavored value of p .

Bayesian afterglow modeling. We now use Bayesian inference to analyze the X-ray, optical, NIR, and radio counterpart. We use two independent pipelines, one using Markov chain Monte Carlo (MCMC) based on the EMCEE Python package¹⁷³ and one using nested sampling based on PyMultinest¹⁷⁴. The pipelines use different priors and implementations, but arrive at consistent results. Both utilize the `afterglowpy` python package¹⁷⁵ to estimate the physical parameters of the multi-wavelength afterglow. `Afterglowpy` is a public, open-source computational tool which models forward shock synchrotron emission from relativistic blast waves as a function of jet structure and viewing angle. Descriptions of the MCMC implementation may be found in Refs.^{176,177} and the nested sampling implementation in Ref.²⁰.

For this work we assume Gaussian statistics for the optical and radio data, while assuming Poissonian statistics (via the C-statistic¹⁷⁹) for the X-ray data due to low detector counts. The `Afterglowpy` model is parametrized by the isotropic kinetic energy, $E_{\text{K,iso}}$; jet collimation angle, θ_c ; viewing angle, θ_v ; the circumburst constant density, n ; the spectral slope of the electron distribution, p ; the fraction of energy imparted to both the electrons, ϵ_e , and to the magnetic field, ϵ_B , by the shock. The redshift and luminosity distance of the source are held fixed.

Our modelling of the host indicates a small galaxy of stellar mass $\sim 5 \times 10^9 M_\odot$ with moderate extinction $A_V = 0.34_{-0.22}^{+0.33}$ (see §3 for details). In the MCMC implementation, we incorporate a Small Magellanic Cloud-like host extinction correction with total-to-selective extinction $R_V = 2.93$ ¹⁸⁰ implemented with the `dust-extinction` software package¹⁸¹. We leave the

color excess $E(B - V)$ as a free parameter, with a prior distribution computed from the A_V posterior found from modelling the host galaxy. This posterior is similar to the distribution of host extinction values observed in LGRBs with well-sampled multiband photometry^{182–184}. The nested sampling implementation performs no extinction correction.

Using a top hat model for the jet structure, we perform a search over the parameter space by allowing all the parameters $E_{K,\text{iso}}$, θ_c , θ_v , n , p , ϵ_e , ϵ_B , and $E(B - V)$ to vary with broad priors. We report only the MCMC results, although the nested sampling results are consistent. Given both the highly degenerate nature of the afterglow fitting and the low number of observations available, some of the parameters are not particularly constrained. One exception is the spectral slope of the electron distribution, for which we find $p = 2.4 \pm 0.04$, consistent with the analytical results. The uncertainty in the circumburst density, $n = 5.5^{+187.3}_{-5.4} \times 10^{-2} \text{ cm}^{-3}$, includes typical ranges for both SGRBs and LGRBs. We find very little reddening from the host, $E(B - V) = 2.5^{+4.8}_{-2.3} \times 10^{-2}$. The posterior probability distributions are shown in Extended Data Figure 5 and the parameter estimates are listed in Extended Data Table 7.

Bayesian model selection. In addition to an afterglow-only model, we consider an afterglow plus a KN and an afterglow plus a SN.

Kilonova. We use SEDs simulated by the multi-dimensional Monte Carlo radiative transfer code POSSIS¹⁸⁵. The simulations are performed over a grid of KN parameters: dynamical ejecta $M_{\text{ej}}^{\text{dyn}}$, disk wind ejecta $M_{\text{ej}}^{\text{wind}}$, opening angle Φ , and the observation angle Θ_{obs} (see Ref. ²⁰ for details). We use Gaussian process regression^{186,187} to interpolate the model, enabling rapid parameter inference.

Supernova. The SN model starts with a bicubic spline in time and frequency, $L_{\nu}^{\text{SN}}(t, \nu)$, that interpolates a K-corrected SN1998bw template²². We apply to the template a scale factor, k , and a stretch factor, s , which are drawn from a bivariate normal distribution that is consistent with a historical sample of GRB-SNe^{23,190}. The model for the observed flux density is

$$F_{\nu}^{\text{SN}}(t_{\text{obs}}, \nu_{\text{obs}}) = \frac{(1+z)k}{4\pi d_L^2} L_{\nu}^{\text{SN}}\left(\frac{t_{\text{obs}}}{(1+z)s}, (1+z)\nu_{\text{obs}}\right).$$

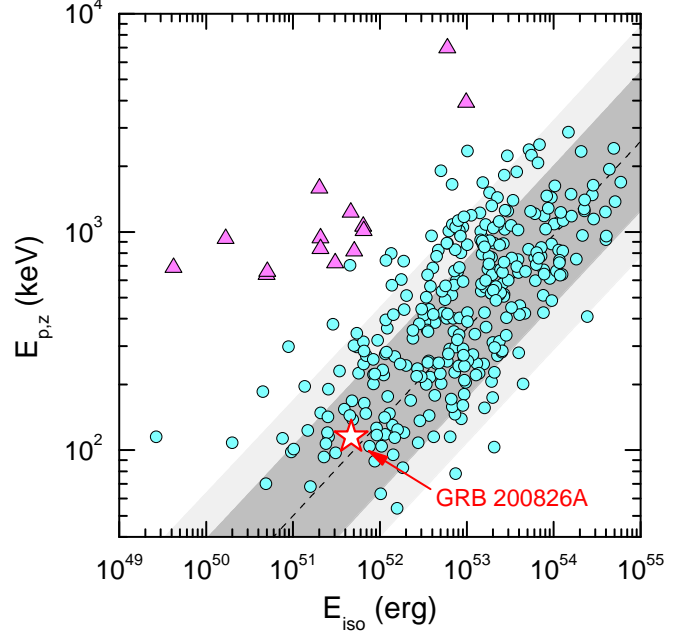
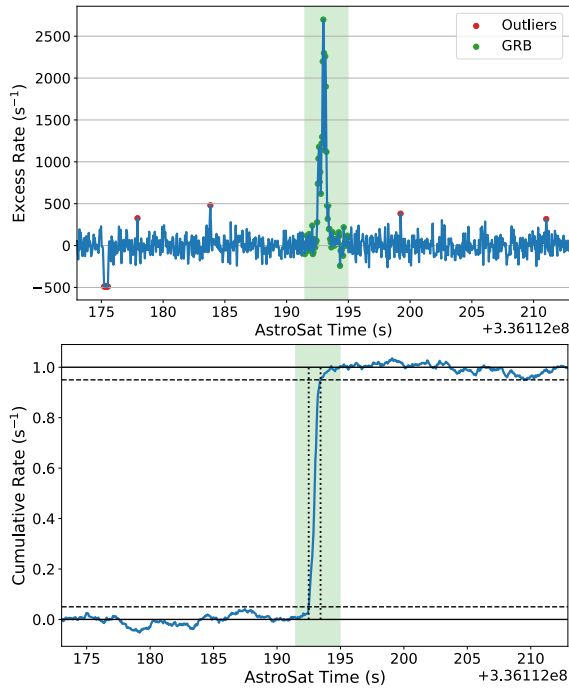
We perform Bayesian model selection to determine which model best explains the data. We used the nested sampling pipeline to calculate the Bayesian evidence for each of the three models. The Bayes factor, or the ratio of the evidences, between the afterglow-plus-KN model and the

afterglow-only model is ~ 1 , indicating that neither model is strongly favored over the other, because the KN contributes negligible flux compared to the afterglow at the time of the GMOS observation. The Bayes factor between the afterglow-plus-SN model and the afterglow-only model is $\sim 10^{5.5}$, strongly favoring the presence of a SN.

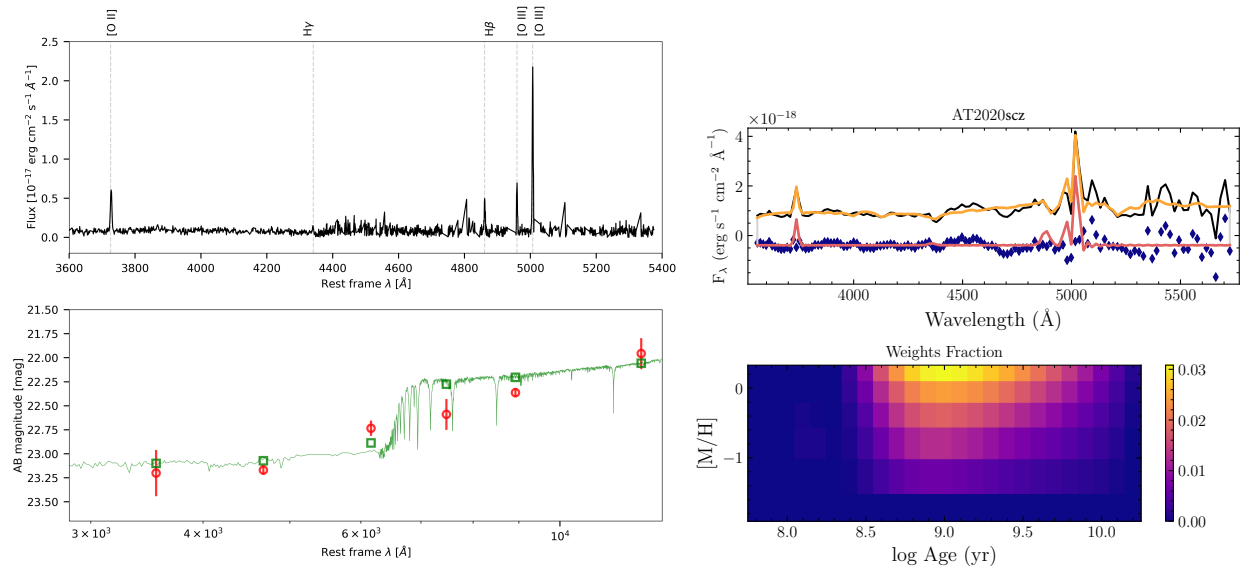
To better understand the source of discriminating power between models, we carry out a posterior predictive check. Here, we sample the posterior while excluding the GMOS i -band data point and then predict its value using the rest of the data. The posterior predictive distribution of the AB magnitude for the GMOS i -band detection both with and without the inclusion of a SN contribution is shown in Extended Data Figure 6. The fit with the SN is consistent with the observation, while the fit without the SN is inconsistent at the $\sim 5\sigma$ level. Therefore, inclusion of the GMOS i -band data point requires a SN component, confirming a collapsar origin.

We compared our GMOS-N detection against extinction-corrected i -band fluxes of three GRB-KN candidates found in the literature¹⁹¹; using photometry from GRB 130603¹⁹² and GRB 160821^{193,194}, and the compiled light curve of AT2017gfo^{5,196}, we correct to the redshift of GRB 200826A. We fit each light curve to the best-fit 2D model of GW170817²⁰ from POSSIS using SNCosmo¹⁹⁷ and extract the corresponding i -band magnitude at a rest-frame time of 16 days. At $z = 0.748$, both AT2017gfo and GRB 160821 would be at $M \approx -10$ mag, ~ 8 mags fainter than our detection. GRB 130603B does not have enough late-time detections at the same phase for comparison.

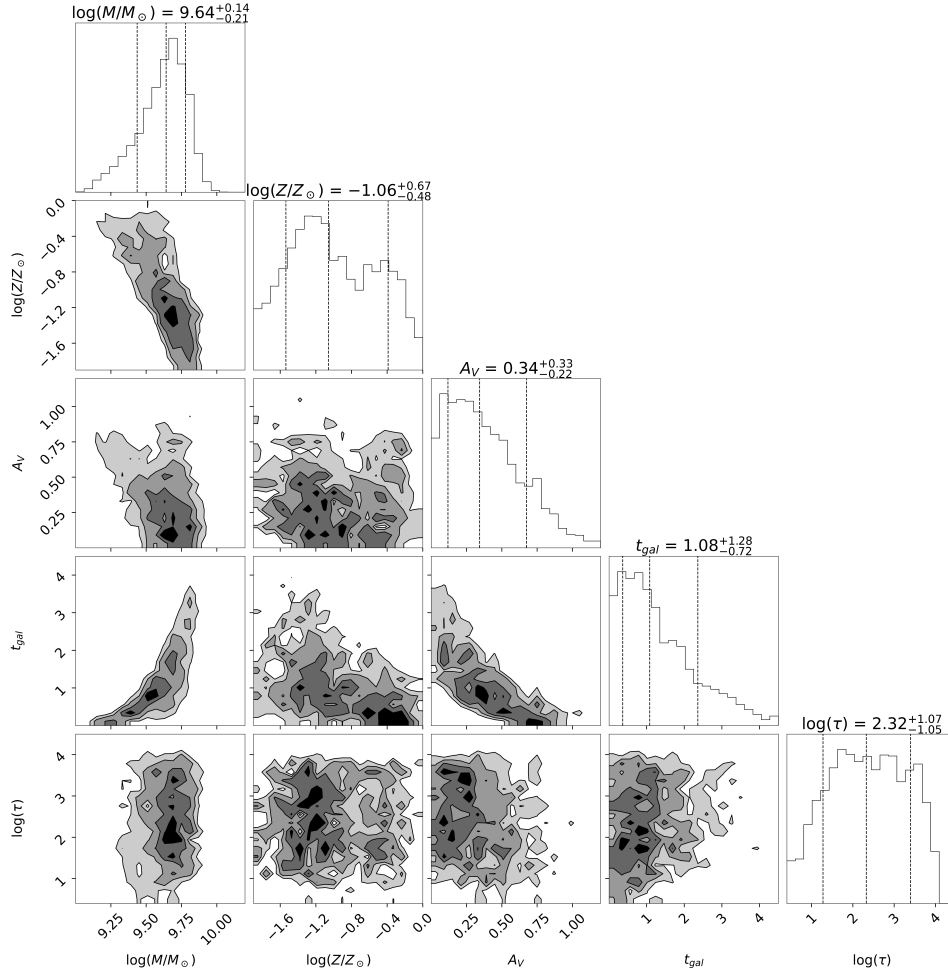
Our detection of a source with an extinction-corrected absolute magnitude in the i -band with $M_i = -18.0$ is consistent with the population of collapsars associated with LGRBs. A typical SN Ic reaches its peak magnitude 10 to 20 days post-burst, at $M_B = -17.66 \pm 1.18$ mag¹⁹⁸. Fig. 3 shows K-corrected light curves for three well-sampled GRB-SNe: SN1998bw, one of the brightest; SN2006aj and SN2010bh, among the faintest.



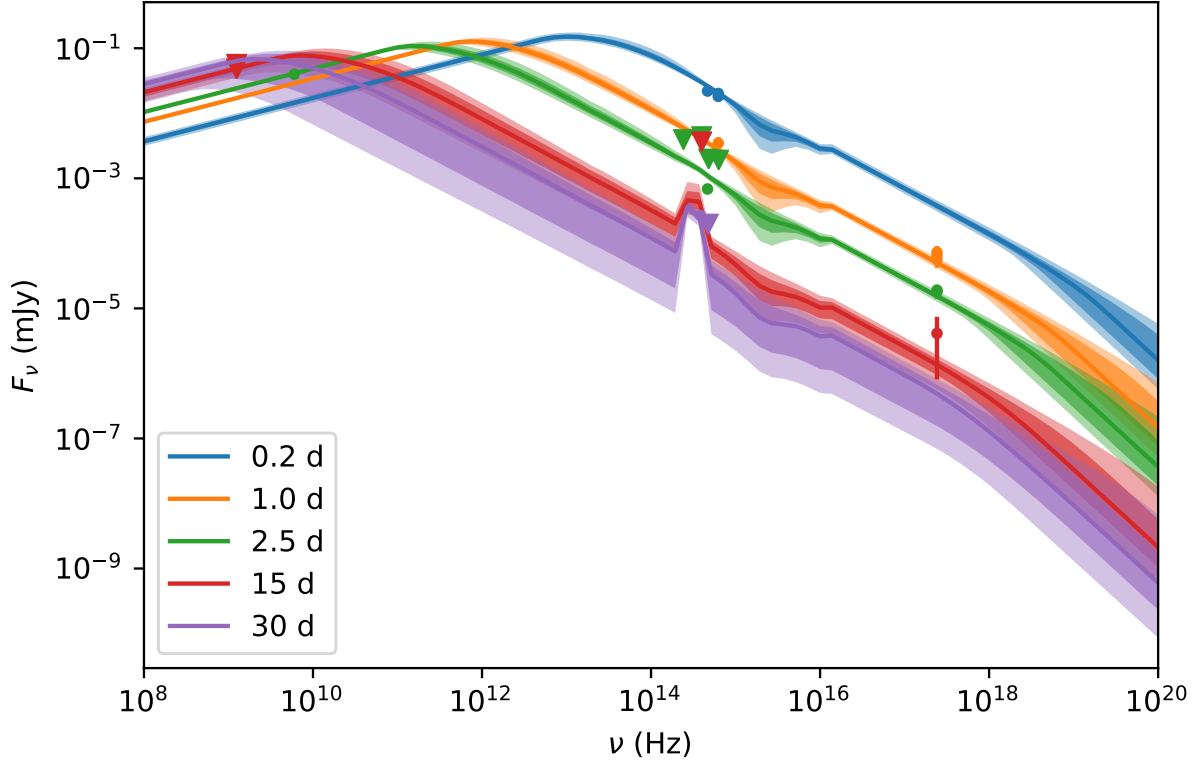
Extended Data Figure 1: **The AstroSat and Konus-Wind gamma-ray detections.** (*left*) Upper panel: The de-trended light curve for GRB 200826A obtained from AstroSat CZTI data. We combined data from all four CZTI quadrants and binned it in 0.05 s bins. We fit and subtract a quadratic trend from the background to obtain zero-mean data. The shaded green region and corresponding green symbols denote a conservative GRB time span excluded from background trend estimation. Similarly red points denote outliers that are automatically flagged and rejected from the background estimate. Lower panel: A cumulative light curve obtained by summing the de-trended data, and normalised such that the median post-GRB value is 1.0. The dashed horizontal lines denote the 5% and 95% intensity levels. The corresponding vertical dotted black lines denote T_{05} and T_{95} , yielding T_{90} of $0.96^{+0.71}_{-0.17}$ s. (*right*) Rest-frame energetics of 331 Konus-Wind GRBs (SGRB: triangles, LGRB: circles) with known redshift in the $E_{\text{iso}}-E_{\text{peak,z}}$ plane, with $E_{\text{peak,z}}$ the rest frame E_{peak} . The hardness-intensity (‘Amati’) relation for LGRBs is plotted with its 68% and 90% prediction intervals (dark and light gray regions, respectively). GRB 200826A, as a red star, appears not to be consistent with the SGRB population.



Extended Data Figure 2: **The Host Galaxy.** (*left*) In the upper panel, we show the GTC spectrum of the host galaxy and the lines used to determine a redshift of 0.748. In the bottom panel, the photometry of the host galaxy (*ugrizJ*, see Extended Data Table 2) in the AB system is presented in red circles. The SED model and photometry from *Prospector* are shown in green. (*right*) The pPXF host galaxy model results described in §3. (top) The integrated spectrum (black) overlaid with the best-fit spectrum (orange), which sums the contributions of stars and gas in the modeled galaxy. The red spectrum shows the gas contribution to the spectrum, and the blue diamonds show the residuals to the fit. The gas is offset by $1.59\text{e-}18 \text{ erg s}^{-1} \text{ cm}^{-2}$. (bottom) The pPXF weights (color bar) of the different stellar population templates used to construct the best-fit galaxy.



Extended Data Figure 3: **Posterior distribution of the Prospector parameters.** The covariances and posterior probability distributions of the parameters for the host galaxy derived from the Prospector modelling described in §3.



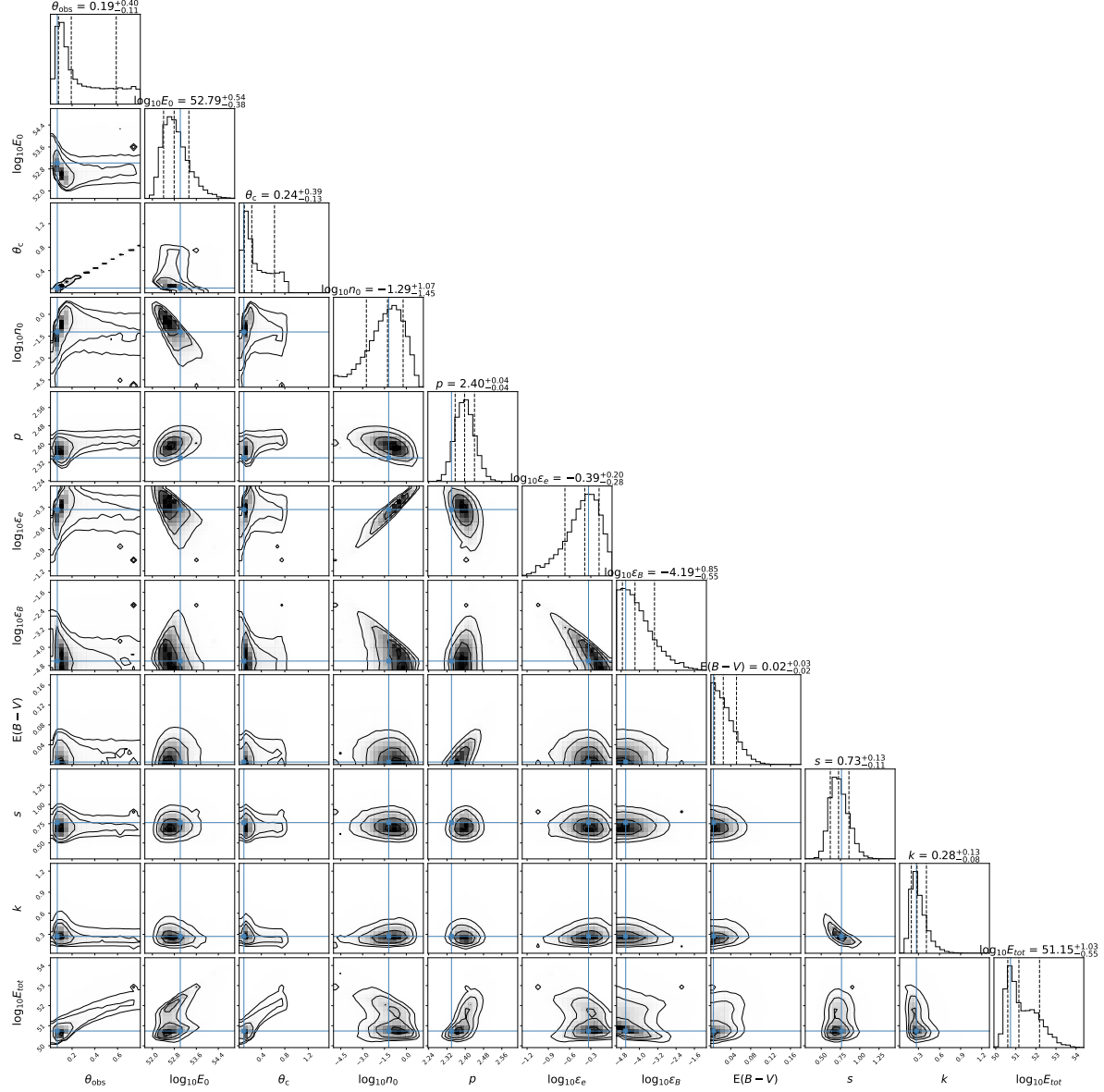
Extended Data Figure 4: **SED sequence of the afterglow.** The SED of our model compared to observations at five epochs. The cooling frequency ν_c is located at frequencies higher than 1 keV (i.e. $\nu_c > 2.4 \times 10^{17}$ Hz). The glitches at optical $\nu \sim 10^{16}$ Hz are the edge of validity of our dust extinction model. The SN makes a large contribution at late times. See the observations in Extended Data Table 1, Extended Data Table 3, and Extended Data Table 4.

5 Data Availability

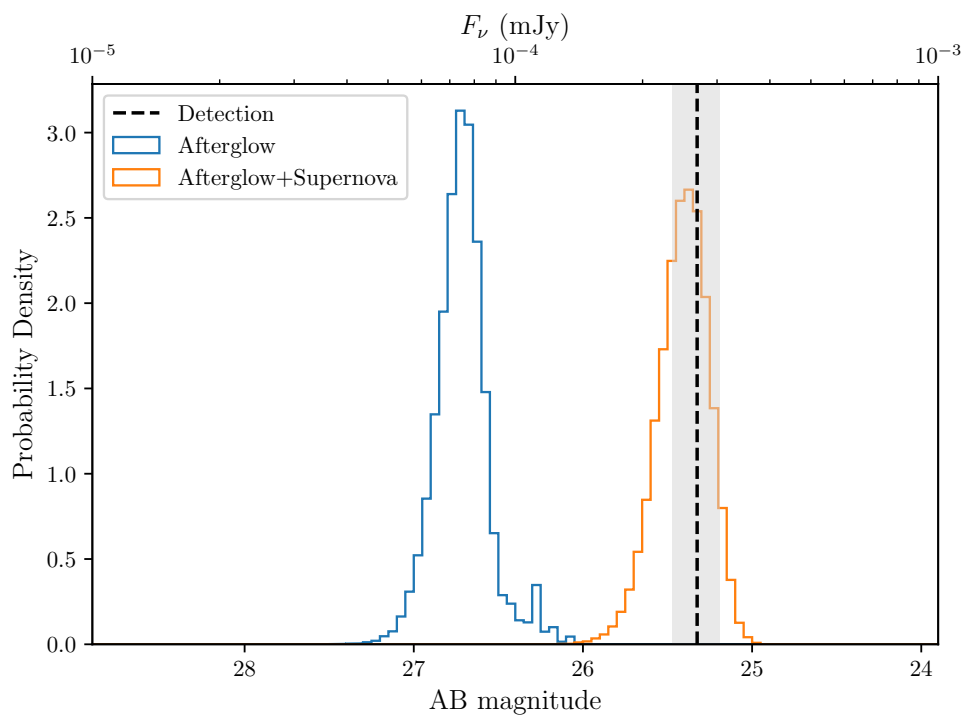
Upon request, the corresponding author will provide data required to reproduce the figures, including light curves and spectra for any objects.

6 Code Availability

Upon request, the corresponding author will provide code (primarily in python) used to produce the figures.



Extended Data Figure 5: **Posterior distribution for the afterglow fit.** The covariances and posterior probability distributions of the parameters for the top hat jet afterglow with 1998bw-like SN model described in §4. The afterglow parameters are the viewing angle θ_{obs} (rad), on-axis isotropic kinetic energy E_0 (erg), opening angle θ_c (rad), circumburst number density n_0 (cm^{-3}), electron spectral index p , fraction of energy in accelerated electrons ϵ_e , and fraction of energy in magnetic field ϵ_B . We assume an SMC-like extinction curve with variable $E(B - V)$. The SN model is a 1998bw template with variable stretch s and scale k . The total beaming-corrected kinetic energy in the jet E_{tot} (erg), computed from the afterglow parameters, is also reported. The histograms denote the 14, 50, and 84 percentiles of the distributions, with blue lines marking the solution with maximum posterior probability density.



Extended Data Figure 6: **Posterior predictive plot for the GMOS *i*-band detection.** The posterior of the AB magnitude estimated at the time of the *i*-band data point (~ 28 days after the trigger) using afterglow only (blue) and afterglow-and-SN (orange) light curves are shown.

Julian day	δt	Instrument	Filter	AB magnitude	σ_{mag}	5 σ Limiting magnitude	A_ν
2459085.8814	-1.81	P48+ZTF	g	—	—	21.50	0.21
2459085.9678	-1.72	P48+ZTF	r	—	—	21.46	0.15
2459086.9011	-0.79	P48+ZTF	r	—	—	21.33	0.15
2459086.9696	-0.72	P48+ZTF	g	—	—	21.31	0.21
2459087.8986	0.21	P48+ZTF	g	20.86	0.04	22.50	0.21
2459087.9181	0.23	P48+ZTF	r	20.69	0.05	22.27	0.15
2459087.9651	0.28	P48+ZTF	g	20.95	0.16	21.23	0.21
2459088.8340	1.15	P48+ZTF	g	22.75	0.26	22.53	0.21
2459088.9038	1.22	P48+ZTF	r	—	—	21.30	0.15
2459088.9761	1.29	P48+ZTF	g	—	—	21.2	0.21
2459089.6256	1.93	P200+WIRC	J	—	—	21.8	0.06
2459089.9585	2.27	FNT+LCO	r	—	—	23.30	0.15
2459089.9698	2.28	FNT+LCO	g	—	—	23.41	0.21
2459090.9200	3.23	LDT+LMI	r	24.46	0.12	26.37	0.15
2459115.9675	28.28	Gemini+GMOS	i	25.45	0.15	25.9	0.11
2459115.9675	28.28	Gemini+GMOS	r	—	—	25.4	0.15
2459133.8039	46.11	Gemini+GMOS	i	—	—	25.5	0.11
2459133.8039	46.11	Gemini+GMOS	r	—	—	25.7	0.15

Extended Data Table 1: **Afterglow panchromatic observations.** Observations of the GRB 200826A afterglow and SN. GRB 200826A was triggered at Julian day 2459087.6874.

Julian day	δt	Instrument	Filter	Host AB magnitude	σ_{mag}	A_ν
2459089.62569	9.73	P200+WIRC	J	21.11*	0.16	0.05
2459105.80457	18.11	LMI+LDT	u	23.45	0.24	0.28
2459105.80457	18.11	LMI+LDT	g	23.36	0.05	0.22
2459105.80457	18.11	LMI+LDT	r	22.86	0.18	0.15
2459105.80457	18.11	LMI+LDT	i	22.66	0.16	0.11
2459105.80457	18.11	LMI+LDT	z	22.13	0.05	0.09

Extended Data Table 2: **Host galaxy panchromatic data.** Observations of the host galaxy of GRB 200826A. *Magnitudes are in the Vega system.

Julian day	δt	Instrument	Energy range	Count rate [s^{-1}]
2459088.386616	0.70	XRT	0.3–10 keV	$2.16_{-0.53}^{0.53} \times 10^{-2}$
2459088.441940	0.75	XRT	0.3–10 keV	$1.91_{-0.49}^{0.49} \times 10^{-2}$
2459088.522328	0.83	XRT	0.3–10 keV	$1.66_{-0.41}^{0.41} \times 10^{-2}$
2459089.437860	1.75	XRT	0.3–10 keV	$5.64_{-1.21}^{1.21} \times 10^{-3}$
2459090.285972	2.60	XRT	0.3–10 keV	$5.36_{-1.22}^{1.22} \times 10^{-3}$
2459094.090850	6.40	XRT	0.3–10 keV	$3.33_{-0.77}^{0.77} \times 10^{-3}$
2459104.352156	16.66	XRT	0.3–10 keV	$1.22_{-0.67}^{0.98} \times 10^{-3}$

Extended Data Table 3: **Afterglow X-ray detections.** X-ray observations of GRB 200826A.

Julian day	δt	Instrument	Frequency	Flux [$\text{erg cm}^{-2} \text{s}^{-1}$]
2459089.967407	2.28	VLA	6 GHz	40 **
2459102.153825	14.46	GMRT	1.256 GHz	< 48.6
2459107.485017	19.79	GMRT	1.256 GHz	< 57.4

Extended Data Table 4: **Radio data.** Radio observations of GRB 200826A. ** VLA data from Ref.¹².

Time Bins (s)	Model	Amplitude	E_{peak} [keV]	α	β	Photon Flux [$\text{ph s}^{-1} \text{cm}^{-2}$]	Photon Fluence [ph cm^{-2}]	Energy Flux [$\text{erg s}^{-1} \text{cm}^{-2}$]	Energy Fluence [erg cm^{-2}]	Fit Merit
0.000-0.180	Comp	0.3 ± 0.04	114.8 ± 6.0	-0.6 ± 0.1	-	32.2 ± 0.9	5.8 ± 0.2	$3.1 \pm 0.2 \times 10^{-6}$	$0.6 \pm 0.02 \times 10^{-6}$	1.1
0.180-0.318	Band	4.2 ± 2.3	61.3 ± 0.5	0.5 ± 0.3	-2.7 ± 0.1	41.5 ± 1.1	11.5 ± 0.2	$3.7 \pm 0.2 \times 10^{-6}$	$1.2 \pm 0.2 \times 10^{-6}$	0.9
0.318-0.414	Band	0.6 ± 0.1	141.9 ± 10.3	-0.5 ± 0.1	-3.0 ± 0.4	61.5 ± 1.5	17.4 ± 0.2	$7.8 \pm 0.4 \times 10^{-6}$	$1.9 \pm 0.3 \times 10^{-6}$	0.9
0.414-0.506	Band	0.7 ± 0.1	145.4 ± 13.5	-0.3 ± 0.1	-2.3 ± 0.2	61.4 ± 1.5	23.0 ± 0.1	$9.6 \pm 0.4 \times 10^{-6}$	$2.8 \pm 0.4 \times 10^{-6}$	1.0
0.506-0.607	Band	1.2 ± 0.4	99.6 ± 7.4	-0.04 ± 0.2	-2.5 ± 0.2	56.4 ± 1.4	28.8 ± 0.1	$6.9 \pm 0.3 \times 10^{-6}$	$3.5 \pm 0.3 \times 10^{-6}$	0.9
0.607-0.747	Band	0.9 ± 0.2	95.3 ± 5.3	-0.2 ± 0.1	-3.3 ± 0.4	41.7 ± 1.1	34.6 ± 0.2	$4.2 \pm 0.2 \times 10^{-6}$	$4.1 \pm 0.3 \times 10^{-6}$	0.9
0.747-1.152	Band	0.6 ± 0.5	41.2 ± 3.9	-0.3 ± 0.4	-2.7 ± 0.2	15.0 ± 0.5	40.7 ± 0.2	$0.9 \pm 0.1 \times 10^{-6}$	$4.5 \pm 0.3 \times 10^{-6}$	1.1

Extended Data Table 5: **Parameters of the *Fermi*-GBM fit.** Fitting parameter of the *Fermi*-GBM gamma-ray spectrum of GRB 200826A.

Line	F_{ν}
\AA	$10^{-17} \text{ erg s}^{-1} \text{ cm}^{-2}$
[OII] ₃₇₂₆	1.268 ± 0.37
[OII] ₃₇₂₉	2.522 ± 0.37
[OIII] ₅₀₀₇	8.21 ± 3.5

Extended Data Table 6: **Host galaxy emission line fluxes.** Fluxes derived with pPXF for the lines detected in the GTC spectrum of the host galaxy.

Parameter	Unit	Value
θ_v	[rad]	$0.20^{+0.53}_{-0.15}$
$E_{K,iso}$	[erg]	$6.0^{+51.3}_{-4.4} \times 10^{52}$
θ_c	[rad]	$0.24^{+0.53}_{-0.17}$
n	[cm ⁻³]	$5.5^{+187.3}_{-5.4} \times 10^{-2}$
p		$2.40^{+0.07}_{-0.07}$
ϵ_e		$0.42^{+0.40}_{-0.28}$
ϵ_B		$6.4^{+197.4}_{-5.2} \times 10^{-5}$
$E(B - V)$		$2.5^{+4.8}_{-2.3} \times 10^{-2}$
stretch s		$0.73^{+0.23}_{-0.16}$
scale k		$0.28^{+0.25}_{-0.12}$
E_k	[erg]	$1.4^{+54.8}_{-1.2} \times 10^{51}$

Extended Data Table 7: **Afterglow properties.** Posterior `afterglowpy` fit model parameters with an SMC extinction curve, a SN1998bw template and including the final Gemini+GMOS detection. Uncertainties are quoted at 90%. E_k is the beamed corrected kinetic energy. See §4 for more details.

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